

LECTURE 23: The QM Postulates (part deux)

What I expect you to learn:

- What are the QM postulates
- How to work with the Dirac notation

(Roughly reviews chapter 5 of the textbook)

RECAP OF POSTULATES USING THE INFINITE WELL EXAMPLE;

SUPPOSE WE PUT A PARTICLE (SAY A He^+ ion) IN AN INFINITE WELL.

A- WE USE A "STATE VECTOR" TO DESCRIBE THIS PARTICLE: $| \psi \rangle$, THIS VECTOR CONTAINS ALL THE INFORMATION ABOUT THIS PARTICLE:

- MOMENTUM
- POSITION
- ANGULAR MOM.
- ETC.

B- IF $| \psi \rangle$ DESCRIBES A He^+ ion IN AN INFINITE WELL AND $| \phi \rangle$ DESCRIBES ANOTHER He^+ ion IN AN INFINITE WELL, THEN:

$| \eta \rangle = C_1 | \psi \rangle + C_2 | \phi \rangle$ could also

DESCRIBE A He^+ ion IN ANOTHER INFINITE WELL

RECAP. CONT.

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C - TO EVERYTHING WE CAN MEASURE ABOUT THE H_{et} low WE ASSOCIATE AN OPERATOR THAT IS LINEAR AND HERMITIAN

AN OPERATOR λ TRANSFORMS A STATE VECTOR INTO A NEW STATE VECTOR: $O|\psi\rangle = \lambda|\psi\rangle$

IF $|\psi\rangle$ IS AN EIGENVECTOR OF O , THEN $O|\psi\rangle = c|\psi\rangle$
complex number \leftarrow

Initial state



Operator



Final state



* IF WE MEASURE SOME PHYSICAL QUANTITY A , THE RESULT CAN ONLY BE AN EIGENVALUE OF \hat{A} , THE OPERATOR ASSOCIATED WITH THE QUANTITY A .

→ AND ← THE SYSTEM WILL BE IN AN EIGENSTATE OF \hat{A} RIGHT AFTER THE MEASUREMENT.

quantum
jackalope



rabbit
eigenstate



⇒ MEASUREMENT
↓
Operator

antelope
eigenstate



RECAP. CONT.

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E - ONE CAN EXPAND A STATE $| \psi \rangle$ IN TERMS OF THE EIGENSTATES ASSOCIATED WITH THE OPERATOR THAT REPRESENTS A PHYSICALLY MEASURABLE QUANTITY (ENERGY, MOM., POSITION, ANGULAR MOM., ETC).

$$| \psi \rangle = \sum_n A_n | E_n \rangle \quad (\text{discrete basis})$$

\hookrightarrow energy eigenstate

$$| \psi \rangle = \int dx c(x) | x \rangle \quad (\text{continuous basis})$$

\hookrightarrow position eigenstate

$\hookrightarrow \psi(x)!$

TO PERFORM THIS EXPANSION, ONE CAN USE A PROJECTION OPERATOR:

$$\sum_n | E_n \rangle \langle E_n | = 1$$

$$\Rightarrow | \psi \rangle = \sum_n | E_n \rangle \langle E_n | \psi \rangle = \sum_n | E_n \rangle \cdot A_n$$

$$| \psi \rangle = \int dx | x \rangle \langle x | \psi \rangle = \int dx | x \rangle \psi(x)$$

RECAP COMT.

⑥

- TO CALCULATE THE PROBABILITY OF FINDING THE PARTICLE WITH AN ENERGY E_M , ONE DOES:

$$|\langle E_M | \psi \rangle|^2 = |\text{complex number}|^2$$

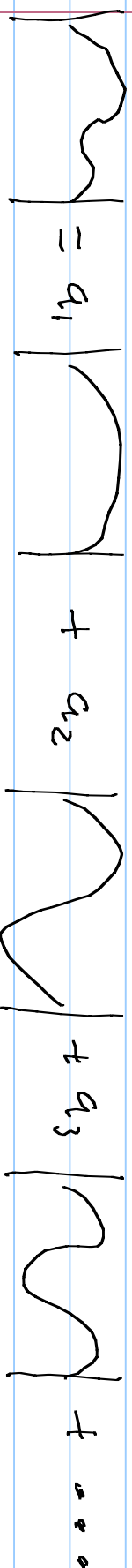
IF WE INSERT THE POSITION EIGENSTATE PROJECTOR:

$$= \left| \int_{-\infty}^{\infty} dx \langle E_M | x \rangle \langle x | \psi \rangle \right|^2 = \left| \int_{-\infty}^{\infty} dx \psi(x) \psi(x) \right|^2 \quad \text{①}$$

$$\psi(x) = \sum_m a_m \psi_m(x) \quad \text{and} \quad \int_{-\infty}^{\infty} \psi_m(x) \psi_n(x) dx = \delta_{mn}$$

$$\text{①} = \left| \int_{-\infty}^{\infty} dx \psi(x) \sum_m a_m \psi_m(x) \right|^2 = |a_M|^2$$

REMEMBER THE FOURIER SERIES:



A diagram illustrating the Fourier series decomposition of a periodic function. On the left, a periodic wave is shown with the equation $\psi(x) = a_1 \psi_1(x) + a_2 \psi_2(x) + a_3 \psi_3(x) + \dots$. The wave is composed of several individual sine wave components, each labeled with a coefficient a_n and a basis function $\psi_n(x)$.

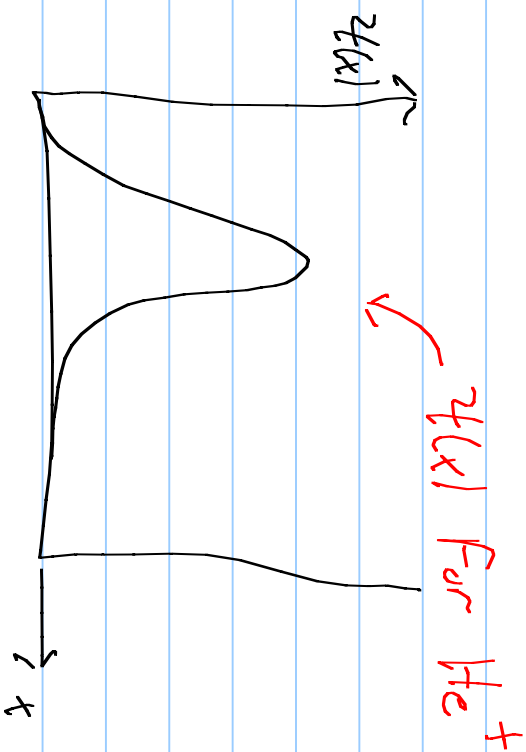
RECAP CONT.

⑦

BACK TO OUR H_e^+ ion in the infinite well described by the state vector $|\psi\rangle$.

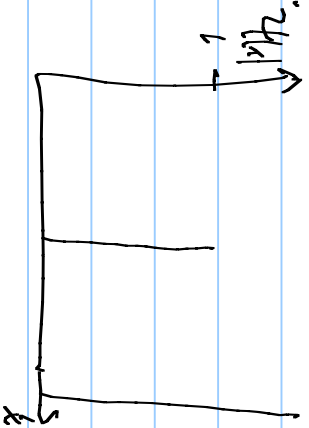
THE PROBABILITY OF MEASURING THE POSITION OF THE PARTICLE BETWEEN x_1 AND x_2 IS:

$$\int_{x_1}^{x_2} |\langle x|\psi\rangle|^2 dx$$
$$= \int_{x_1}^{x_2} |\psi(x)|^2 dx$$

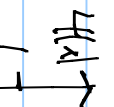


AFTER THE MEASUREMENT, $\psi(x)$ COULD LOOK LIKE

THIS:



or



or

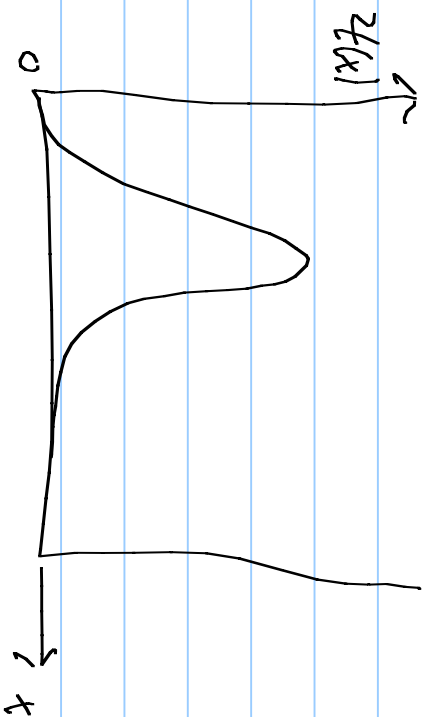


RECAP CONT.

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BEFORE THE MEASUREMENT
WE HAD FOR $\psi(x)$:

THIS FUNCTION CAN BE
EXPANDED USING A FOURIER
SERIES



$$\psi(x) \approx \sum_n A_n \psi_n(x) = \sum_n A_n \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a}$$

NOW WHEN WE MEASURE ENERGY, WE CAN ONLY GET
ONE OF THE FOLLOWING EIGENVALUES: $E_n = \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2 \pi^2 n^2}{2m a^2}$

THE PROBABILITY OF MEASURING A GIVEN EIGENVALUE IS:

$$P_n = |\langle E_n | \psi \rangle|^2 = \left| \int_{-\infty}^{\infty} dx \langle E_n | x \rangle \psi(x) \right|^2$$
$$= \left| \int_{-\infty}^{\infty} dx \psi_n(x) \psi(x) \right|^2$$

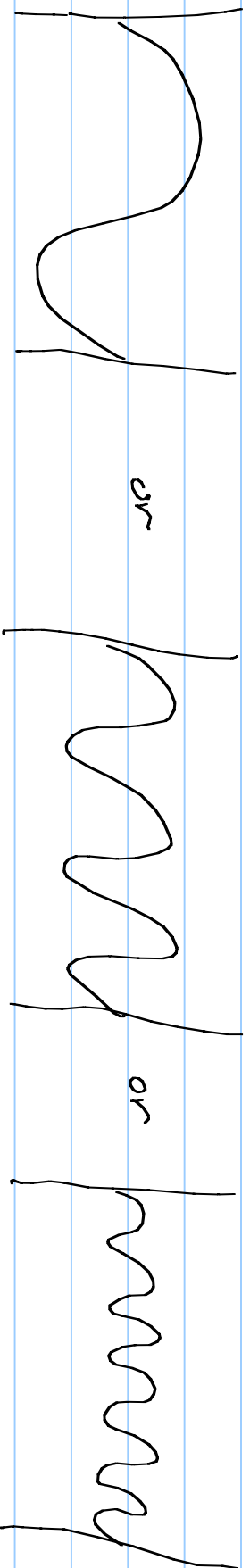
RECAP CONT.

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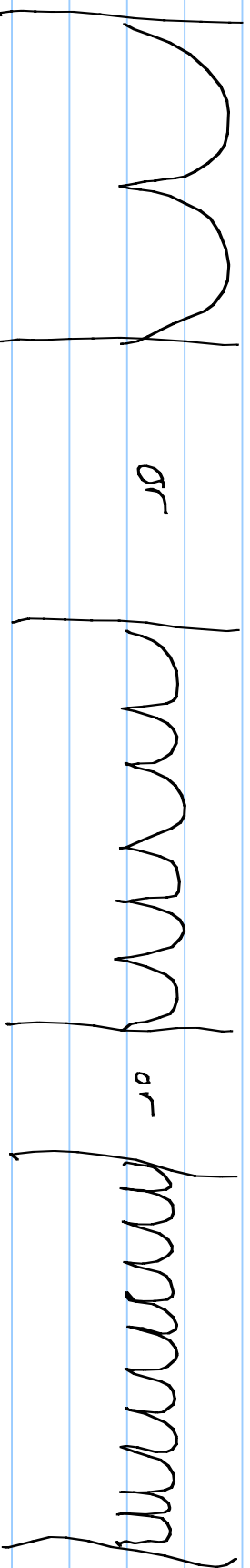
NOTE THAT $E_M = \frac{\hbar^2 k^2}{2m} = \frac{p^2}{2m}$ ($V=0$)

\Rightarrow IF YOU KNOW p , YOU KNOW E . THE OPERATORS ASSOCIATED WITH p AND E COMMUTE

AFTER AN ENERGY MEASUREMENT, $\psi(x)$ MUST BE ONE OF THE $\psi_n(x)$: eg



$$|\psi_n(x)|^2 =$$



RECAP (cont)

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WE STARTED WITH $\psi(x)$ \rightarrow

IF WE MEASURE E (or P)

FIRST WE WILL GET

ONE OF THE $\psi_n(x)$,

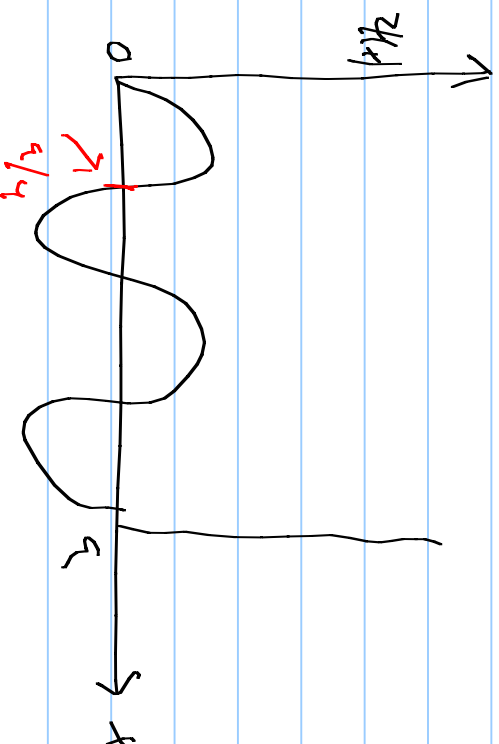
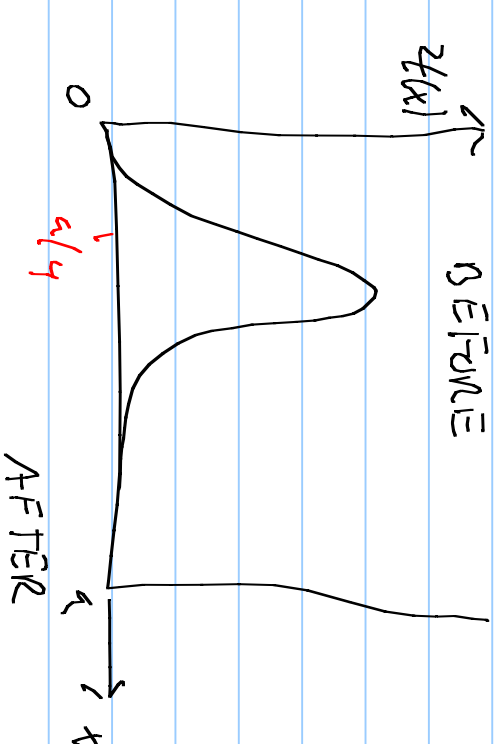
SAY: $A_n \cdot \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a}$

IF I MEASURE THE

POSITION NOW, THE

PROBABILITY OF FINDING THE

PARTICLE AT $x = \frac{a}{4}$ WOULD BE: 0

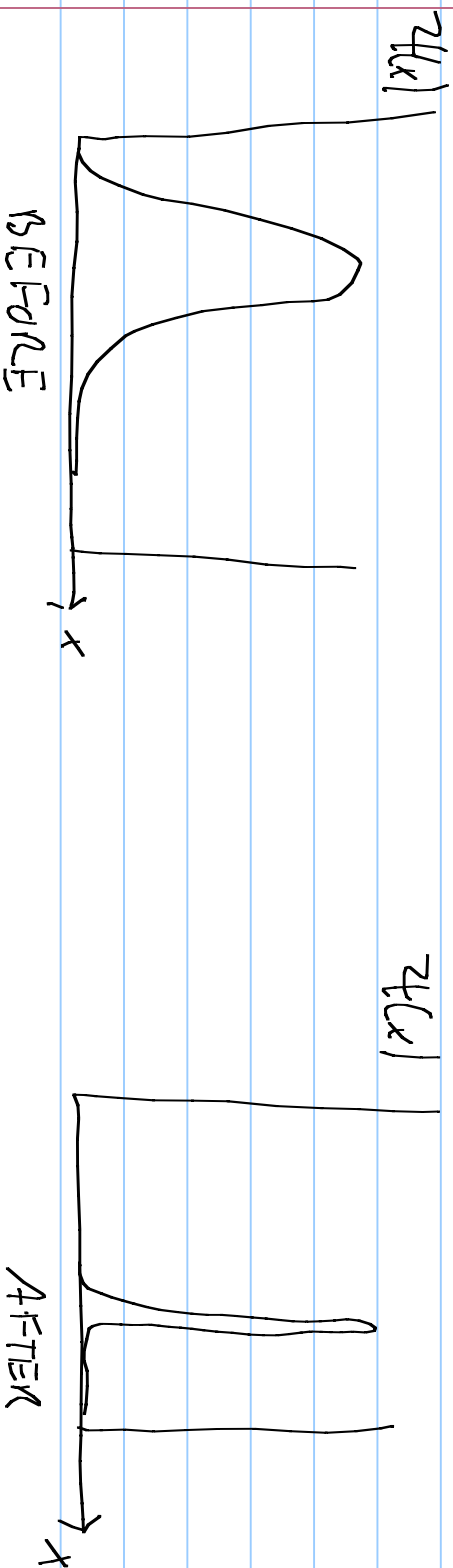


IF YOU MEASURE THE POSITION FIRST, $P(a/4) \neq 0$.

RECAP (cont.)

III

ALSO $P(E_n) = |A_n|^2$ WHEN ENERGY WAS MEASURED
FIRST BUT $P(E_n)$ AFTER A POSITION MEASUREMENT
THE A_n COEFFICIENT WILL HAVE CHANGED:



FOURIER SERIES OF "BEFORE" \neq FOURIER SERIES OF "AFTER"

$$\text{BEFORE } \sum_n B_n \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a} \neq \text{AFTER } \sum_n A_n \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a}$$

$\Rightarrow B_n \neq A_n \Rightarrow |B_n|^2 \neq |A_n|^2 \Rightarrow$ PROBABILITY
OF MEASURING THE ENERGY E_n BEFORE MEASURING THE
POSITION \neq PROB. OF MEASURING E_n AFTER POSITION MEASUREMENT.

