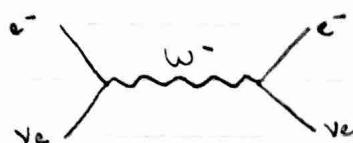


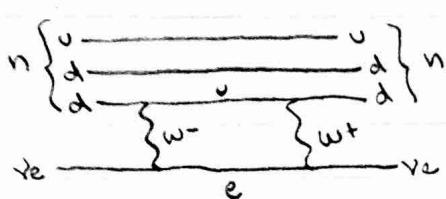
## Question #2

a) Quark line diagrams (pr):

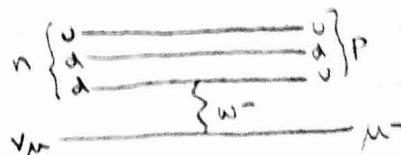
(1)  $\nu_e e^- \rightarrow \nu_e e^-$



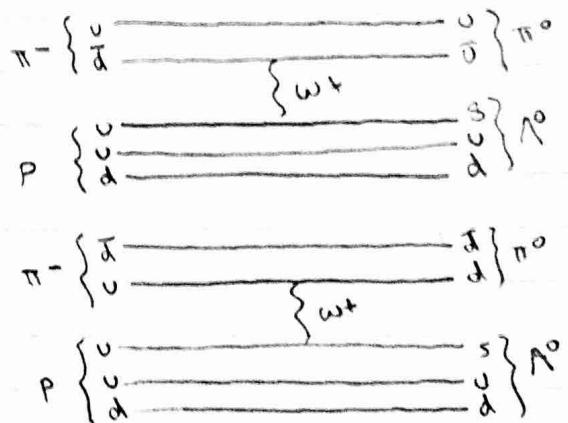
(2)  $\nu_e n \rightarrow \nu_e n$



(3)  $\nu_\mu n \rightarrow \mu^- p$

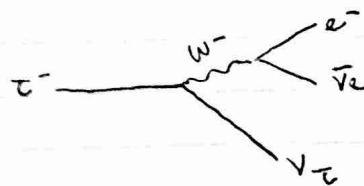


(4)  $\pi^- p \rightarrow \Lambda \pi^0$

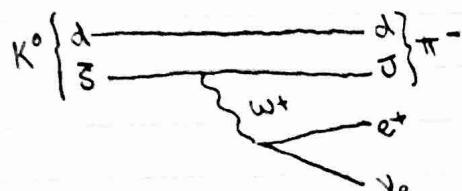


b) Quark line diagrams (pr):

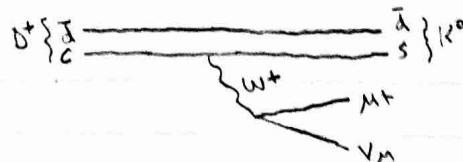
(1)  $\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$



(2)  $K_0 \rightarrow \pi^- e^+ \nu_e$



(3)  $D^+ \rightarrow \bar{K}^0 \mu^+ \nu_\mu$



(4)  $\tau^+ \rightarrow \pi^+ \nu_\tau$

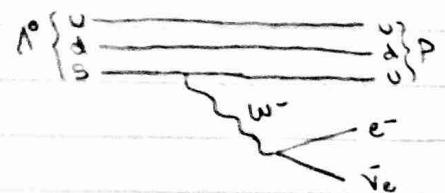


## Problem Set #4

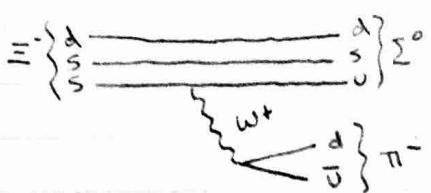
Physics 357S

Question #2 cont

$$15) N^0 \rightarrow p e^- \bar{\nu}_e$$



$$\omega \Xi^- \rightarrow \Sigma^0 \pi^-$$



## Question #2

$$(1) \pi^- + p \rightarrow \pi^- + \pi^+ + n$$

$d\bar{d} + uud \rightarrow d\bar{d} + u\bar{d} + uud \rightarrow$  strong force (flavour/baryon  $\ell$  conserved)

$$(2) \gamma + p \rightarrow \pi^0 + p$$

$\gamma + uud \rightarrow u\bar{d}/d\bar{d} + uud \rightarrow$  EM force since  $\gamma \rightarrow q\bar{q}$  is EM

$$(3) \bar{\nu}_\mu + p \rightarrow \mu^+ + n$$

$\bar{\nu}_\mu + uud \rightarrow \mu^+ + \bar{u}d\bar{d} \rightarrow$  must be weak force b/c we have a neutrino

$$(4) \pi^0 \rightarrow e^+ + e^- + e^+ + e^-$$

$d\bar{d} \rightarrow e^+ + e^- + e^+ + e^- \rightarrow$  EM since  $\pi^0 \rightarrow \gamma\gamma \rightarrow e^+e^-e^+e^-$

$$(5) p + \bar{p} \rightarrow \pi^+ + \pi^- + \pi^0$$

$uud + \bar{s}\bar{d} \rightarrow u\bar{d} + \bar{s}d + u\bar{d}/d\bar{d} \rightarrow$  strong force

$$(6) \tau^- \rightarrow \pi^- + \bar{\nu}_\tau$$

$\tau^- \rightarrow \bar{u}d + \bar{\nu}_\tau \rightarrow$  weak force

$$(7) D^- \rightarrow K^+ + \pi^- + \pi^-$$

$d\bar{d} \rightarrow u\bar{s} + \bar{d}s + \bar{d}s \rightarrow$  weak since strange/dcharm not conserved

$$(8) \pi^- \rightarrow \pi^0 + e^- + \bar{\nu}_e$$

$\bar{s}d \rightarrow u\bar{d}/d\bar{d} + e^- + \bar{\nu}_e \rightarrow$  weak

$$(9) \Lambda^+ + p \rightarrow K^- + \rho^+ + p$$

$uds + uud \rightarrow s\bar{d} + uud + uud \rightarrow$  strong

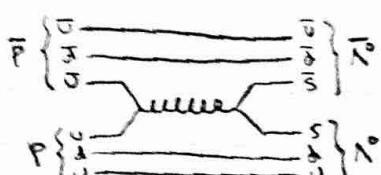
## Question #3

a)  $\bar{p} + p \rightarrow A + \Lambda$

$\bar{u}\bar{d} + uud \rightarrow A + uds \rightarrow$  strong force (see part b)

b)

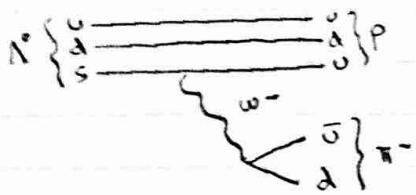
	$\bar{p}$	$p$	$\Lambda$	$A$
charge	-1	+1	0	0
baryon	-1	+1	+1	-1



The anti-lambda conserves charge and baryon number. Strangeness is also conserved.

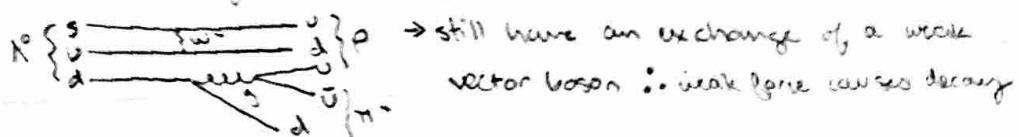
c) The  $\Lambda^0$  decays into a  $\pi^-$  and a proton.

The  $\bar{\Lambda}^0$  decays into a  $\pi^+$  and anti-proton.



so the weak force causes the neutral particle decay at the secondary vertices. We can infer this because of the displaced vertices - due to the longer lifetimes, the neutral particles had time to travel in the bubble chamber before decaying.

Also can have (from class notes):



note this conclusion holds for  $\bar{\Lambda}^0$  as well, Feynman diagrams are analogous.

d) See part b) for the quark line diagram.

e) See the Feynman diagrams in part c), but take the antiquarks.

Question #3 cont...

f) interaction at B:  $\bar{p} + p \rightarrow \pi^+ + \pi^- + \pi^+ + \pi^-$ 

$$\bar{u}\bar{d} + u\bar{d} \rightarrow \bar{u}\bar{d} + \bar{d}d + u\bar{d} + \bar{d}d$$

(where) (baryon #: (-2)(-1) (+1)(+1) (+1)(0) (-1)(0) (+1)(0) (-1)(0))

Baryon # and charge are conserved. No surprising new quarks.  $\Rightarrow$  strong interaction.

g) The rest mass of  $\Lambda$  can be determined based on the momenta of the charged particles, which can be determined from the curvature of their tracks.(i) To get the momenta:  $m\vec{v} = q\vec{J} \times \vec{B}$ 

$$\frac{mv^2}{R} = qvB$$

$$mv = qBR = \vec{p}_\perp$$

(ii) To get the mass:  $(p_\Lambda^m)^2 = (p_p^m + p_{\pi^+})^2$ 

$$m_\Lambda^2 = (E_p + E_{\pi^+})^2 - (\vec{p}_p + \vec{p}_{\pi^+})^2$$

$$m_\Lambda^2 = (E_p + E_{\pi^+})^2 - (\vec{p}_p^2 + \vec{p}_{\pi^+}^2 + 2\vec{p}_p \cdot \vec{p}_{\pi^+})$$

$$m_\Lambda^2 = (E_p + E_{\pi^+})^2 - (\vec{p}_p^2 + \vec{p}_{\pi^+}^2 + 2|\vec{p}_p||\vec{p}_{\pi^+}| \cos\theta)$$

angle between  
p and  $\pi^+$

h) The distribution expected from observing many neutral decays, given that the mean lifetime of  $\Lambda$  is  $2.6 \times 10^{-10}$  s and the precision of a bubble chamber is  $\sim$  MeV.

The width of a decay is given by  $\Gamma = \frac{\lambda}{T} = \frac{\lambda}{2.6 \times 10^{-10}} \text{ s}$   
 $= 2.53 \times 10^{-6} \text{ eV}$

$\therefore$  since the width of  $\Lambda$  is much smaller than the resolution of the bubble chamber, we expect a mass peak with a width of  $\sim$  MeV.

## Question #4

a) General expression for form factor in spherical coordinates:  $F(q^2) = \frac{1}{2e} \int \rho(r) e^{iq^2 r/\hbar} dV$   
 where  $dV = r^2 \sin\theta dr d\theta d\phi$

$$= r^2 \cos\theta d\phi dr \quad (\text{for spherically-symmetric distribution})$$

Then,

$$\begin{aligned} F(q^2) &= \int_0^{2\pi} d\phi \int_{-\pi}^{\pi} \int_0^{\infty} \left(\frac{1}{2e}\rho(r)\right) e^{iq^2 r/\hbar} r^2 \cos\theta dr \\ &= \frac{2\pi}{2e} \int \rho(r) e^{iq^2 r/\hbar} r^2 \cos\theta dr \\ &= \frac{2\pi}{2e} \int \rho(r) \left(\frac{1}{2} \sin(q^2 r/\hbar)\right) (e^{iq^2 r/\hbar} - e^{-iq^2 r/\hbar}) r^2 dr \\ &= \frac{4\pi}{2e} \int \rho(r) \left(\frac{1}{2} \sin(q^2 r/\hbar)\right) \sin(q^2 r/\hbar) r^2 dr \end{aligned}$$

Expanding the sine gives:

$$\begin{aligned} \sin(q^2 r/\hbar) \left(\frac{1}{2} \sin(q^2 r/\hbar)\right) &= \left(\frac{q^2 r}{\hbar} - \frac{q^4 r^3}{6\hbar^3} + \dots\right) \left(\frac{q^2 r}{\hbar}\right) \\ &= 1 - \left(\frac{1}{6}\right) \left(\frac{q^2 r}{\hbar}\right)^2 + \dots \end{aligned}$$

Then, we write the form factor as

$$F(q^2) = \frac{4\pi}{2e} \int r^2 \rho(r) \left(1 - \left(\frac{1}{6}\right) \left(\frac{q^2 r}{\hbar}\right)^2 + \dots\right) dr$$

Imposing the normalization condition, we define the root-mean-square radius:

$$\sqrt{\frac{4\pi}{2e} \int r^2 \rho(r) dr} = 1 \quad \text{and} \quad \langle r^2 \rangle = \frac{4\pi}{2e} \int r^2 \rho(r) r^2 dr$$

∴ the form factor can be written as,

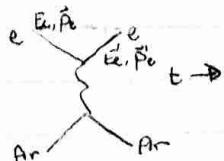
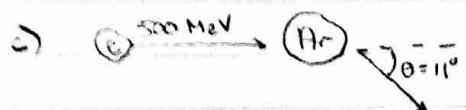
$$\begin{aligned} F(q^2) &= \frac{4\pi}{2e} \int r^2 \rho(r) dr - \frac{4\pi}{2e} \int r^2 \rho(r) r^2 \left(\frac{1}{6}\right) \left(\frac{q^2 r}{\hbar}\right)^2 dr \\ &= 1 - \langle r^2 \rangle \left(\frac{1}{6}\right) \left(\frac{q^2 r}{\hbar}\right)^2 \\ &= 1 - \frac{q^2 \langle r^2 \rangle}{6R^2} + O(q^4) \end{aligned}$$

b) Now we have the charge distribution  $\rho(r) = \rho_0 e^{-r/a} \left(\frac{1}{r}\right)$

As in part a), for spherically-symmetric charge distributions we have:

$$\begin{aligned} F(q^2) &= \frac{4\pi}{2e} \int r^2 \rho(r) \sin(q^2 r/\hbar) \left(\frac{1}{r}\right) dr \quad \text{normalization} \\ &= \frac{4\pi}{2e} \int r^2 \rho(r) \sin(q^2 r/\hbar) \left(\frac{1}{r}\right) q^2 r dr \times \left[ \frac{1}{2e} \int r^3 dr \rho(r) \right]^{-1} \\ &= \frac{4\pi}{2e} \int r^2 \rho_0 e^{-r/a} \left(\frac{1}{r}\right) \sin(q^2 r/\hbar) \left(\frac{1}{r}\right) q^2 r dr \times \left[ \frac{4\pi}{2e} \int r^2 dr \rho_0 e^{-r/a} \left(\frac{1}{r}\right) \right]^{-1} \\ &= \frac{4\pi}{2e} \rho_0 \left(\frac{1}{a}\right) \int_0^\infty e^{-r/a} \sin(q^2 r/\hbar) dr \times \left[ \frac{4\pi}{2e} \rho_0 \int_0^\infty r e^{-r/a} dr \right]^{-1} \\ &= \frac{4\pi}{2e} \int_0^\infty e^{-r/a} \sin(q^2 r/\hbar) dr \times \left[ \int_0^\infty r e^{-r/a} dr \right]^{-1} \\ &= \frac{4\pi}{2e} \left[ \frac{q^2 r}{a(a^2 + (q^2 r/\hbar)^2)} \right] \times [a^2]^{-1} \\ &= (1 + (q^2 r/\hbar)^2)^{-1} \end{aligned}$$

Question #11 cont'd...



(i) Momentum transfer:  $q^2 = (E_e + E'_e)^2 - (\vec{p}_e + \vec{p}'_e)^2 \rightarrow (\text{upper vertex})$   
 $= 2m_e^2 - 2E_e E'_e + 2\vec{p}_e \cdot \vec{p}'_e \cos\theta \quad \textcircled{*}$

$p_{Ar}^2 = (\vec{p}_{Ar} + q)^2 \rightarrow (\text{lower vertex})$

$m_{Ar}^2 = m_{Ar}^2 + q^2 + 2\vec{p} \cdot \vec{q}$

$\Rightarrow q^2 = m_{Ar}^2 - m_{Ar}^2 - 2\vec{p} \cdot \vec{q}$

$= -2\vec{p} \cdot \vec{q}$

$= -2(m_{Ar}, \vec{0}) \cdot (E_e - E'_e, \vec{p}_e - \vec{p}'_e)$

$= -2m_{Ar}(E_e - E'_e) \quad \textcircled{**}$

$\textcircled{*}$  and  $\textcircled{**}$  give,

$E'_e = -q^2 / (2E_e(1-\cos\theta)) = q^2 / (2m_{Ar}) + E_e$

$\Rightarrow q^2 = -E_e \left[ \frac{1}{2m_{Ar}} + \frac{1}{2E_e(1-\cos\theta)} \right]^{-1}$

$= -500 \text{ MeV} \left[ \frac{1}{2 \times 40 \times 938} + \frac{1}{2 \times 500 \times (1-\cos 11^\circ)} \right]^{-1}$

$\approx -9189 \text{ MeV}^2$

(ii) Reduced deBroglie wavelength:  $\frac{\hbar c}{pc} = \frac{197 \text{ MeV fm}}{500 \text{ MeV}} \approx 0.394 \text{ fm}$

(iii) Mott differential cross section:  $\left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} = 4(2\alpha)^2 \frac{E^2}{q^2} (1 - \beta^2 \sin^2 \theta/2)$   
 $= 4(18 \cdot \frac{1}{137})^2 (500 \text{ MeV})^2 / (9189 \text{ MeV}^2) (1 - \sin^2 11^\circ) \text{ (assume } \beta \approx 1)$   
 $\approx 7.86 \text{ fm}^2$   
 $\approx 78.6 \text{ mb}$

(iv) Change in  $\frac{d\sigma}{d\Omega}$  is found using form factor of part (a):  $f(q^2) = 1 - \frac{1}{6} \frac{q^2}{\pi^2} \langle r^2 \rangle$   
 $= 1 - \frac{1}{6} (-9189 \text{ MeV}^2) (1.2 \times 40^{1/3})^2 (\frac{1}{197^2})$   
 $\approx 1.66$

$\therefore$  the cross section increases by  $\times 1.66^2 = 2.76$

## Question #5

$$\text{a) B.E.} = -a_1 A + a_2 A^{2/3} + a_3 \frac{Z^2}{A^{1/3}} + a_4 \frac{(N-Z)^2}{A} \pm a_5 A^{-3/4}$$

$$M(A, Z) c^2 = (A-Z)m_n c^2 + Zm_p c^2 - a_1 A + a_2 A^{2/3} + a_3 \frac{Z^2}{A^{1/3}} + a_4 \frac{(N-Z)^2}{A} \pm a_5 A^{-3/4}$$

Differentiate  $M(A, Z)$  for fixed  $A$ :

$$\frac{dM}{dZ} = -m_n + m_p + a_3 \frac{2Z}{A^{1/3}} + a_4 \left(\frac{1}{A}\right) 2(A-2Z) (-2)$$

$$0 = -m_n + m_p + 2a_3 \left(\frac{1}{A^{1/3}}\right) Z - a_4 \left(\frac{4}{A}\right) (A-2Z)$$

$$m_n - m_p = 2a_3 \left(\frac{1}{A^{1/3}}\right) Z - a_4 \left(\frac{4}{A}\right) A + a_4 \left(\frac{4}{A}\right) 2Z$$

$$m_n - m_p + 4a_4 = [2a_3 \left(\frac{1}{A^{1/3}}\right) + 8a_4 \left(\frac{1}{A}\right)] Z$$

$$\Rightarrow Z = \frac{m_n - m_p + 4a_4}{2a_3/A^{1/3} + 8a_4/A}$$

For the plot of  $M$  vs  $Z$ , we can see that for constant  $A$ ,  $M(Z)$  would be a parabola since we have  $Z^2$  terms in the semi-empirical mass formula.

$$\text{For } A=16; Z = \frac{m_n - m_p + 4a_4}{2a_3/A^{1/3} + 8a_4/A}$$

$$\text{note: } a_4 = 23.70 \text{ MeV}$$

$$a_3 = 0.71 \text{ MeV}$$

$$= \frac{1.293 \text{ MeV} + 4(23.70 \text{ MeV})}{2(0.71 \text{ MeV})/(16)^{1/3} + 8(23.70 \text{ MeV})/16}$$

$$\approx 8 \quad \therefore \text{the isotope is } {}^{16}\text{O}, \text{ a stable isotope of Oxygen.}$$

Similarly, for  $A=208$ ,  $Z \approx 83$   $\therefore$  the isotope is  ${}^{208}\text{Bi}$ . This is off by 1 from Bi's "most stable" isotope of  ${}^{209}\text{Bi}$  (also could be close to  ${}^{208}\text{Pb}$ ).

b) Total binding energy using equation for B.E. given above. ✓ pairing term  $\begin{cases} + & Z, N \text{ even} \\ 0 & A \text{ odd} \\ - & Z, N \text{ odd} \end{cases}$

$$\text{For } {}^8\text{Be}, \text{B.E.} = -a_1 A + a_2 A^{2/3} + a_3 \frac{Z^2}{A^{1/3}} + a_4 \frac{(N-Z)^2}{A} \pm a_5 A^{-3/4}$$

$$= -(5.8)(8) + (18.3)(8^{2/3}) + (0.714)(4^2)(8^{-1/3}) + (23.2)\left(\frac{1}{8}(4-4)^2\right) + 12(8^{-3/4})$$

$$\approx -44.96 \text{ MeV}$$

Similarly, for  ${}^{12}\text{C}$ , B.E.  $\approx -80.59$  MeV

${}^{56}\text{Fe}$ , B.E.  $\approx -483.57$  MeV

${}^{208}\text{Pb}$ , B.E.  $\approx -1617.53$  MeV

note: used slightly different values

for "a"s, but your values should be close.

Question #5 cont...

Binding energies using atomic masses of the nuclides are:

$$\begin{aligned} \text{for } {}^4\text{Be, } BE &= \Delta m = (N m_n + 2 m_p) - (m_{Be}) \\ &= (4 \cdot m_n + 4 \cdot m_p) - m_{Be} \\ &\approx 8.06642 \text{ u} - 8.005305 \text{ u} \\ &\approx 0.061115 \text{ u} \\ &\approx 56.93 \text{ MeV} \end{aligned}$$

Similarly,  ${}^{12}\text{C, } BE \approx 92.8 \text{ MeV}$

$${}^{56}\text{Fe, } BE \approx 495.0 \text{ MeV}$$

$${}^{208}_{82}\text{Pb, } BE \approx 1645.2 \text{ MeV}$$

## Question #6

a) Electrostatic potential energy of a uniformly charged sphere,

$$\text{For } r < R, \text{ the amount of charge is } q(r) = Qr^3/R^3$$

$$\text{For } r > R, \text{ the amount of charge is } q(r) = Q$$

$$\text{The electric field is given by: } \vec{E}(r) = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$

$$\text{For } r < R, \vec{E}(r) = \frac{1}{4\pi\epsilon_0} \frac{Qr}{R^3} \hat{r}$$

$$\text{For } r > R, \vec{E}(r) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \hat{r}$$

$$\text{Then the electrostatic potential is given by: } V = \frac{\epsilon_0}{2} \int E^2 dV$$

$$= 2\pi\epsilon_0 \int E^2 r^2 dr$$

$$\begin{aligned} \text{Integrate from } \{0, R\} \text{ and } \{R, \infty\}, V &= \frac{1}{8\pi\epsilon_0} \left[ \int_0^R \frac{Q^2}{R^6} r^4 dr + \int_R^\infty \frac{Q^2}{r^2} dr \right] \\ &= \frac{1}{8\pi\epsilon_0} \left[ \frac{Q^2}{R^6} \left( \frac{1}{5} R^5 \right) + \frac{Q^2}{r} \Big|_R^\infty \right] \\ &= \frac{1}{8\pi\epsilon_0} \left[ \frac{Q^2}{5R} + \frac{Q^2}{R} \right] \\ &= \frac{3Q^2}{20\pi\epsilon_0 R} \end{aligned}$$

b) Let  $Q = Ze$ , then using the Coulomb term of the semi-empirical mass formula we have,

$$ac^2/A^{1/3} = \frac{3Ze^2}{20\pi\epsilon_0 R}$$

$$\Rightarrow ac = 3A^{1/3}e^2 (20\pi\epsilon_0 \cdot 1.24 \text{ fm})^{1/3}$$

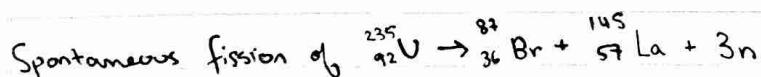
$$= 3e^2 (20\pi\epsilon_0 \cdot 1.24 \text{ fm})^{1/3}$$

$$= \frac{3}{5} \left( \frac{e^2}{4\pi\epsilon_0} \right) (1.24 \text{ fm})$$

$$= \frac{3}{5} \left( \frac{10^{-2}}{137} \text{ MeV fm} \right) (1.24 \text{ fm})$$

$$= 0.696 \text{ MeV}$$

Solving the BE formula for  $ac$  in the case of  ${}^{181}\text{Ta}$  gives  $ac \approx 0.693 \text{ MeV}$  - not bad!



$$\begin{aligned} \text{The energy released can be estimated: } E_{\text{released}} &= BE({}^{235}\text{U}) - BE({}^{87}\text{Br}) - BE({}^{145}\text{La}) \\ &\approx 153 \text{ MeV} \end{aligned}$$

Note: binding energy of 3 neutrons are zero, the binding energies of the rest of the nuclei are found using the formula from S(a).