

PHENOMENOLOGY OF THE STANDARD MODEL PHY24085

I will cover:

- Introduction To gauge theories
- Spontaneous symmetry breaking
- Weak interactions
- Strong interactions
- CP Violation
- Neutrino physics

Suggested Texts:

Main - Thomson Modern Particle Physics

Others - Quigg
 - Burgess and Moore
 - Halzen Martin
 - Griffiths
 - Aitchison and Hey

Grading

- Weekly* problem sets	25%	**
- Long problem sets	50%	**
- Presentation (in April)	25%	

* almost weekly

** - 20% per week late

The Standard Model

A quantum field theory. Lagrangian obeys local gauge invariance

$$\begin{aligned} \psi &\rightarrow \psi' = e^{i\theta} \psi & \text{global} \\ \psi &\rightarrow \psi' = e^{i\alpha(x)} \psi & \text{local} \end{aligned}$$

Lagrangian invariant under a continuous group of local transformations (gauge transformations). They form a Lie Group whose generators have associated vector fields (gauge fields)

- 12 Fermions
 - 6 quarks
 - 6 leptons
- + 1 scalar : Higgs

- 3 Forces
 - EM
 - Weak
 - Strong
- spin 1 bosons
 - W^+, W^-, Z^0
 - 8 gluons

Elementary Particles

	Leptons			Quarks			Force Carriers		
I	ν_e electron neutrino	e electron	ν_e electron neutrino	u up	d down	ν_e electron neutrino	u up	d down	γ photon
II	ν_μ muon neutrino	μ muon	ν_μ muon neutrino	c charm	s strange	ν_μ muon neutrino	c charm	s strange	g gluon
III	ν_τ tau neutrino	τ tau	ν_τ tau neutrino	t top	b bottom	ν_τ tau neutrino	t top	b bottom	Z Z boson
									W W boson

Three Families of Matter

⊕

In more detail: SM based on the fields below
and $SU(3) \times SU(2)_L \times U(1)_Y$ gauge symmetries.

$$Q_L^i = \begin{pmatrix} u_L \\ d_L \end{pmatrix} \quad \begin{pmatrix} c_L \\ s_L \end{pmatrix} \quad \begin{pmatrix} T_L \\ b_L \end{pmatrix} \quad \begin{matrix} SU(3) \\ 3 \\ 3 \end{matrix} \quad \begin{matrix} SU(2)_L \\ 2 \end{matrix} \quad \begin{matrix} U(1)_Y \\ 1/6 \end{matrix}$$

$$u_R^i \quad u_R \quad c_R \quad t_R \quad \begin{matrix} 3 \\ 3 \\ 1 \\ 1 \end{matrix} \quad \begin{matrix} 2/3 \\ 2/3 \end{matrix}$$

$$d_R^i \quad d_R \quad s_R \quad b_R \quad \begin{matrix} 3 \\ 3 \\ 1 \\ 1 \end{matrix} \quad \begin{matrix} -1/3 \\ -1/3 \end{matrix}$$

$$L_L^i \quad \begin{pmatrix} \nu_{eL} \\ e_L \end{pmatrix} \quad \begin{pmatrix} \nu_{\mu L} \\ \mu_L \end{pmatrix} \quad \begin{pmatrix} \nu_{\tau L} \\ \tau_L \end{pmatrix} \quad \begin{matrix} 1 \\ 1 \\ 1 \\ 2 \end{matrix} \quad \begin{matrix} -1/2 \\ -1/2 \end{matrix}$$

$$e_R^i \quad e_R \quad \mu_R \quad \tau_R \quad \begin{matrix} 1 \\ 1 \\ 1 \\ 1 \end{matrix} \quad \begin{matrix} -1 \\ -1 \end{matrix}$$

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} \quad \begin{matrix} 1 \\ 1 \end{matrix} \quad \begin{matrix} 2 \\ 2 \end{matrix} \quad \begin{matrix} 1/2 \\ 1/2 \end{matrix}$$

Now we should add neutrinos above (and we will)

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We can break the SM into the following:

$$\mathcal{L}_{SM} = \mathcal{L}_{GAUGE} + \mathcal{L}_{MATTER} + \mathcal{L}_{YUKAWA} + \mathcal{L}_{HIGGS}$$

A brief look ahead and an answer to a question sometimes asked to graduate students.

- in the late 80s, early 90s, physicists started designing machines to probe physics at the TeV scale. The argument:

"we are sure to find something; either the Higgs or something else,"
why??

Let's look at \mathcal{L}_{HIGGS} :

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$$\mathcal{L}_{\text{Higgs}} = (D^\mu \varphi)^\dagger D_\mu \varphi + \underbrace{m^2 \varphi^\dagger \varphi - \lambda (\varphi^\dagger \varphi)^2}_{\text{potential}}$$

m^2 - is the only dimensionful parameter of the
S_A - sign chosen to provided nonzero VEV

Higgs acquisition of
non-zero VEV breaks
electroweak symmetry and gives
mass of W, Z bosons and
fermions

$$\langle \varphi_0 \rangle = \frac{M}{2\sqrt{2}\lambda} \equiv \frac{V}{\sqrt{2}}$$

$$V \approx 246 \text{ GeV}$$

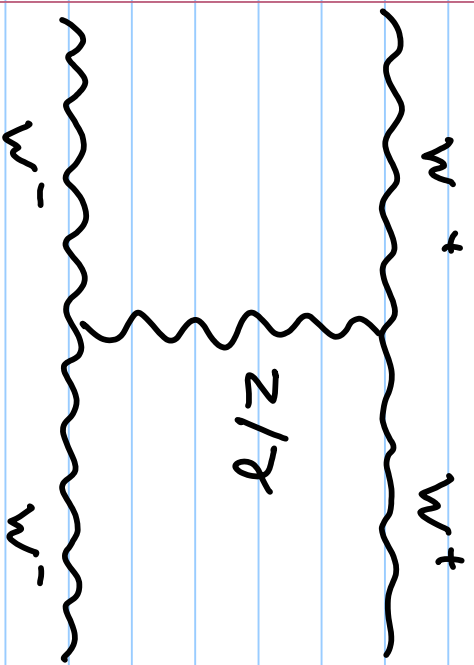
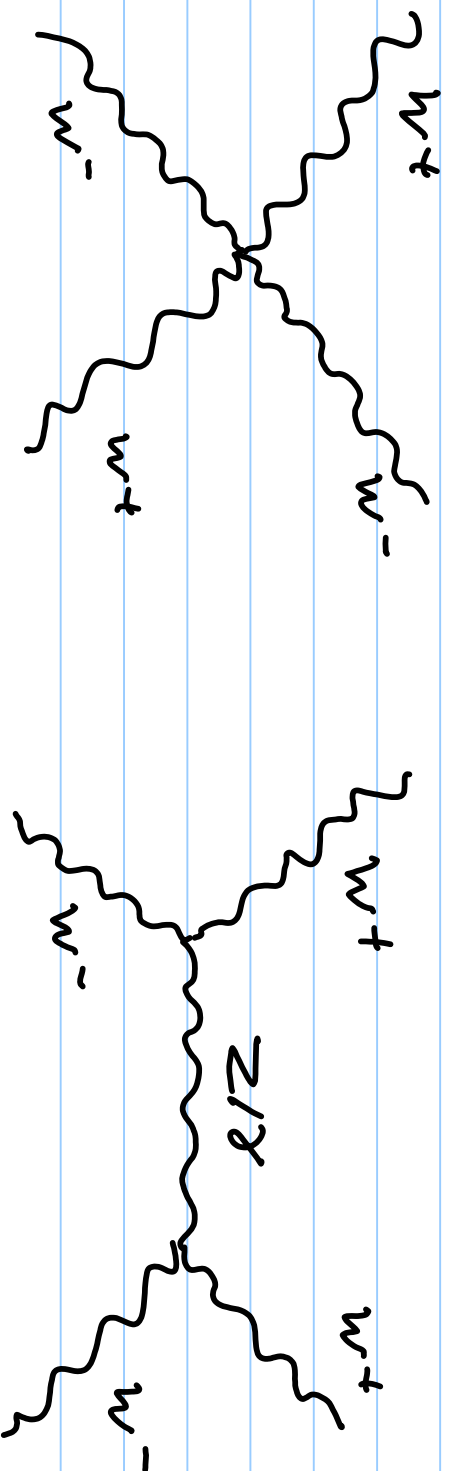
Now SU(2) x U(1) part of $\mathcal{L}_{\text{Gauge}}$:

$$-\frac{1}{4} F_{\mu\nu}^i F^{\mu\nu i} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu}$$

→ self interactions (non-Abelian theory)

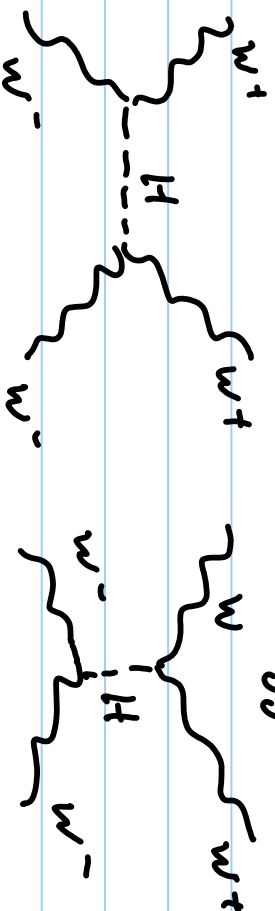
Have diagrams for WW scattering:

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$\rightarrow \sigma$ grows with s (E_{cm}^2)

and violates unitarity unless we add the Higgs



$\Rightarrow M_H \lesssim 1 \text{ TeV}$

New Physics

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→ problem with our Fundamental scalar:

Higgs mass receives quadratically-divergent corrections from loop diagrams.

$$\delta M_H^2 \propto \Lambda^2$$

in contrast $\delta m_F \propto \ln \left(\frac{\Lambda}{m_F} \right)$

⇒ "Natural" value for Higgs mass is at Λ , the cutoff energy

⇒ Keeping Higgs mass at weak scale while setting Λ at GUT scale ($\sim 10^{16}$ GeV) requires very precise Fine-Tuning.

To prevent this Fine Tuning, new physics should be close to weak scale,