

# ELECTROWEAK INTERACTIONS

④  
①

RECAP WHAT WE KNOW ABOUT  
QED (EM)

$$\psi(x) \rightarrow e^{i\alpha(x)} \psi(x)$$

$x, y, z, t$   
LOCAL GAUGE INVARIANCE

INVARIANT

UNDER  $U(1)$  1-D ROTATIONS

→ CONSERVED CHARGE — GLOBAL GAUGE INVARIANCE

$$\bar{\psi} (i\gamma^\mu \partial_\mu - m) \psi \leftarrow \text{DIRAC}$$

$$\downarrow$$
$$\bar{\psi} (i\gamma^\mu \partial_\mu - m) \psi + e \bar{\psi} \gamma^\mu \psi A_\mu$$

PHOTON  
EM FIELD.



MAYBE WEAK INTERACTION IS  
INVARIANCE UNDER ROTATIONS  
IN 3-d !  $\Rightarrow SU(2)$

• FOR WEAK INTERACTION

• LEFT HANDED DOUBLETS

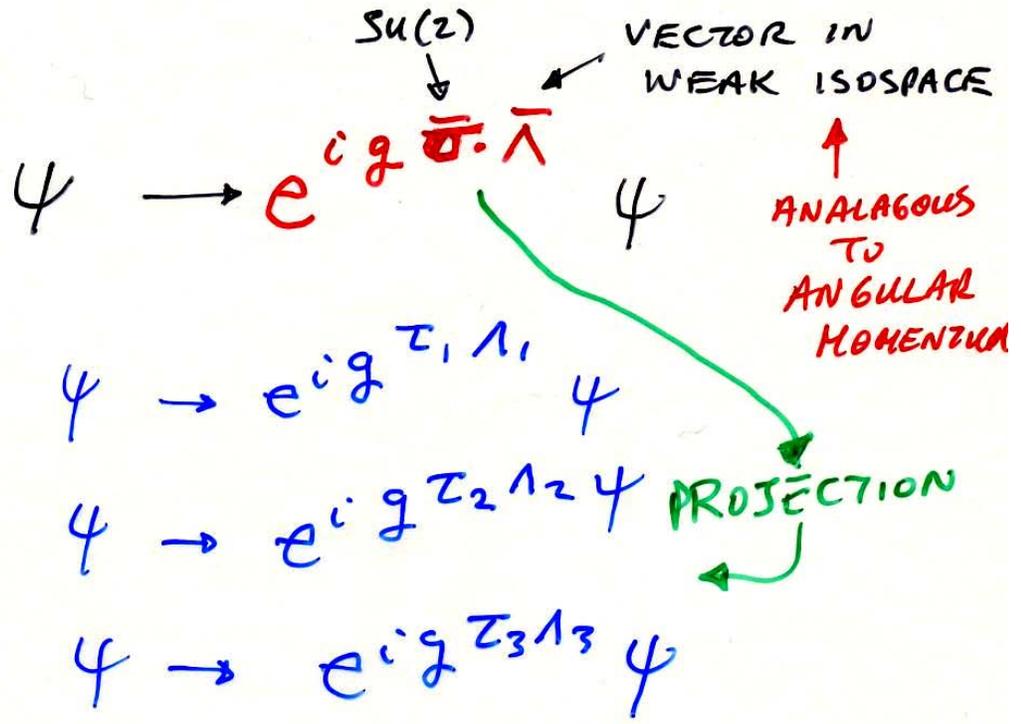
$$\begin{pmatrix} \nu_e \\ e \end{pmatrix} \quad \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix} \quad \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}$$

WEAK ISOSPIN UP AND DOWN ?

$$\begin{pmatrix} W^+ \\ Z^0 \\ W^- \end{pmatrix}$$

ISOTRIplet  
OF WEAK  
BOSONS ?

GUESS :



QED GAUGE INVARIANCE

$$D_\mu \equiv \partial_\mu - ie A_\mu$$

$\gamma_\mu$   
 PHOTON

$$A_\mu \rightarrow A_\mu + \frac{1}{e} \partial_\mu \alpha$$

WEAK GAUGE INVARIANCE

$$D_\mu \rightarrow \partial_\mu - ig \vec{\sigma} \cdot \vec{W}_\mu$$

$$\partial_\mu - ig \tau_i W_\mu$$

$W^+$   
 $W^0$   
 $W^-$

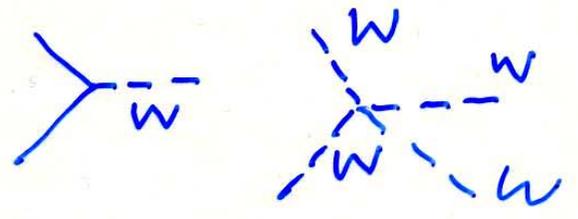
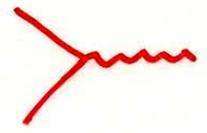
WEAK CHARGE

$$W_\mu \rightarrow W_\mu + \partial_\mu \bar{\Lambda} - g \frac{\bar{\Lambda} \times W_\mu}{\Lambda}$$

cf  $-ie A_\mu$        $-ig \tau_i W_\mu$

U(1) GENERATORS  
COMMUTE

SU(2) GENERATORS  
DO NOT COMMUTE

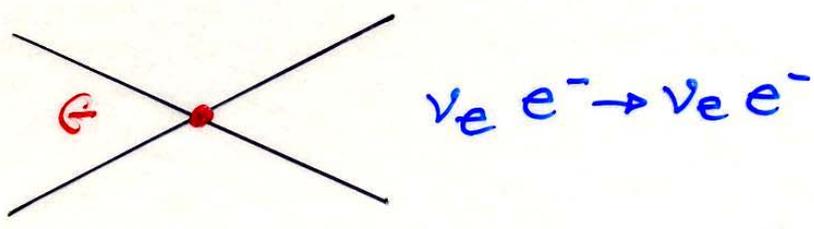


ABELIAN

NON ABELIAN

# DIVERGENCES IN WEAK INTERACTIONS

• IN THE 4-FERMION V-A WEAK THEORY



PERKINS CH.5

$$\int \frac{d\sigma}{dq^2} = \int \frac{G^2}{\pi}$$

FERMI WEAK COUPLING  
CM ENERGY SQUARED

$$\sigma_{TOT} = \frac{G^2}{\pi} q_{MAX}^2 = \frac{2 G^2 m E_\nu}{\pi} = \frac{G^2 s}{\pi}$$

∴  $\sigma \propto G \times$  DENSITY OF FINAL STATES

$\propto G \times s$

$\propto G \times p^{*2} \times 4$

RISES INDEFINITELY  
⋮?

UNITARITY

$$\sigma_{MAX} = \frac{\pi}{2} \left( \frac{\hbar}{p^*} \right)^2 \quad 1$$

$$\sigma_{\text{TOT}} = \frac{4G^2 p^{*2}}{\pi}$$

$$\sigma_{\text{MAX}} = \frac{\pi}{2p^{*2}}$$

WHEN DOES  $\sigma_{\text{TOT}}$  EXCEED  $\sigma_{\text{MAX}}$ ?

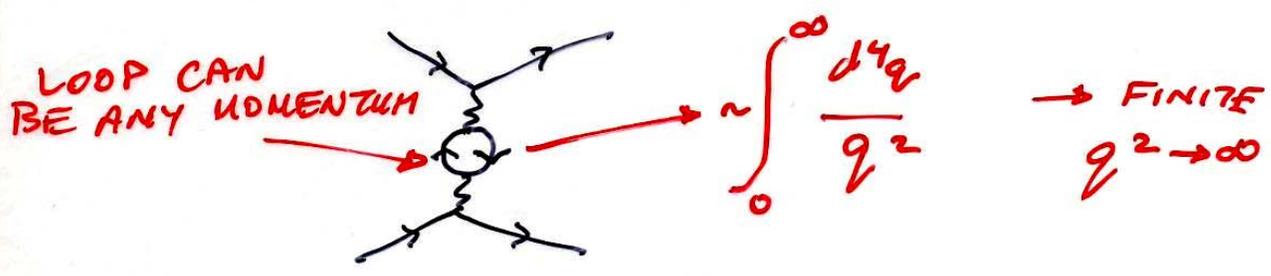
$$\frac{4G^2 p^{*2}}{\pi} > \frac{\pi}{2p^{*2}}$$

$$p^* > \left( \frac{\pi}{G} \cdot \frac{1}{\sqrt{8}} \right)^{\frac{1}{2}} \approx 300 \frac{\text{GeV}}{c}$$

- AT 300 GeV/c HAVE UNITARITY VIOLATION
  - PROBABILITY NOT CONSERVED
  - SCATTERED FLUX > INCIDENT FLUX
- ⇒ CANNOT BE TRUE.

# RENORMALIZATION & QED

- RENORMALIZATION IS A BIT OUT OF OUR SCOPE
- IN QED EVERY LOOP  $\rightarrow \infty$



- $\infty$  CAN BE CONTROLLED IN QED  
 ↳ ABSORBS INTO PHYSICAL QUANTITIES  
 $m_e, e$
- RENORMALIZATION IS A PROPERTY OF GAUGE INVARIANT THEORIES.

↓

INTIMATELY CONNECTED TO  
 CANCELLATIONS INHERENT  
 IN GAUGE INVARIANCE

# POINT 4-FERMI WEAK

(7-3)

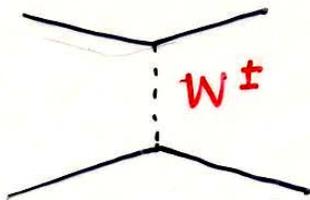
- BAD HIGH ENERGY BEHAVIOUR

WORSE AS HIGHER ORDER IN G

$$[G] = [E]^{-2}$$



- INTRODUCE  $W^\pm$



$$\frac{1}{\left(1 + \frac{q^2}{M_W^2}\right)}$$

→ NO LONGER POINT INTERACTION

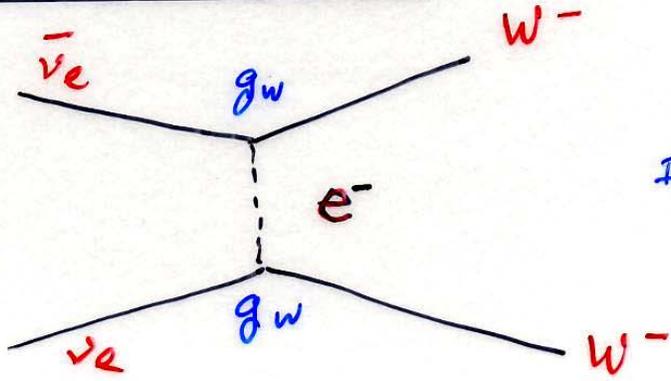
$$\sigma_{TOT} \rightarrow \frac{G^2 M_W^2}{\pi} \quad \text{FOR S LARGE}$$

- STILL HAVE DIVERGENCES IN HIGHER ORDER - LOGARITHMIC

- NOT A GAUGE THEORY, NOT RENORMALIZABLE

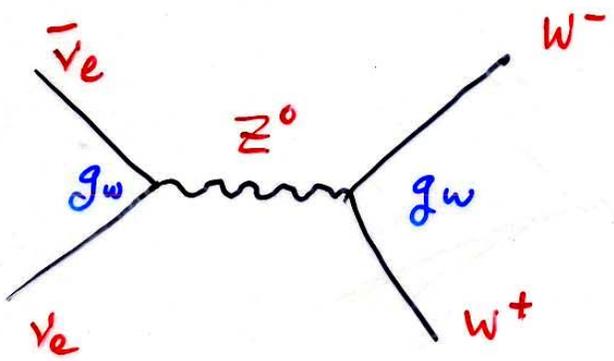
↳ NEED SOME CANCELLATION MECHANISM

# NEUTRAL CURRENTS

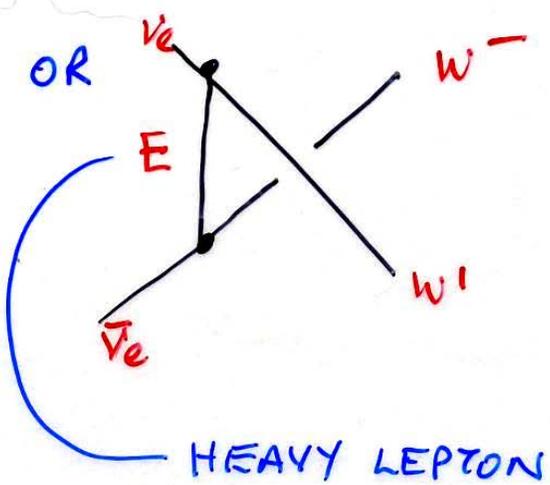


DIVERGES

CANCEL WITH



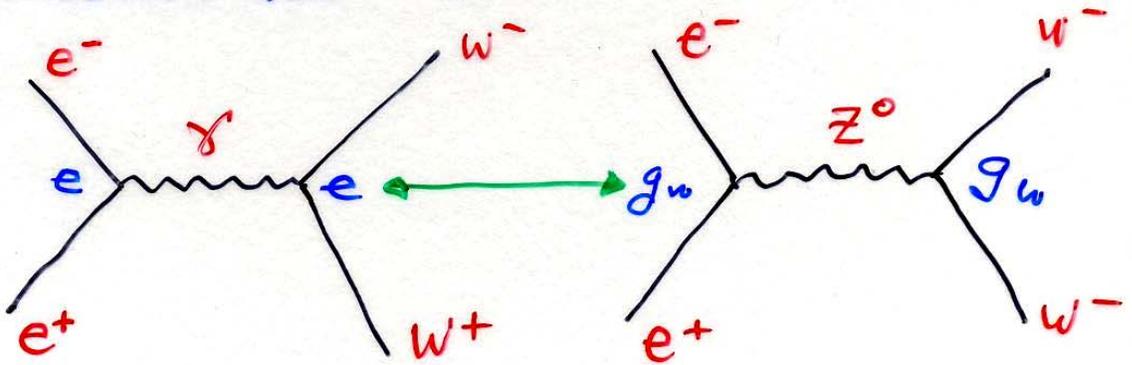
NEUTRAL CURRENT  
IN GIM  
MECHANISM



HEAVY LEPTON  
WITH ELECTRON  
FLAVOUR

↑  
NO SIGN OF  
THIS

CANCELLATION OF



OBSERVE THAT THERE WILL BE SOME MIXING  $\gamma \leftrightarrow Z^0 \Rightarrow g_w$

• FOR CANCELLATION  $e \sim g_w$

RECALL

$$\lim_{q^2 \rightarrow 0} \frac{g_w^2}{(q^2 + M_w^2)} \equiv \frac{G}{\sqrt{2}}$$

PUT  $g_w = e$

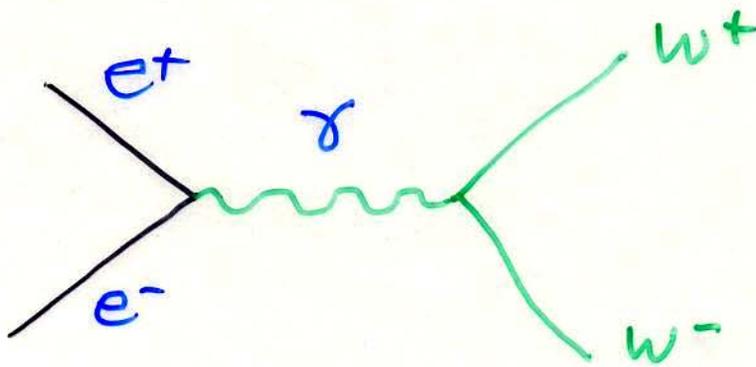
$$M_w^2 = \frac{g_w^2 \sqrt{2}}{G} = \frac{e^2 \sqrt{2}}{G}$$

$$\Rightarrow M_{w, Z^0} \approx 100 \text{ GeV}$$

BUT STILL HAVE DIVERGENCES DUE TO LEPTON MASSES  $\neq 0$

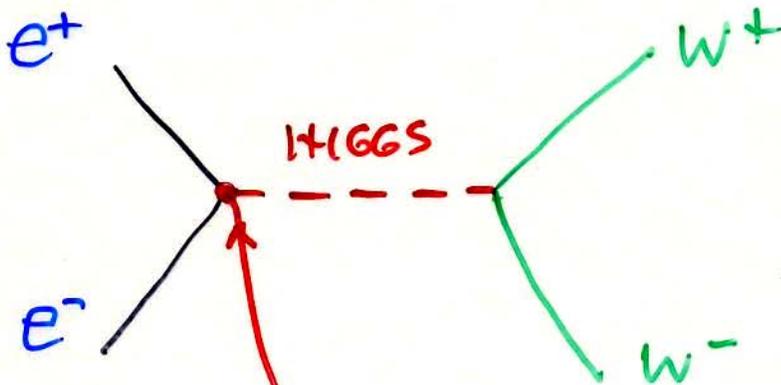
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# WHAT ABOUT HEAVY LEPTON?



DIVERGES  $\sim \frac{1}{M}$   
 ( $m_e \neq 0$ )

IN STANDARD ELECTRO WEAK MODEL  
 THIS DIVERGENCE EXACTLY  
 CANCELLED BY :



HIGGS PREDI  
 FOR MASSIVE  
 $W, Z$

THEORY  $-im_e \sqrt{\sqrt{2}G_F}$

THIS CANCELLATION ONLY WORK  
 IN STANDARD MODEL

HEAVY LEPTON WOULD HAVE TO  
 "JUST HAPPEN" TO HAVE CORRE

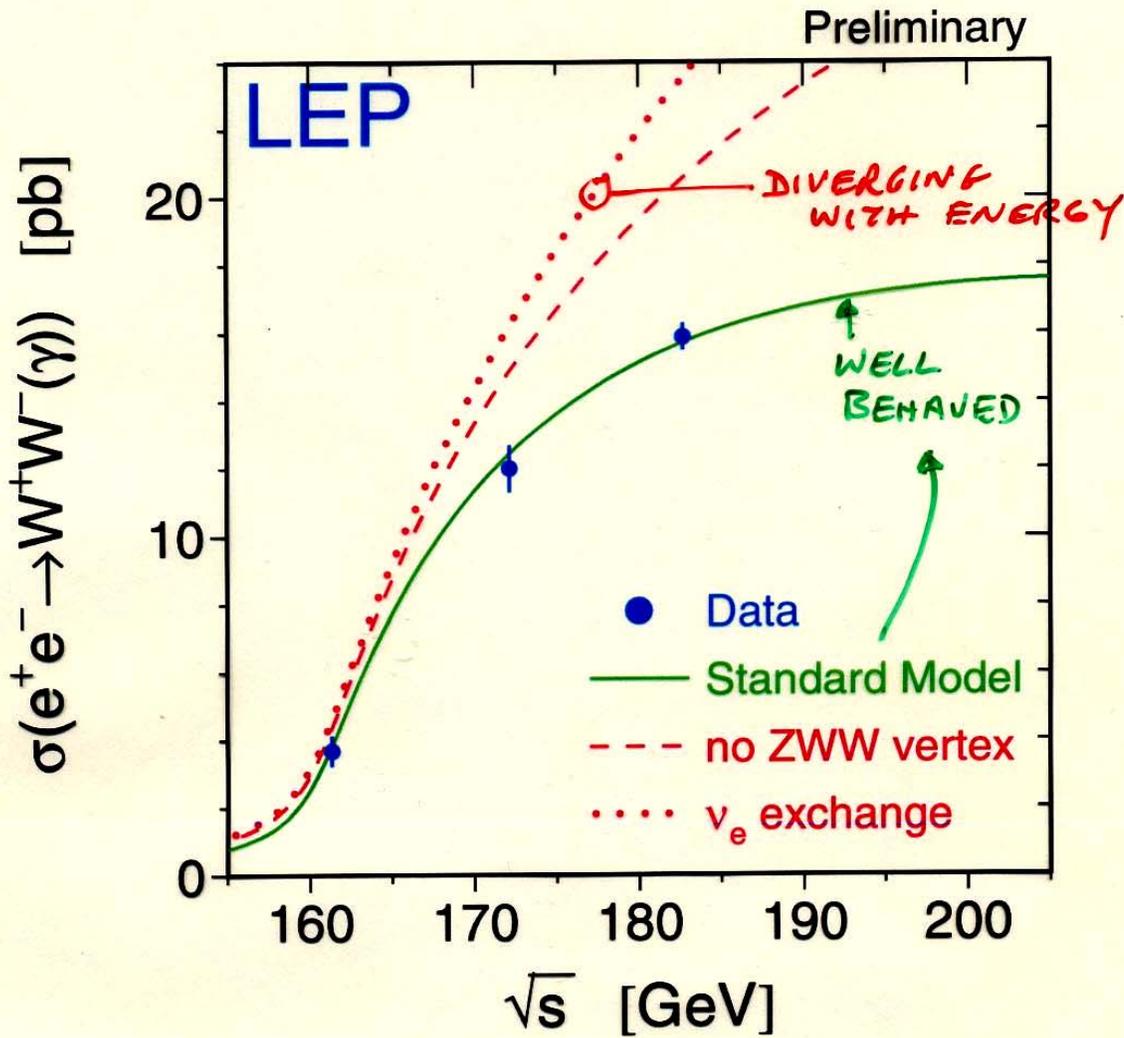


Figure 5: The W-pair cross-section as a function of the centre-of-mass energy. The data points are the LEP averages. Also shown is the Standard Model prediction (solid line), and for comparison the cross-section if the ZWW coupling did not exist (dashed line), or if only the  $t$ -channel  $\nu_e$  exchange diagram existed (dotted line).

The W decay branching ratios have been determined from individual channels cross-section values measured by the four experiments (at all centre-of-mass energies [91–94]), with and without lepton universality assumption and are shown in Table 24. Correlated errors between the individual channels have been taken into account.  $B(W \rightarrow \tau\nu)$  is 25% anti-correlated with the other two leptonic branching ratios,  $B(W \rightarrow \mu\nu)$  and  $B(W \rightarrow e\nu)$  are 1.0% correlated. Under the assumption of lepton universality,