LECTURE 23: CP Violation (Part II)

Overview:

- Meson mixing (cont.)
- K meson system
- B meson system

(I used mostly Burgess, and Cheng Li as references, and notes from Frank Wuerthwein)

Meson Mixing (continued from lecture 22)

We define the states $|\varphi\rangle$ and $|\bar{\varphi}\rangle$ (destroyed by the fields $\varphi(x)$ and $\bar{\varphi}(x)$) and define $|\psi_\pm\rangle$ (destroyed by $\psi(x)$ and $\bar{\psi}(x)$). The normalized states are:

$$|\psi_+\rangle = \frac{1}{\sqrt{|p|^2+|q|^2}}\left[p|\varphi\rangle + q|\bar{\varphi}\rangle\right]$$

$$|\psi_-\rangle = \frac{i}{\sqrt{|p|^2+|q|^2}}\left[p|\varphi\rangle - q|\bar{\varphi}\rangle\right]$$

$$\frac{p}{q} = z^2 = \sqrt{\frac{C}{B}}$$

If CP is a symmetry: $C = B \Rightarrow z = 1$

$$\Rightarrow \varphi^+ = \frac{\varphi + \bar{\varphi}}{\sqrt{2}}, \qquad \varphi^- = \frac{i(\bar{\varphi} - \varphi)}{\sqrt{2}}$$

$$M^2_\pm \simeq m^2_\pm - i m \Gamma_\pm = A \pm B$$

Meson Mixing (cont.)

eigenstates: $$|\psi_+\rangle = \frac{1}{\sqrt{1+|\hat{\varepsilon}|^2}} \left[|\varphi_+\rangle - i\hat{\varepsilon}\, |\varphi_-\rangle \right]$$

$$|\psi_-\rangle = \frac{1}{\sqrt{1+|\hat{\varepsilon}|^2}} \left[|\varphi_-\rangle + i\hat{\varepsilon}\, |\varphi_+\rangle \right]$$

with $\hat{\varepsilon} = \frac{(z^2-1)}{1+z^2} = \frac{p-q}{p+q}$

If we start with state $|\varphi\rangle$:

$$\langle \varphi | \varphi \rangle_t = \frac{1}{2}\left[e^{-iE_+t} + e^{-iE_-t} \right], \quad \langle \bar{\varphi} | \varphi \rangle = \frac{q}{2p}\left[e^{-iE_-t} - e^{-iE_+t} \right]$$

$$P_r\left[\varphi(k) \to \varphi(k)\right] = \frac{1}{4}\left[e^{-\Gamma_+(k)t} + e^{-\Gamma_-(k)t} + 2e^{-\left[\frac{\Gamma_+(k)+\Gamma_-(k)}{2}\right]t} \cos \Omega_k t \right]$$

$$P_r\left[\varphi(k) \to \bar{\varphi}(k)\right] = \frac{1}{4}\left|\frac{q}{p}\right|^2 \left[e^{-\Gamma_+(k)t} + e^{-\Gamma_-(k)t} - 2e^{-\left[\frac{\Gamma_+(k)+\Gamma_-(k)}{2}\right]t} \cos \Omega_k t \right]$$

$$\Omega_k = \sqrt{k^2+m_+^2} - \sqrt{k^2+m_-^2} \approx m_+ - m_- \quad \text{if } k \ll m$$

$$\approx \frac{m_+^2 - m_-^2}{2k} \quad \text{if } k \gg m$$

$$\rightarrow \quad K\left[\sqrt{1+\frac{m_+^2}{K^2}} - \sqrt{1+\frac{m_-^2}{K^2}}\right]$$

Taylor $\quad (1+x)^{1/2} = 1 + \frac{x}{2}$

$$\Rightarrow K\left[1+\frac{m_+^2}{2K^2} - 1 - \frac{m_-^2}{2K^2}\right] = \frac{m_+^2 - m_-^2}{2K}$$

$K^0 - \bar{K}$ mixing (revisited)

With Kaons we will have $|K_S\rangle = |\psi_+\rangle$

$$|K_L\rangle = |\psi_-\rangle$$

As we saw, to a good approx. (since $|\hat{\varepsilon}| \ll 1$),

$\Rightarrow |K_L\rangle \simeq |K_-\rangle = |q_-\rangle$

$K^0 - \bar{K}^0$ mixing (cont.)

For Kaons, we have:

$m_K = 498$ MeV $\qquad (3 m_\pi = 420$ MeV$)$

$\Gamma_s^{-1} = 9 \times 10^{-11}$ sec , $\Gamma_L^{-1} = 5 \times 10^{-8} \Rightarrow \Gamma_s = 580\, \Gamma_L$

With $\Gamma_+ \gg \Gamma_-$ and assuming slow-moving Kaons:

$$P_t[K(k) \to K(k)] \approx \frac{1}{4} e^{-\Gamma_L(k)T} \left[1 + 2 e^{-[\Gamma_s(k) - \Gamma_L(k)]t/2} \cos(\Delta m t)\right]$$

$$P_t[K(k) \to \bar{K}(k)] \approx \frac{1}{4} e^{-\Gamma_L(k)T} \left|\frac{q}{p}\right|^2 \left[1 - 2 e^{-[\Gamma_s(k) - \Gamma_L(k)]t/2} \cos(\Delta m t)\right]$$

$\Delta m = 3.5 \times 10^{-6}$ eV $\qquad (m_{K_L} - m_{K_s}) > 0$

Frequency comparable to K_s lifetime (how would you measure K^0 vs $\bar{K}^0$?)

--

CP violation in Kaon decays

- CP violation can occur through mixing:

$$|K_L\rangle = \frac{1}{\sqrt{1+|\tilde{\varepsilon}|^2}}\left[\,|\varphi_-\rangle + i\tilde{\varepsilon}\,|\varphi_+\rangle\,\right]$$

With $\mathcal{L}_w$ denoting weak interaction Lagrangian ($\Delta S = \pm 1$) we write $\langle \pi\pi | \mathcal{L}_w | K_+ \rangle$

- CP violation can also occur at the <u>decay</u>:

$\langle \pi\pi | \mathcal{L}_w | K_- \rangle$ i.e. $\mathcal{L}_w$ itself breaks CP (known as "direct" CP violation.

To determine relative size of these contributions we can use the observables:

$$\eta_{00} = \frac{\langle \pi^0\pi^0 | \mathcal{L}_w | K_L \rangle}{\langle \pi^0\pi^0 | \mathcal{L}_w | K_S \rangle} \qquad \eta_{+-} = \frac{\langle \pi^+\pi^- | \mathcal{L}_w | K_L \rangle}{\langle \pi^+\pi^- | \mathcal{L}_w | K_S \rangle}$$

CP violation in Kaon decays

If mixing is sole source of CP violation: $\eta_{00} = \eta_{+-}$

Note that the Kaons decaying to two pions have either isospin 0 or 2. Amplitudes are:

$$\langle \pi^+ \pi^- | \mathcal{L}_w | K^0 \rangle = A_0 e^{i\xi_0} + A_2 e^{i\xi_2}$$

$$\langle \pi^+ \pi^- | \mathcal{L}_w | \bar{K}_0 \rangle = A_0 e^{-i\xi_0} + A_2 e^{-i\xi_2}$$

A_0, A_2 are CP-conserving strong interaction matrix elements for pion isospin channel, ξ_0, ξ_2 are the CP-violating phases due to $\mathcal{L}_w$ (assuming CP-violating $\mathcal{L}_w$)

$\rightarrow$ Physical decay rates prop. to $|A_0 e^{i\xi_0} + A_2 e^{i\xi_2}|^2$

$\Rightarrow$ relative phase is relevant

CP violation in Kaon decays (cont.)

with $\varepsilon = -\hat{\varepsilon} + \xi_0$

$$\varepsilon' = \left(\frac{A_2}{A_0 + A_2}\right) \boxed{(\xi_2 - \xi_0)}$$

we get : $\eta_{+-} = \varepsilon + \varepsilon'$, $\eta_{00} = \varepsilon - 2\varepsilon'$

Experiments (KTEV, NA48) have confirmed that

$\frac{\varepsilon'}{\varepsilon} \neq 0$, $\mathrm{Re}\left|\frac{\varepsilon'}{\varepsilon}\right| = 1.7 \times 10^{-3}$

$\longrightarrow \varepsilon'$ <u>small</u> relative to ε

Measurements with B mesons also have observed direct CP violation

$B-\bar{B}$ Mixing

Similar to $K-\bar{K}$ mixing but b quark mass $\gg$ s quark mass and well above QCD scale:

- Theoretical uncertainties are reduced.
- much larger phase space eliminates (essentially) the lifetime difference. Simplifies expressions ...
- B decays are more CKM suppressed

With $\Gamma_- \simeq \Gamma_+$, we'll denote the two states by "H" for heavy, and "L" for light. Previous oscillation probabilities become (non-relativistic B's):

$$P_T[B^0 \to B^0] = e^{-\Gamma T} \cos^2\left(\frac{\Delta m T}{2}\right)$$

$$P_T[B^0 \to \bar{B}^0] = \left|\frac{q}{p}\right|^2 e^{-\Gamma T} \sin^2\left(\frac{\Delta m t}{2}\right)$$

$B-\bar{B}$ Mixing (cont.)

In our example (non-rel, e.g. CLEO), it is difficult to measure time t elapsed since the B was in a pure B^0 eigenstate. We start from:

$$e^+e^- \to \gamma^* \to B\bar{B}$$

$\to$ relative angular momentum $\ell = 1$

$$CP|B\bar{B}\rangle = -|B\bar{B}\rangle$$

initial state $|B\bar{B}\rangle = \frac{1}{\sqrt{2}}[|B(k)\bar{B}(-k)\rangle - |B(-k)\bar{B}(k)\rangle$

We can then "Tag" the flavour of the B using semi-leptonic decay of one B (there are other ways to tag).

By reconstructing decay vertex, we can determine prob. of observing (for instance) same-sign leptons vs distance i.e. versus time

$\to$ asymmetric B factories make this easier

$B-\bar{B}$ Mixing (in more detail)

$$|\langle f|H|B^0\rangle|^2 = |\langle f|H|B^0(t)\rangle|^2$$

$$= \frac{1}{4|p|^2}|\langle f|B_L(t)\rangle + \langle f|B_H(t)\rangle|^2$$

$$= \frac{1}{4|p|^2}|pAe^{(-im_L-\Gamma_L/2)t} + pAe^{(-im_H-\Gamma_H/2)t}|^2$$

$$= \frac{1}{4}|A|^2\left(e^{-\Gamma_L t} + e^{-\Gamma_H t} + 2e^{-(\Gamma_H+\Gamma_L)t/2}\cos\Delta m t\right)$$

$$|\langle \bar{f}|H|B^0\rangle|^2 = |\langle \bar{f}|B^0(t)\rangle|^2$$

$$= \frac{1}{4|p|^2}|\langle \bar{f}|B_L(t)\rangle + \langle \bar{f}|B_H(t)\rangle|^2$$

$$= \frac{1}{4|p|^2}|q\bar{A}e^{(-im_L-\Gamma_L/2)t} - q\bar{A}e^{(-im_H-\Gamma_H/2)t}|^2$$

$$= \frac{1}{4}\left|\frac{p}{q}\right|^2|\bar{A}|^2\left(e^{-\Gamma_L t} + e^{-\Gamma_H t} - 2e^{-(\Gamma_H+\Gamma_L)t/2}\cos\Delta m T\right)$$

$B-\bar{B}$ Mixing (in more detail, cont.)

Remember we have:

$$|B_H\rangle = p|B^0\rangle + q|\bar{B}^0\rangle$$

$$|B_L\rangle = p|B^0\rangle - q|\bar{B}^0\rangle$$

$\rightarrow$ $\langle f|B^0\rangle = A$, $\langle \bar{f}|B^0\rangle = 0$

$\langle \bar{f}|\bar{B}^0\rangle = \bar{A}$, $\langle f|\bar{B}^0\rangle = 0$

$B_H \rightarrow f$ provides $p\,A$

$B_L \rightarrow \bar{f}$ provides $-q\,\bar{A}$

$\rightarrow$ if $A \neq \bar{A}$, CP violation in decay

$\rightarrow$ if $\left|\frac{q}{p}\right| \neq 1$, CP violation in mixing

Without CP violation we would have

$$|\langle f|H|B^0\rangle|^2 + |\langle \bar{f}|H|B^0\rangle|^2 = \frac{1}{2}|A|^2\left(e^{-\Gamma_L t} + e^{-\Gamma_H t}\right)$$

CKM Matrix

In SM, flavour changing and CP-violating physics is due to 4 parameters in unitary CKM matrix.

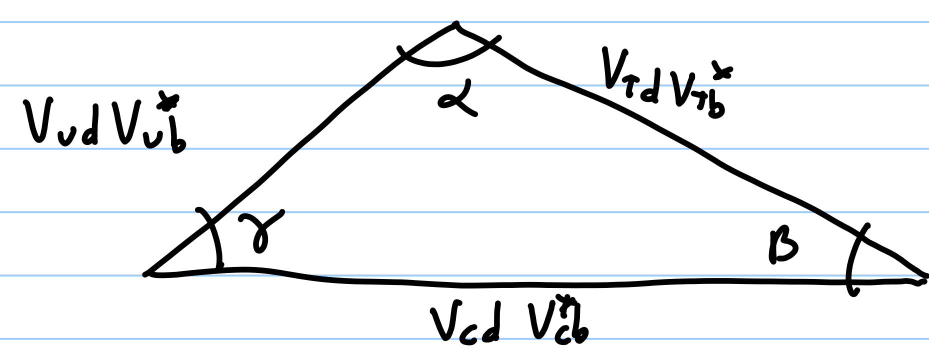
Physics beyond SM could provide new contributions. B physics provides many opportunities of testing SM.

Unitarity conditions $\rightarrow$ $\sum_i V_{in} V_{im}^* = \delta_{mn}$

With B_d mesons, we have $m = b$, $n = d$, which implies:

$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$

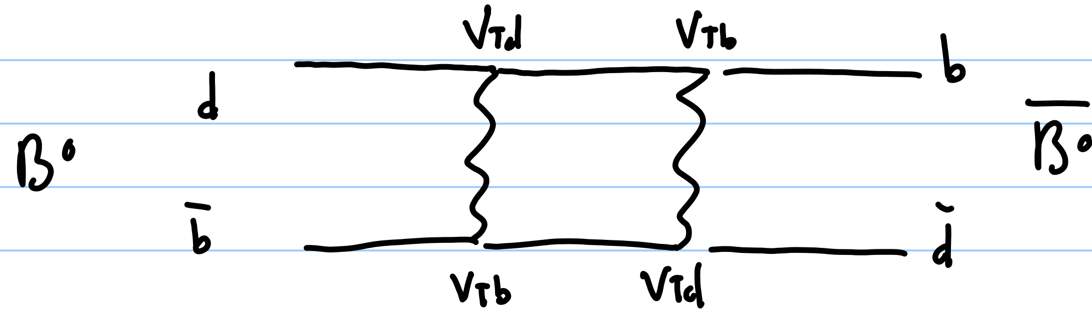
$\rightarrow$ sum of 3 complex numbers vanish.
$\rightarrow$ can be expressed 3 vectors in complex plane



$$\alpha = \arg\left(-\frac{V_{td} V_{tb}^*}{V_{ud} V_{ub}^*}\right)$$

$$\beta = \arg\left(\frac{V_{cd} V_{cb}^*}{V_{td} V_{tb}^*}\right)$$

CKM Matrix and B physics



$$\beta = \arg\left(\frac{V_{cd} V_{cb}^*}{V_{Td} V_{Tb}^*}\right)$$

$B^0 \to \psi K_s$: $V_{Td} V_{Tb}$ from mixing

V_{cb} from $b \to c$

V_{cd} from K^0 mixing

In asymmetric B factories:

- reconstruct one B in CP eigenstate e.g. ψK_s
- reconstruct decay of other B → determine flavour
- measure distance between mesons and convert to proper time

- Fit decay time difference to

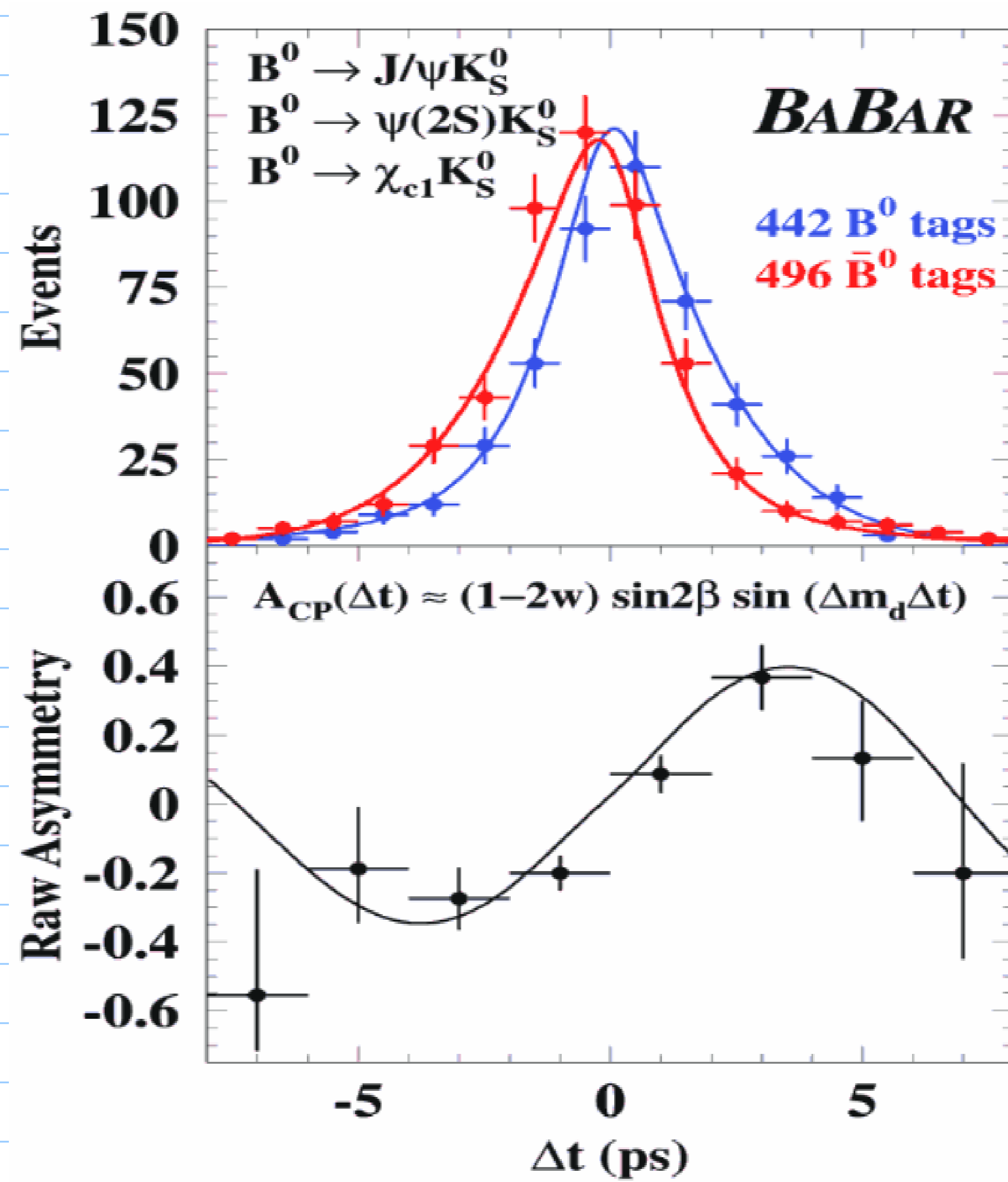
$$\frac{dW}{dt} \propto e^{-\Gamma|t|} \left[1 \pm F S_{CP} \left(\sin 2\varphi \right) \sin\left(\Delta m t \right) \right]$$

S_{CP} ↳ CP of final state

- Determine φ

Note that: $A = \dfrac{N(B_1 \to f_{CP}, \bar{B}_2 \to \bar{F}) - N(B_1 \to f_{CP}, B_2 \to F)}{\text{sum}} = 0$

But $A(t_1, t_2) \propto \sin 2\varphi_{CP}$



$B^0 \to J/\psi K^0_S$
$B^0 \to \psi(2S) K^0_S$
$B^0 \to \chi_{c1} K^0_S$
BABAR
442 B^0 tags
496 $\bar{B}^0$ tags
Events
150
125
100
75
50
25
0
$A_{CP}(\Delta t) \approx (1-2w) \sin 2\beta \sin(\Delta m_d \Delta t)$
Raw Asymmetry
0.6
0.4
0.2
0
-0.2
-0.4
-0.6
-5
0
5
Δt (ps)

Other sources of CP violation

CP violation due to CKM matrix is too small to explain observed matter anti-matter asymmetry in the Universe.

Are there other potential sources of CP violation in SM?

Yes:

- lepton sector (next lecture)
- Strong CP violation (very small if it exists)

Physics beyond SM could also contribute

Strong CP Problem

Most general renorm. Lagrangian involving SM fields includes:

$$\mathcal{L}_{F_S} = -\frac{1}{4} G^{\alpha}_{\mu\nu} G^{\alpha\mu\nu} - \frac{1}{4} W^{\alpha}_{\mu\nu} W^{\alpha\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu}$$

$$- \frac{g_3^2 \Theta_3}{64\pi^2} \varepsilon_{\mu\nu\lambda\rho} G^{\alpha\mu\nu} G^{\alpha\lambda\rho} \quad - \text{ similar Terms for } W, B$$

→ This Term is effectively a Total derivative and would be expected To have no physical implications

$$\propto \partial^{\mu} K_{\mu}$$

Term determined by change in charge $\int d^3x\, K^0 \equiv N_{CS}$

Change in charge need not be zero in a vacuum To vacuum process → QCD vacuum Topologically non-Trivial

Axions are a potential solution (need To add scalar field)