

LECTURE 27: The Last Postulate

What I expect you to learn:

- How quantum systems evolve with time
- What is the Evolution Operator
- What are the Heisenberg and Schrodinger "pictures"

(Corresponds to chapter 5.7 of the textbook)

(Problem set 4: due Wednesday Nov. 22nd)

GENERAL SOLUTION FOR A TIME INDEPEND. POTENTIAL (2)

WE SAW THAT WE COULD EXPRESS $\psi(x,t) = \psi(x)F(t)$

$$\text{WHERE } F(t) = C \exp(-iEt/\hbar) \quad \text{or } C(T_0) \exp\left[-\frac{iE}{\hbar}(t-T_0)\right]$$

$$\text{WE HAVE } \psi(x) = \sum_n A_n \psi_n(x)$$

$$\Rightarrow \psi(x,t) = \sum_n A_n \psi_n(x) C(T_0) \exp\left[-\frac{iE_n}{\hbar}(t-T_0)\right]$$

we could rewrite $\psi(x,t) = \sum_n A_n \psi_n(x) \exp(-iEt/\hbar)$

$$\text{with } A_n = C_n(T_0) \exp(iE_n T_0/\hbar)$$

IF WE KNOW THE WAVE FUNCTION ψ AT A PARTICULAR TIME T_0 , we can determine ψ FOR ALL VALUES OF T :

$$A_n = \exp(iE_n T_0/\hbar) \int \psi_n^*(x) \psi(x, T_0) dx$$

$$\psi(x,t) = \sum_n \left[\int \psi_n^*(x) \psi(x, T_0) dx' \right] \psi_n(x) \exp\left[-\frac{iE_n}{\hbar}(t-T_0)\right]$$

GENERAL SOLUTION FOR A TIME INDEP. POTENTIAL (3)

So we have: $\psi(x,t) = \sum_n A_n \psi_n(x) \exp(-iEt/\hbar)$

$$\begin{aligned} P(x,t) &= \psi^*(x,t) \psi(x,t) \\ &= \sum_n \sum_m A_m^* A_n \exp[-i(E_n - E_m)t/\hbar] \psi_m^*(x) \psi_n(x) \end{aligned}$$

$$= \sum_n |A_n|^2 |\psi_n|^2 + \sum_{n \neq m} A_m^* A_n \exp\left[-\frac{i}{\hbar}(E_n - E_m)t\right] \psi_m^*(x) \psi_n(x)$$

→ Pure eigenstates have no time dep.

→ The larger the ΔE , the faster the probability evolves versus time

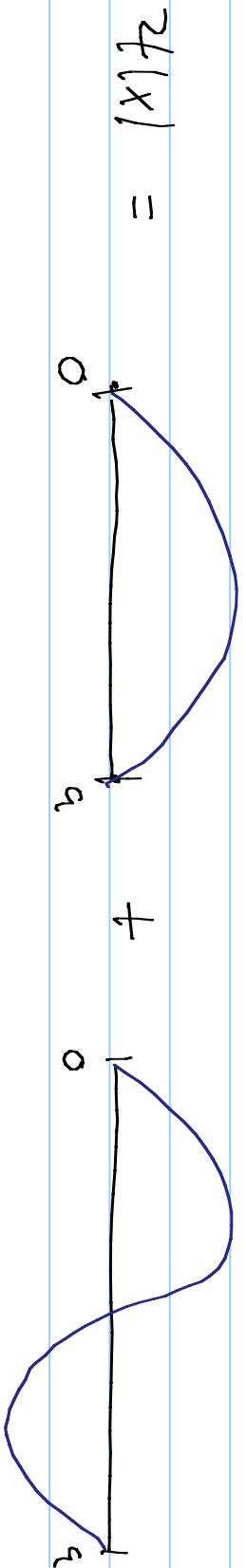
→ $P(x,t)$ is time dependent, physical quantities will in general depend on time. Show however that $\langle H \rangle$ does not depend on time (exercise).

BACK TO THE INFINITE WELL

(4)

CONSIDER THE FOLLOWING WAVE FUNCTION (at $T=0$)

$$\psi(x) = A \sin\left(\frac{\pi x}{a}\right) + A \sin\left(\frac{2\pi x}{a}\right)$$



$n=1$

$n=2$

$$\psi(x, t) = \underbrace{A \sin\left(\frac{\pi x}{a}\right)}_{F_1} e^{-i\frac{E_1}{\hbar}t} + A \sin\left(\frac{2\pi x}{a}\right) \underbrace{e^{-i\frac{E_2}{\hbar}t}}_{F_2}$$

$$\psi(x, t) = F_1 e^{-iE_1 t/\hbar} + F_2 e^{-iE_2 t/\hbar}$$

$$P(x, t) = (F_1 e^{iE_1 t/\hbar} + F_2 e^{iE_2 t/\hbar})(F_1 e^{-iE_1 t/\hbar} + F_2 e^{-iE_2 t/\hbar})$$

$$= F_1^2 + F_2^2 + F_1 F_2 \left[e^{i(E_1 - E_2)t/\hbar} + e^{-i(E_1 - E_2)t/\hbar} \right]$$

$$= F_1^2 + F_2^2 + 2F_1 F_2 \cos\left(\frac{t}{\hbar}(E_1 - E_2)\right)$$

THE INFINITE WELL (cont.)

(5)

$$P(x,t) = F_1^2 + F_2^2 + 2F_1F_2 \cos\left(\frac{T}{\hbar}(E_1 - E_2)\right)$$

oscillates between $(F_1 + F_2)^2$ and $(F_1 - F_2)^2$

we have $E_1 = \frac{\hbar^2 \pi^2}{2ma^2}$, $E_2 = \frac{4\hbar^2 \pi^2}{2ma^2}$, $E_1 - E_2 = -3E_1$

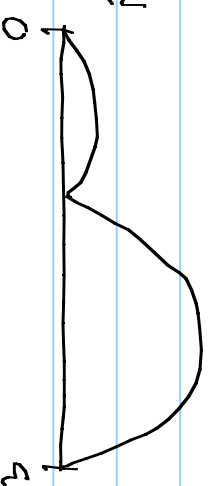
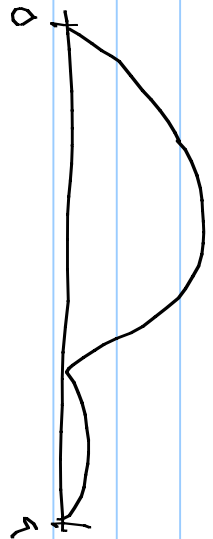
$$\Rightarrow \frac{3E_1}{\hbar} T = \pi, 3\pi, 5\pi, \dots$$

\rightarrow obtain same config. when $\frac{3E_1}{\hbar} T = 2\pi, 4\pi, \dots$

$$\psi(x) = \begin{array}{c} \text{graph of } \psi(x) \text{ from } 0 \text{ to } a \\ + \\ \text{graph of } \psi(x) \text{ from } 0 \text{ to } a \end{array}$$

$$P(x,t) \quad T = \frac{\hbar \pi}{3E_1}$$

$$P(x,0) =$$



THE LAST POSTULATE

(5)

A SUMMARY OF THE POSTULATES WE'VE SEEN

- 1- INTRODUCED THE STATE VECTOR $|\psi\rangle$
- * 2- THE SUPERPOSITION PRINCIPLE: $|\psi\rangle = c_1|\psi_1\rangle + c_2|\psi_2\rangle$
- 3- ASSOCIATES A LINEAR HERMITIAN OPERATOR TO EVERY OBSERVABLE

- * 4- A MEASUREMENT OF AN OBSERVABLE YIELDS AN EIGENVALUE OF THE OPERATOR AND LEAVES THE SYSTEM IN AN EIGENSTATE OF THE OPERATOR
- 5- $\langle A \rangle = \langle \psi | A | \psi \rangle$
- 6- $|\psi\rangle = \sum_n a_n |\psi_n\rangle$, $|\psi\rangle = \int a(x) |\psi_x\rangle dx$

- 7- THE TIME EVOLUTION OF $|\psi\rangle$ IS DETERMINED BY THE TIME-DEPENDENT SCHRÖDINGER EQUATION:

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = \hat{H} |\psi(t)\rangle$$

FIRST-ORDER DIFF. LINEAR EQUATION IN TIME.
IF YOU KNOW $|\psi\rangle$ AT $T=0$, YOU KNOW $|\psi\rangle$ FOR ALL T

THE EVOLUTION OPERATOR

(7)

WE INTRODUCE AN EVOLUTION OPERATOR $U(t, t_0)$:

$$|\psi(t)\rangle = U(t, t_0) |\psi(t_0)\rangle$$

↳ BANNED IN KANSAS ...

$$\hat{U}(T_0, T_0) = I$$

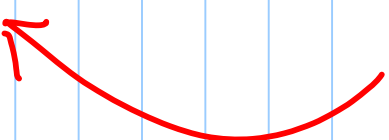
(1)

$$\hat{U}(T, T_0) = \hat{U}(T, T_1) \hat{U}(T_1, T_0)$$

(2)

$$\hat{U}^{-1}(T, T_0) = \hat{U}(T_0, T)$$

(3)



\hat{U}



$\rangle = |$



\rangle


THE EVOLUTION OPERATOR (cont.) (8)

$$\rightarrow i\hbar \frac{d}{dt} (U(t, T_0) | \psi(t_0) \rangle) = \hat{H} (U(t, T_0) | \psi(t_0) \rangle) \quad (9)$$

$$\rightarrow i\hbar \frac{d}{dt} U(t, T_0) = \hat{H} U(t, T_0) \quad (5)$$

using (1) AND (5)

$$\rightarrow \ddot{U}(t, T_0) = I - \frac{i}{\hbar} \int_{T_0}^t \hat{H} \ddot{U}(t', T_0) dt' \quad (6)$$


$$\frac{d}{dt} \ddot{U}(t, T_0) = \frac{i}{\hbar} \hat{H} \ddot{U}(t, T_0) = -\frac{i}{\hbar} U(t, T_0)$$

$$\Rightarrow \int_{t_0}^T \frac{d}{dt} \ddot{U}(t, T_0) dt = \ddot{U}(t, T_0) - \ddot{U}(T_0, T_0) = \ddot{U}(t, T_0) - I$$

EVOLUTION OPERATOR (CONT.)

(9)

CONSERVATION OF PROB. IMPLIES:

$$\langle \psi(t_0) | \psi(t_0) \rangle = \langle \psi(t) | \psi(t) \rangle$$

$$\langle \psi(t) | \psi(t) \rangle = \langle \hat{U}(t, t_0) \psi(t_0) | \hat{U}(t, t_0) \psi(t_0) \rangle$$

$$= \langle \psi(t_0) | \hat{U}^\dagger(t, t_0) \hat{U}(t, t_0) \psi(t_0) \rangle$$

$$\Rightarrow \hat{U}^\dagger \hat{U} = I \quad \text{AND} \quad \hat{U} \hat{U}^\dagger = I$$

HOW DOES \hat{U} CHANGE WHEN $t \rightarrow t + \delta t$?

$$\hat{U}(t_0 + \delta t, t_0) = \hat{U}(t_0, t_0) = I \quad \leftarrow \text{FROM (6)}$$

$$\hat{U}(t_0 + \delta t, t_0) = I - \frac{i}{\hbar} \hat{H} \delta t$$

\rightarrow \hat{H} generates an infinitesimal Time Translation

Evolution OPERATOR CDMT.

(10)

A solution to: $i\hbar \frac{d}{dt} U(t, T_0) = \hat{H} U(t, T_0)$

$$\begin{aligned} \text{is: } U(t, T_0) &= \exp \left[-i\hbar \hat{H} (t - T_0) \right] \\ &= \sum_{n=0}^{\infty} \frac{1}{n!} \left(\frac{-i}{\hbar} \right)^n \hat{H}^n (t - T_0)^n \end{aligned}$$

FIRST TWO TERMS:

$$I - \frac{i}{\hbar} \hat{H} (t - T_0)$$

$$\Rightarrow |\psi(t)\rangle = \exp \left[\frac{-i}{\hbar} \hat{H} (t - T_0) \right] |\psi(T_0)\rangle$$

$\hat{H} \rightarrow$ Time indep. potential

AN EXAMPLE:

$$|2\rangle\langle 0| = \frac{1}{5} \begin{pmatrix} 3 \\ 0 \\ 4 \end{pmatrix}, \quad \hat{H} = \epsilon \begin{pmatrix} 3 & 0 & 0 \\ 0 & 0 & 5 \\ 0 & 5 & 0 \end{pmatrix}$$

- 1- FIND WHAT VALUES CAN BE MEASURED FOR THE ENERGY AND WITH WHAT PROBABILITIES.
- 2- FIND $\langle 2|\psi\rangle$ AND EXPAND IN TERMS OF THE EIGENSTATES OF \hat{H} .
- 3- FIND THE EXPECTATION VALUE OF \hat{H} AT TIMES 0 AND T .

(SOLUTION ON BLACKBOARD, PARTS OF THE SOLUTIONS FOLLOW)

(12)

FIRST STEP, DIAGONALISE THE MATRIX AND FIND THE EIGENVALUES:

$$E_1 = -5\varepsilon, \quad E_2 = 3\varepsilon, \quad E_3 = 5\varepsilon$$

STEP TWO, FIND THE EIGENVECTORS:

$$|q_1\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ -1 \\ 1 \end{pmatrix}, \quad |q_2\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \quad |q_3\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix}$$

DETERMINE THE PROBABILITIES FOR OBTAINING THE E_n :

$$P(E_1) = |\langle q_1 | \psi(0) \rangle|^2 = \left| \frac{1}{5\sqrt{2}} (0, -1, 1) \begin{pmatrix} 3 \\ 0 \\ 4 \end{pmatrix} \right|^2 = \frac{8}{25}$$

$$P(E_2) = 9/25, \quad P(E_3) = 8/25$$

EXPAND $|\psi_0\rangle$ IN TERMS OF EIGENVECTORS:

$$|\psi_0\rangle = \frac{1}{5} \begin{pmatrix} 3 \\ 0 \\ 4 \end{pmatrix} = \frac{\sqrt{8}}{5} |q_1\rangle + \frac{3}{5} |q_2\rangle + \frac{\sqrt{8}}{5} |q_3\rangle$$

(13)

APPLY TIME EVOLUTION OPERATOR

$$|2(t)\rangle = \exp\left[-\frac{i}{\hbar} \hat{H}(t)\right] |2(0)\rangle$$

EXPAND $|2(t)\rangle$ IN TERMS OF EIGENVECTORS:

$$|2(t)\rangle = \frac{\sqrt{8}}{5} e^{-iE_1 T/\hbar} |e_1\rangle + \frac{3}{5} e^{-iE_2 T/\hbar} |e_2\rangle + \frac{\sqrt{8}}{5} e^{-iE_3 T/\hbar} |e_3\rangle$$

$$= \frac{\sqrt{8}}{5} \begin{pmatrix} 0 \\ -1/\sqrt{2} e^{+i5Et/\hbar} \\ 1/\sqrt{2} \cdot e^{+i5Et/\hbar} \end{pmatrix} + \frac{3}{5} \begin{pmatrix} e^{-i3Et/\hbar} \\ 0 \\ 0 \end{pmatrix} + \frac{\sqrt{8}}{5} \begin{pmatrix} 0 \\ 1/\sqrt{2} e^{-i5Et/\hbar} \\ 1/\sqrt{2} e^{-i5Et/\hbar} \end{pmatrix}$$

$$= \begin{pmatrix} 3/5 e^{-i3Et/\hbar} \\ -2/5 e^{i5Et/\hbar} + 2/5 e^{-i5Et/\hbar} \\ 2/5 e^{i5Et/\hbar} + 2/5 e^{-i5Et/\hbar} \end{pmatrix} = \begin{pmatrix} 3/5 e^{-i3Et/\hbar} \\ 4/5 i \sin(5Et/\hbar) \\ 4/5 \cos(5Et/\hbar) \end{pmatrix}$$

(19)

Is $|z(t)| >$ normalised? i.e. $\langle z(t) | z(t) \rangle = 1$?

$$\left(\frac{3}{5} e^{i3\epsilon t/\hbar}, \frac{4}{5i} \sin(5\epsilon t/\hbar), \frac{4}{5} \cos(5\epsilon t/\hbar) \right)$$

$$\left(\frac{3}{5} e^{-i3\epsilon t/\hbar}, \frac{-4}{5i} \sin(5\epsilon t/\hbar), \frac{4}{5} \cos(5\epsilon t/\hbar) \right)$$

$$= \frac{9}{25} e^{i3\epsilon t/\hbar} \cdot e^{-i3\epsilon t/\hbar}$$

$$+ \frac{16}{25} \cdot \sin^2(5\epsilon t/\hbar) + \frac{16}{25} \cos^2(5\epsilon t/\hbar)$$

$$= \frac{9}{25} + \frac{16}{25} \left(\sin^2(5\epsilon t/\hbar) + \cos^2(5\epsilon t/\hbar) \right)$$

$$= \frac{9}{25} + \frac{16}{25} = \frac{25}{25}$$

Question 3

$$\langle \psi_0 | \hat{H} | \psi_0 \rangle = (3/5, 0, 4/5) \begin{pmatrix} 3\epsilon & 0 & 0 \\ 0 & 0 & 5\epsilon \\ 0 & 5\epsilon & 0 \end{pmatrix} \begin{pmatrix} 3/5 \\ 0 \\ 4/5 \end{pmatrix} \quad (15)$$

$$\text{or } \boxed{= 27/25 \epsilon}$$

$$\text{or } \langle \psi \rangle = \sum_n P_n E_n = \frac{8}{25} \cdot -5\epsilon + \frac{9}{25} \cdot 3\epsilon + \frac{8}{25} \cdot 5\epsilon = \boxed{\frac{27\epsilon}{25}}$$

$$\text{or } \frac{8}{25} \langle \psi_1 | \hat{H} | \psi_1 \rangle + \frac{9}{25} \langle \psi_2 | \hat{H} | \psi_2 \rangle + \frac{8}{25} \langle \psi_3 | \hat{H} | \psi_3 \rangle$$

$$= \frac{8}{25} \cdot -5\epsilon + \frac{9}{25} \cdot 3\epsilon + \frac{8}{25} \cdot 5\epsilon = \boxed{27/25 \epsilon}$$

$$\langle \psi | \hat{H} | \psi \rangle =$$

$$\left(\frac{3}{5} e^{i3\epsilon t/\hbar}, \frac{4}{5} \sin(5\epsilon t/\hbar), \frac{4}{5} \cos(5\epsilon t/\hbar) \right)$$

$$\begin{pmatrix} 3\epsilon & 0 & 0 \\ 0 & 0 & 5\epsilon \\ 0 & 5\epsilon & 0 \end{pmatrix}$$

$$\begin{pmatrix} \frac{3}{5} e^{-i3\epsilon t/\hbar} \\ -\frac{4}{5} \sin(5\epsilon t/\hbar) \\ \frac{4}{5} \cos(5\epsilon t/\hbar) \end{pmatrix}$$

$$\boxed{= 27/25 \epsilon}$$

TIME VARIATION OF EXPECTATION VALUES

(16)

$$\frac{d}{dt} \langle \hat{A} \rangle = \frac{d}{dt} \langle \psi | \hat{A} | \psi \rangle$$

$$= \left\langle \frac{d\psi}{dt} | \hat{A} | \psi \right\rangle + \left\langle \psi | \frac{d\hat{A}}{dt} | \psi \right\rangle + \left\langle \psi | \hat{A} | \frac{d\psi}{dt} \right\rangle$$

$$i\hbar \frac{d}{dt} = \hat{H} \rightarrow \frac{d}{dt} = \frac{\hat{H}}{i\hbar}$$

$$= -\frac{1}{i\hbar} \langle \hat{H} \psi | \hat{A} | \psi \rangle + \left\langle \psi | \frac{d\hat{A}}{dt} | \psi \right\rangle + \frac{1}{i\hbar} \langle \psi | \hat{A} | \hat{H} \psi \rangle$$

$$= -\frac{1}{i\hbar} \langle \psi | \hat{H} \hat{A} | \psi \rangle + \frac{1}{i\hbar} \langle \psi | \hat{A} | \hat{H} \psi \rangle + \left\langle \psi | \frac{d\hat{A}}{dt} | \psi \right\rangle$$

$$= \frac{1}{i\hbar} \langle \psi | \hat{A} | \hat{H} \psi \rangle - \langle \psi | \hat{H} \hat{A} | \psi \rangle + \left\langle \psi | \frac{d\hat{A}}{dt} | \psi \right\rangle$$

$$= \frac{1}{i\hbar} \langle [\hat{A}, \hat{H}] \rangle + \left\langle \frac{d\hat{A}}{dt} \right\rangle$$

TIME VARIATION OF EXPECTATION VALUES (V2)

(17)

$$\begin{aligned}
 \langle \hat{A} \rangle (t) &= \langle \psi(t) | \hat{A} | \psi(t) \rangle \\
 &= \langle e^{-i\hat{H}t/\hbar} \psi(0) | \hat{A} | e^{i\hat{H}t/\hbar} \psi(0) \rangle \\
 &= \langle \psi(0) | e^{i\hat{H}t/\hbar} \hat{A} e^{-i\hat{H}t/\hbar} | \psi(0) \rangle \\
 &= \langle \psi(0) | \hat{A}(t) | \psi(0) \rangle
 \end{aligned}$$

With no explicit time dependence of \hat{A} , we have

$$\begin{aligned}
 \rightarrow \frac{d}{dt} \langle \hat{A}(t) \rangle &= \frac{i}{\hbar} \langle \psi(0) | e^{i\hat{H}t/\hbar} \hat{A} e^{-i\hat{H}t/\hbar} - \frac{i}{\hbar} e^{i\hat{H}t/\hbar} \hat{A} e^{-i\hat{H}t/\hbar} \hat{H} \\
 &= \frac{i}{\hbar} \langle \psi(0) | [\hat{H}, \hat{A}(t)] \rangle
 \end{aligned}$$

SCHRÖDINGER AND HEISENBERG PICTURES

(18)

SCHRÖDINGER PICTURE :

TIME-DEPEND WAVE FUNCTION THAT SATISFIES SCHRÖDINGER'S EQUATION

GET $\psi(t)$ FROM $\psi(0)$ USING : $\psi(t) = U(T, T_0) \psi(T_0)$

HEISENBERG PICTURE :

$$\begin{aligned} \psi_H &\equiv \psi_0 \\ &= U^\dagger(T, T_0) \psi(t) = U(T_0, t) \psi(t) = \psi(0) \end{aligned}$$

WE PUT THE TIME EVOLUTION IN THE OPERATOR :

$$\begin{aligned} A_H(t) &= U^\dagger(T, T_0) \hat{A} U(T, T_0) \\ &= U(T_0, t) \hat{A} U(T_0, t) \end{aligned}$$

