

LECTURE 31: Spin 1/2

What I expect you to learn:

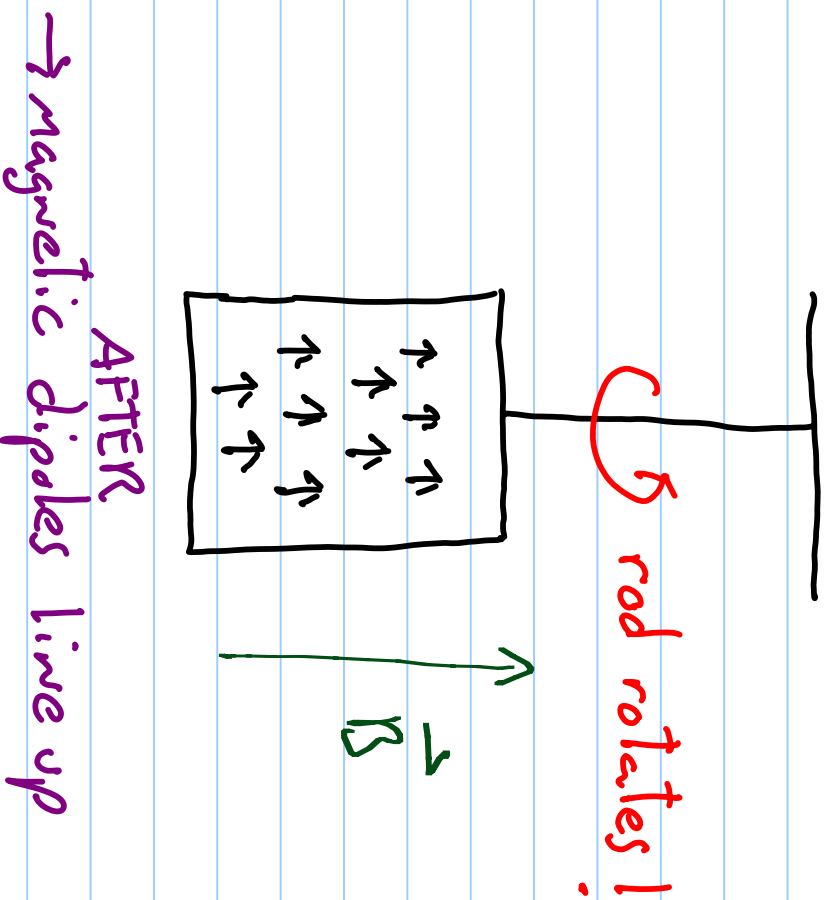
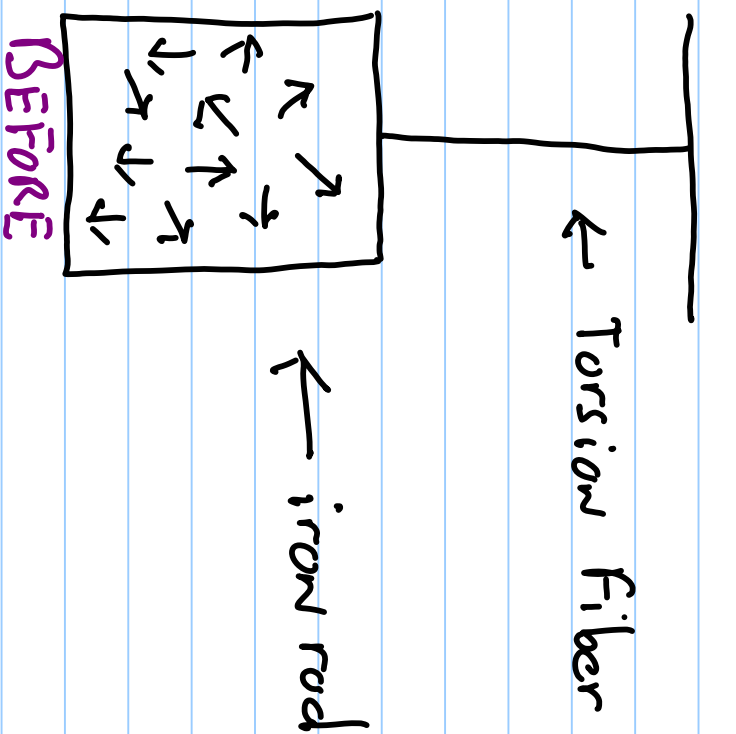
- What is "intrinsic spin"
- What is a gyromagnetic ratio
- How this relates to the Stern-Gerlach results
- What are the relevant operators for spin
- How to derive and work with the Pauli matrices (example problem)

(Roughly covers chapter 6.8 of the textbook)

MAGNETISM AND ANGULAR MOMENTUM

(2)

Einstein - de Haas EFFECT



→ this effect links angular momentum and magnetism. It does not demonstrate quantisation

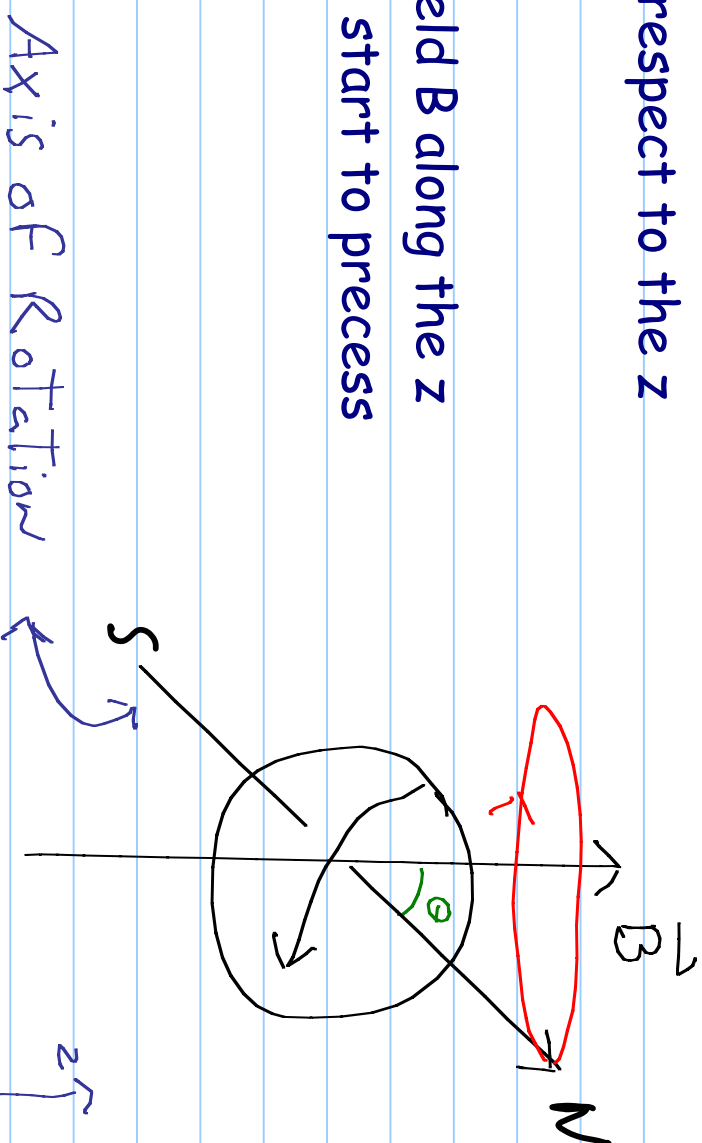
→ magnetic dipoles line up

RECAP FROM LECTURE 7

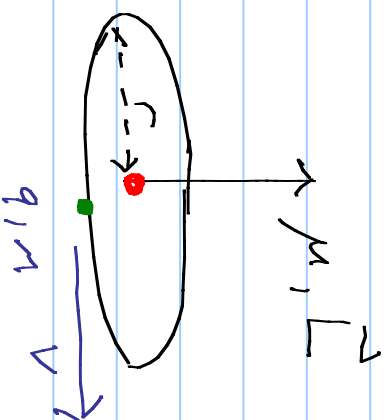
(3)

Consider a charged sphere whose rotation axis is at an angle θ with respect to the z axis

By applying a magnetic field B along the z direction, the sphere will start to precess about the z axis

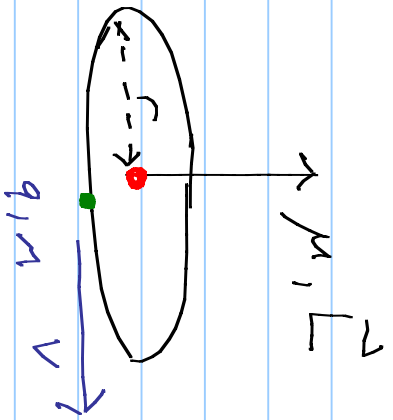


Now consider a one-electron atom in the Bohr Model:



- q : charge of electron = $-e$
- m : mass of electron = m_e
- r : radius of orbit
- v : velocity of electron
- μ : magnetic moment of atom
- L : angular momentum

How long does it take for the electron to orbit? (4)



$$\rightarrow \frac{2\pi r}{v} = T_{\text{orbit}}$$

$$\Rightarrow \text{current } I = \frac{q}{T_{\text{orbit}}} = \frac{qv}{2\pi r}$$

You'll see in Eq 1 that the magnetic moment for a dipole is

$$\rightarrow \mu = \pi r^2 \cdot \frac{qv}{2\pi r} = \boxed{\frac{qvr}{2}} \quad \mu = I \cdot \text{area} = \pi r^2 I$$

The orbital angular momentum: $L = mvr$

$$\Rightarrow \mu = \frac{q}{2m} L, \text{ with } \hbar \text{ natural unit of } L,$$

We can write:

$$\boxed{\mu_B = \frac{e\hbar}{2m_e}}$$

$\rightarrow \mu_B$ is "Bohr magneton"

Gyromagnetic Ratio

(5)

$\mu = \frac{q}{2m} L$ Tells us how the magnetic moment and the angular momentum are related

$\frac{q}{2m}$: gyromagnetic ratio

The electron's intrinsic angular momentum i.e "spin" is equal to $\frac{1}{2} \hbar$. Using this value, we should get

$$\mu = \frac{e}{2m_e} \hbar/2$$

But we get $\mu = \frac{e \hbar}{2m_e}$ or $\mu = g \cdot \frac{e}{2m_e} \hbar/2$

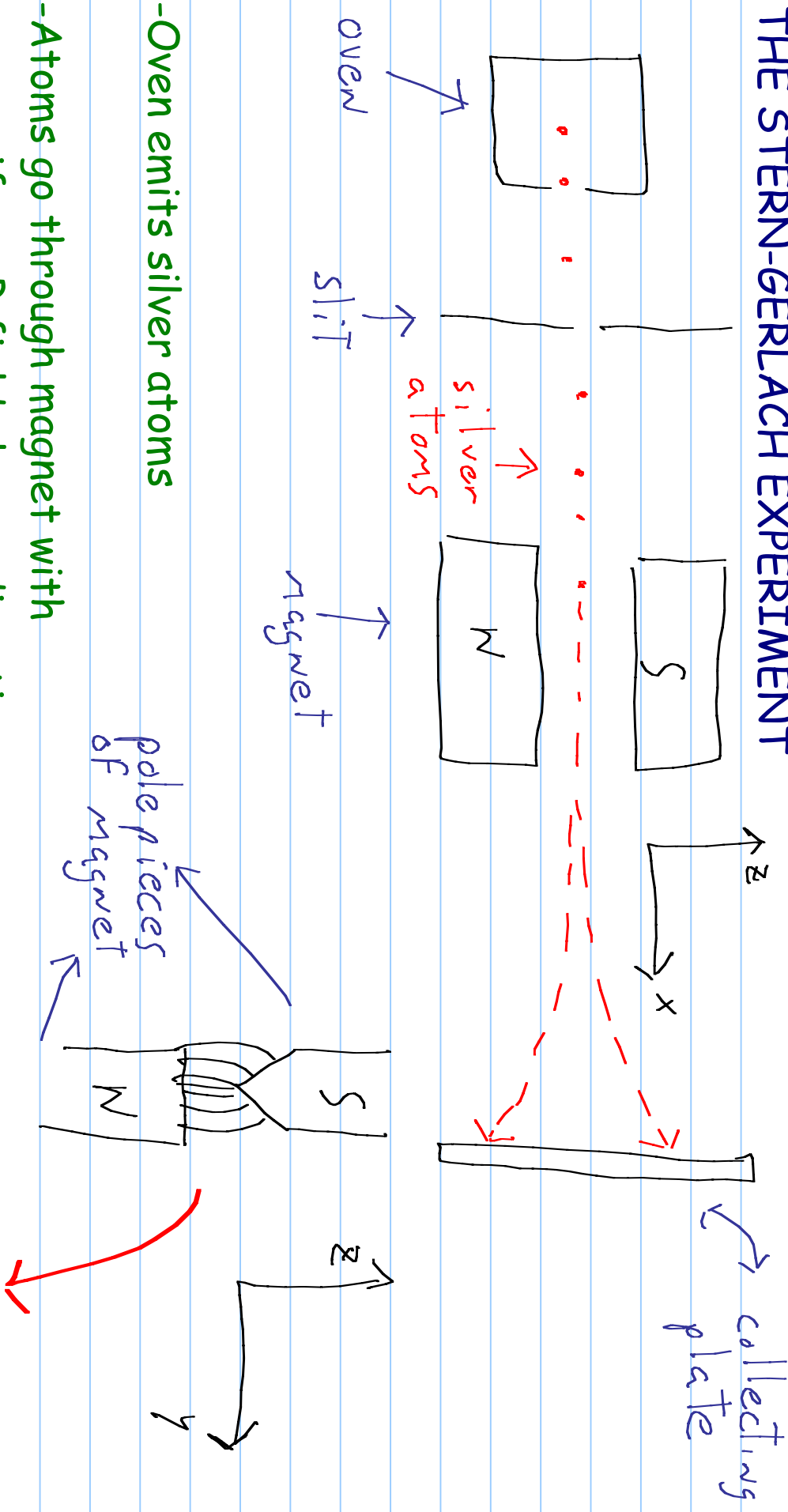
$$g = 2.0023193043766 \quad (87) \quad \rightarrow \text{measured}$$

\rightarrow see electron spin resonance in PHY 225

(See lecture 7)

©

THE STERN-GERLACH EXPERIMENT



-Oven emits silver atoms

-Atoms go through magnet with non-uniform B field along z direction

-Trajectory of atoms is bent is proportion to M_z

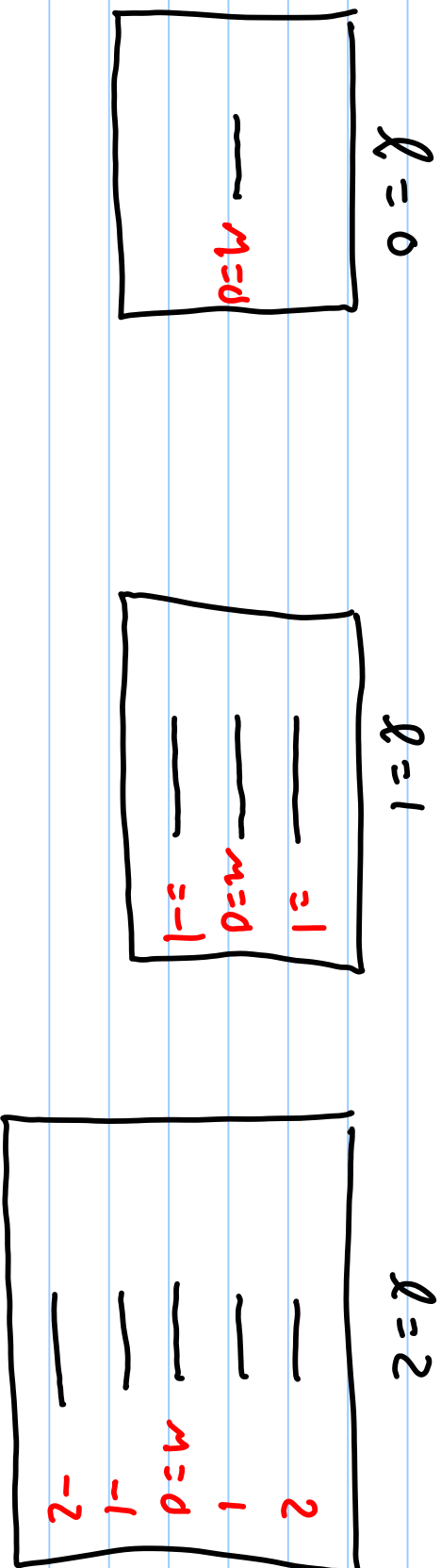
-Observe pattern of silver atoms on collecting plate

$M_z = \mu_B m_l$
B field!

SPIN ANGULAR MOMENTUM

⑦

FROM WHAT WE HAVE SEEN UP UNTIL NOW WITH ORBITAL ANGULAR MOMENTUM, WE SHOULD OBSERVE:



WHAT WE ARE MISSING IS "SPIN". IT HAS NO CLASSICAL ANALOG

WE HAVE FOUND "ELEMENTARY" PARTICLES THAT HAVE SPIN $1/2$ AND SPIN 1 .

↓
FERMIONS (half integer) ↓ BOSONS (integer)

Gravitons have in theory spin 2. We are looking for a spin 0 particle (Higgs Boson).

BUT we were using silver 1?

⑧

Silver (Ag) has 47 electrons

47th electron occupies 5s orbital.

All other electrons are paired

SPLITTING OF THE BEAM INTO TWO LINES is DUE TO THIS UNPAIRED ELECTRON

NOTE: - SPACE QUANTISATION (ORBITAL MOMENTUM) EXPERIMENT: ZEEMAN EFFECT

Spin $1/2$ EXPERIMENT: ELECTRON SPIN RESONANCE

Periodic Table of the Elements

1	H	2	He																																		
3	Li	4	Be	5	B	6	C	7	N	8	O	9	F	10	Ne																						
11	Na	12	Mg	13	Al	14	Si	15	P	16	S	17	Cl	18	Ar																						
19	K	20	Ca	21	Sc	22	Ti	23	V	24	Cr	25	Mn	26	Fe	27	Co	28	Ni	29	Cu	30	Zn	31	Ga	32	Ge	33	As	34	Se	35	Br	36	Kr		
37	Rb	38	Sr	39	Y	40	Zr	41	Nb	42	Mo	43	Tc	44	Ru	45	Rh	46	Pd	47	Ag	48	Cd	49	In	50	Sn	51	Sb	52	Te	53	I	54	Xe		
55	Cs	56	Ba	57	La	58	Hf	59	Ta	60	W	61	Re	62	Os	63	Ir	64	Pt	65	Au	66	Hg	67	Tl	68	Pb	69	Bi	70	Po	71	At	72	Rn		
87	Fr	88	Ra	89	Ac	90	Rf	91	Ha	92	Sg	93	Hs	94	Mt	95	110	111	112	113	114	115	116	117	118	119	120	121	122	123	124	125	126	127	128	129	130

*Lanthanide Series: Ce, Pr, Nd, Pm, Sm, Eu, Gd, Tb, Dy, Ho, Er, Tm, Yb, Lu
*Actinide Series: Th, Pa, U, Np, Pu, Am, Cm, Bk, Cf, Es, Fm, Md, No, Lr

SPIN ANGULAR MOMENTUM

(9)

WE REPEAT WHAT WE DID IN LECTURES 29 AND 30

→ WE SAW THAT ORBITAL ANGULAR MOMENTUM WAS CHARACTERIZED BY TWO QUANTUM NUMBERS: l, m

→ FOR SPIN WE WILL WORK WITH TWO NUMBERS ALSO: s, m_s AND TWO OPERATORS \hat{S}_x, \hat{S}_z

→ SINCE m_s TAKES ON TWO VALUES, $s = 1/2, m_s \pm 1/2$

WITH $|2\rangle$ BEING EIGENSTATE OF \hat{S}_z, \hat{S}_z

WE HAVE: $\hat{S}_z^2 |2\rangle = \hbar^2 s(s+1) |2\rangle$

$$\hat{S}_z |2\rangle = m_s \hbar |2\rangle$$

$$\Rightarrow [\hat{S}_z, \hat{S}_z] = 0, [\hat{S}_x, \hat{S}_y] = i\hbar \hat{S}_z$$

Spin (cont.)

(10)

We re-introduce NEW RAISING AND LOWERING OPERATORS:

$$\begin{aligned} S_+ &= S_x + iS_y \\ S_- &= S_x - iS_y \end{aligned}$$

$$\begin{aligned} S_+^\dagger &= S_- \\ S_-^\dagger &= S_+ \end{aligned}$$

We skip the DERIVATION (see lecture 30)

$$\left. \begin{aligned} S_- |s, m_s\rangle &= \alpha |s, m_s - 1\rangle \\ S_+ |s, m_s\rangle &= \alpha |s, m_s + 1\rangle \end{aligned} \right\}$$

s stays the same
 m_s is raised or lowered by 1

$$S_- |s, m_s\rangle = c_- |s, m_s - 1\rangle$$

$$\langle s, m_s | S_-^\dagger = \langle s, m_s - 1 | c_-^* = \langle s, m_s | S_+^\dagger$$

$$\Rightarrow \langle s, m_s | S_+^\dagger S_- |s, m_s\rangle = |c_-|^2$$

Spin (cont.)

(11)

$$S_+ S_- = (S_x + iS_y)(S_x - iS_y)$$

$$= S_x^2 + S_y^2 - S_x iS_y + iS_y S_x$$

$$= S_x^2 + S_y^2 - i[S_x, S_y]$$

$$= S_x^2 + S_y^2 - i \cdot i \hbar S_z$$

$$= S_x^2 + S_y^2 + \hbar S_z$$

$$\langle s, m_s | S_x^2 + S_y^2 + \hbar S_z | s, m_s \rangle = |c_1|^2$$

$$\langle s, m_s | s, m_s \rangle \cdot [s(s+1)\hbar^2 - m_s^2\hbar^2 + m_s\hbar^2] = |c_1|^2$$

$$= \hbar^2 [s(s+1) - m_s^2 + m_s] = |c_1|^2$$

$$c_1 = \hbar \sqrt{s(s+1) - m_s(m_s - 1)}$$

$$\Rightarrow \boxed{S_- |s, m_s\rangle = \hbar \sqrt{s(s+1) - m_s(m_s - 1)} |s, m_s - 1\rangle}$$

①

Spin (cont.)

(12)

$$S_+^1 |s, m_s\rangle = \hbar \sqrt{s(s+1) - m_s(m_s+1)} |s, m_s+1\rangle \quad (2)$$

$$\Rightarrow \sqrt{\frac{1}{2}(\frac{1}{2}+1) - (-\frac{1}{2}(-\frac{1}{2}+1))} = \sqrt{\frac{3}{4} + \frac{1}{2}} = 1$$

In The S_z basis: \rightarrow Same result for C_+
 \rightarrow S_z eigenstate

$$S_z^1 |+\rangle = \frac{\hbar}{2} |+\rangle$$

$$S_z^1 |-\rangle = -\frac{\hbar}{2} |-\rangle$$

$$S_z^1 : \begin{pmatrix} \frac{\hbar}{2} & 0 \\ 0 & -\frac{\hbar}{2} \end{pmatrix}, \text{ eigenvectors: } |+\rangle : \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$|-\rangle : \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

From (1)(2) : $S_+^1 |-\rangle = \hbar |+\rangle$

$$S_-^1 |+\rangle = \hbar |-\rangle$$

Spin (cont.)

(13)

Matrix Elements of S_+^1 :

$$\langle +1 | S_+^1 | + \rangle = 0, \quad \langle +1 | S_+^1 | - \rangle = \hbar$$

$$\langle -1 | S_+^1 | + \rangle = 0, \quad \langle -1 | S_+^1 | - \rangle = 0$$

$$: \quad \hbar \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}, \quad S_-^1 : \quad \hbar \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

Back to the definition of S_+^1 and S_-^1 :

$$\hat{S}_x = \frac{1}{2} (S_+^1 + S_-^1), \quad S_y^1 = \frac{1}{2i} (S_+^1 - S_-^1)$$

$$S_x^1 : \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad S_y^1 : \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

Pauli Matrices: $\sigma_x : \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\sigma_y : \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $\sigma_z : \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

Spin (cont.)

(14)

WHAT ARE THE EIGENVECTORS OF S_x ?

$$\begin{vmatrix} -\lambda & \hbar/2 \\ \hbar/2 & -\lambda \end{vmatrix} = 0 \quad \lambda^2 - \frac{\hbar^2}{4} = 0 \Rightarrow \lambda = \pm \hbar/2$$

• with $\lambda = +\hbar/2$:

$$\begin{pmatrix} 0 & \hbar/2 \\ \hbar/2 & 0 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \hbar/2 \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} \rightarrow$$

$$\begin{aligned} \hbar/2 x_2 &= \hbar/2 x_1 \\ \hbar/2 x_1 &= \hbar/2 x_2 \end{aligned}$$

$$\Rightarrow |x\rangle_+ = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

• with $\lambda = -\hbar/2$

$$\begin{pmatrix} 0 & \hbar/2 \\ \hbar/2 & 0 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = -\hbar/2 \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} \rightarrow \begin{aligned} \hbar/2 x_2 &= -\hbar/2 x_1 \\ \hbar/2 x_1 &= -\hbar/2 x_2 \end{aligned}$$

(15)

$$\Rightarrow |-\rangle_+ = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

$$|+\rangle_x = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}, \quad |-\rangle_x = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

NOTE THAT WITH: $|+\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$, $|-\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

WE CAN WRITE:

$$|+\rangle_x = \frac{1}{\sqrt{2}} (|+\rangle + |-\rangle)$$

$$|-\rangle_x = \frac{1}{\sqrt{2}} (|+\rangle - |-\rangle)$$

i.e. $|+\rangle_x$ eigenstate even mixture of $|+\rangle$ and $|-\rangle$

EXERCISE: FIND EIGENVALUES AND EIGENVECTORS OF S_y^1

EXAMPLE PROBLEM:

(16)

- FIND THE EIGENVALUES AND EIGENSTATES OF THE SPIN OPERATOR \hat{S} OF AN ELECTRON IN THE DIRECTION OF A UNIT VECTOR \vec{n} , WHERE \vec{n} LIES IN THE XZ PLANE.

WE WANT TO SOLVE:

$$\vec{n} \cdot \hat{S} |\varphi\rangle = \frac{\hbar}{2} \lambda |\varphi\rangle$$

$$\vec{n} \cdot \hat{S} = (\sin \theta \hat{x} + \cos \theta \hat{z}) \cdot (\hat{S}_x \hat{x} + \hat{S}_y \hat{y} + \hat{S}_z \hat{z})$$

\downarrow
unit vectors

$$\hat{S}_x = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \hat{S}_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \hat{S}_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\vec{n} \cdot \hat{S} = \frac{\hbar}{2} \begin{pmatrix} \cos \theta & \sin \theta \\ \sin \theta & -\cos \theta \end{pmatrix}$$

(17)

$$\begin{vmatrix} \frac{h}{2} \cos \theta - \lambda & \frac{h}{2} \sin \theta \\ \frac{h}{2} \sin \theta & -\frac{h}{2} \cos \theta - \lambda \end{vmatrix} = 0$$

$$-\frac{h^2}{4} \cos^2 \theta - \lambda \frac{h}{2} \cos \theta + \lambda \frac{h}{2} \cos \theta + \lambda^2 - \frac{h^2}{4} \sin^2 \theta = 0$$

$$= -\frac{h^2}{4} + \lambda^2 = 0, \quad \lambda = \pm \frac{h}{2}$$

Eigenvectors:

$$\frac{h}{2} \begin{pmatrix} \cos \theta & \sin \theta \\ \sin \theta & -\cos \theta \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \frac{h}{2} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

$$\begin{aligned} \cos \theta x_1 + \sin \theta x_2 &= x_1 \\ \sin \theta x_1 - \cos \theta x_2 &= x_2 \end{aligned} \quad \begin{matrix} \textcircled{4} \\ \textcircled{5} \end{matrix}$$

PROBLEM (Cont.)

(18)

$$x_2 \sin \theta = x_1 (1 - \cos \theta)$$

$$x_1 \sin \theta = x_2 (1 + \cos \theta)$$

(6)

(7)

$$(6) \rightarrow \frac{x_2}{x_1} = \frac{(1 - \cos \theta)}{\sin \theta}, \quad \tan \frac{\theta}{2} = \frac{(1 - \cos \theta)}{\sin \theta}$$

$$\frac{x_2}{x_1} = \tan \frac{\theta}{2} = \frac{\sin \theta/2}{\cos \theta/2}, \quad x_1 = \cos \theta/2, \quad x_2 = \sin \theta/2$$

$$| \phi_1 \rangle = \begin{pmatrix} \cos \theta/2 \\ \sin \theta/2 \end{pmatrix} \rightarrow \text{normalised}$$

PROCEEDING THE SAME WAY WE WOULD GET:

$$| \phi_2 \rangle = \begin{pmatrix} -\sin \theta/2 \\ \cos \theta/2 \end{pmatrix}, \quad \langle \phi_1 | \phi_2 \rangle = 0 \rightarrow \text{orthogonal}$$

- WHAT IS THE PROBABILITY OF MEASURING $S_z = \pm \hbar/2$ IN THE SYSTEM IS IN THE $| \phi_1 \rangle$ EIGENSTATE?

PROBLEM CONT.

(19)

LET'S WRITE $|q_1\rangle$ IN TERMS OF THE S_z^1 EIGENSTATES:

$$|q_1\rangle = \cos \frac{\theta}{2} |+\rangle + \sin \frac{\theta}{2} |-\rangle$$

$$\downarrow \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\downarrow \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$P = |\langle + | q_1 \rangle|^2$$

$$= |\langle + | \cos \frac{\theta}{2} |+\rangle + \sin \frac{\theta}{2} |-\rangle|^2$$

$$= \cos^2 \frac{\theta}{2}$$

