

# Higgs Boson Physics: The Past, the Present, and the Future

Pierre Savard

University of Toronto and TRIUMF

TRIUMF Colloquium

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# OUTLINE

Introduction: Higgs Physics

The Past: what we learned in Run 1

The Present: what we learned in Run 2

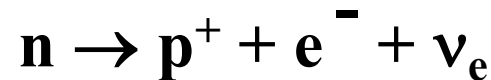
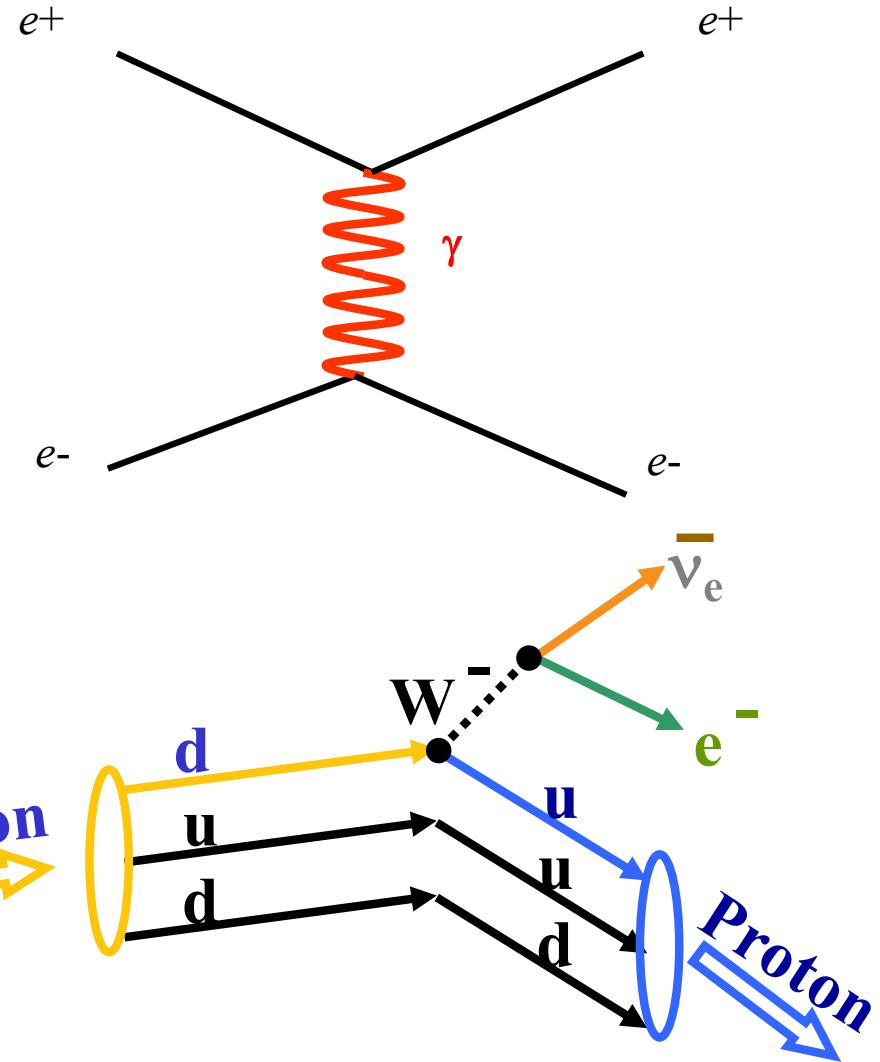
The Future: Run 3 and High Luminosity LHC

Conclusions

# Higgs Boson Physics

# The Electromagnetic and the Weak Forces

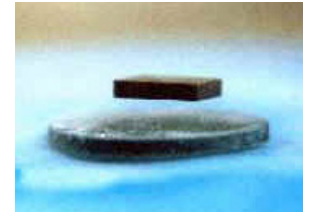
- A quantum theory of electromagnetism (QED) was developed at the end of the 40s
  - The force mediator is the massless photon
  - Theory is renormalizable
- The weak force appeared to require massive bosons
  - Massive mediators makes the theory non-renormalizable
  - This is the theoretical problem that Higgs\* and others solved
  - Mass of W boson acquired by invoking a new scalar field



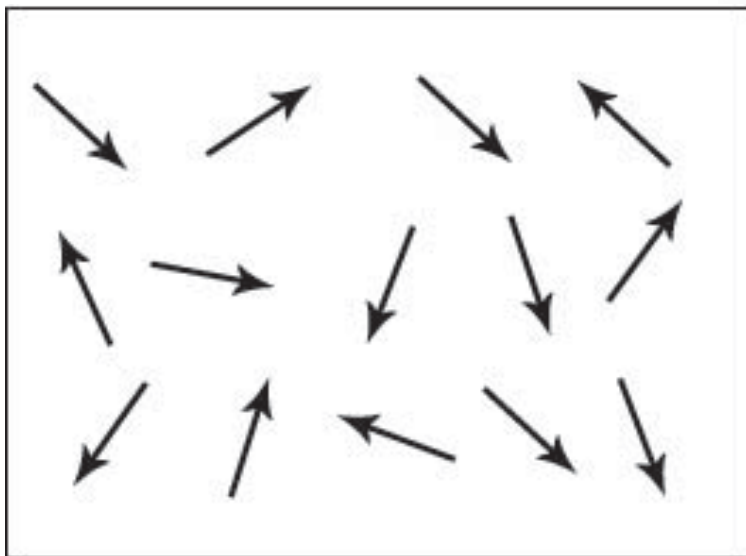
\*Higgs, Brout, Englert, Kibble, Hagen, Guralnik, and also Anderson

# Spontaneous Breaking of Electroweak Symmetry

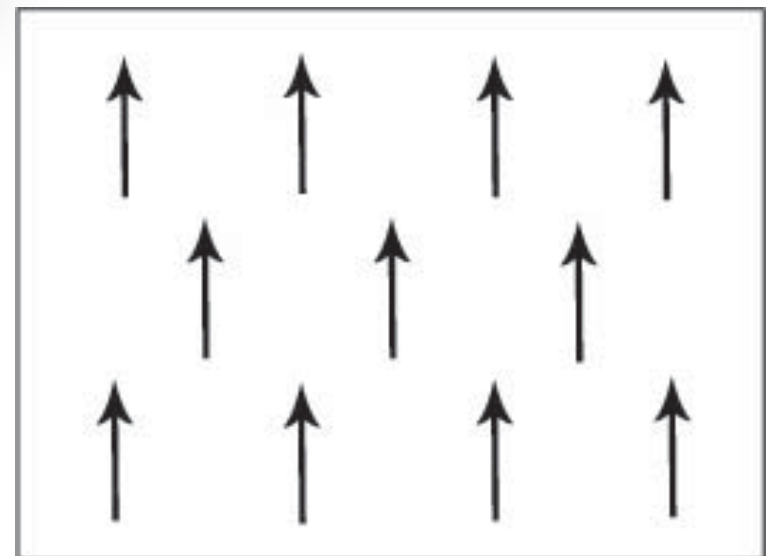
- Hypothesis: a phase transition occurred in the very early Universe: the Higgs field acquired a non-zero value
  - Similar to a phase transition in ferromagnets or a closer analogy would be the Meissner effect in superconductors



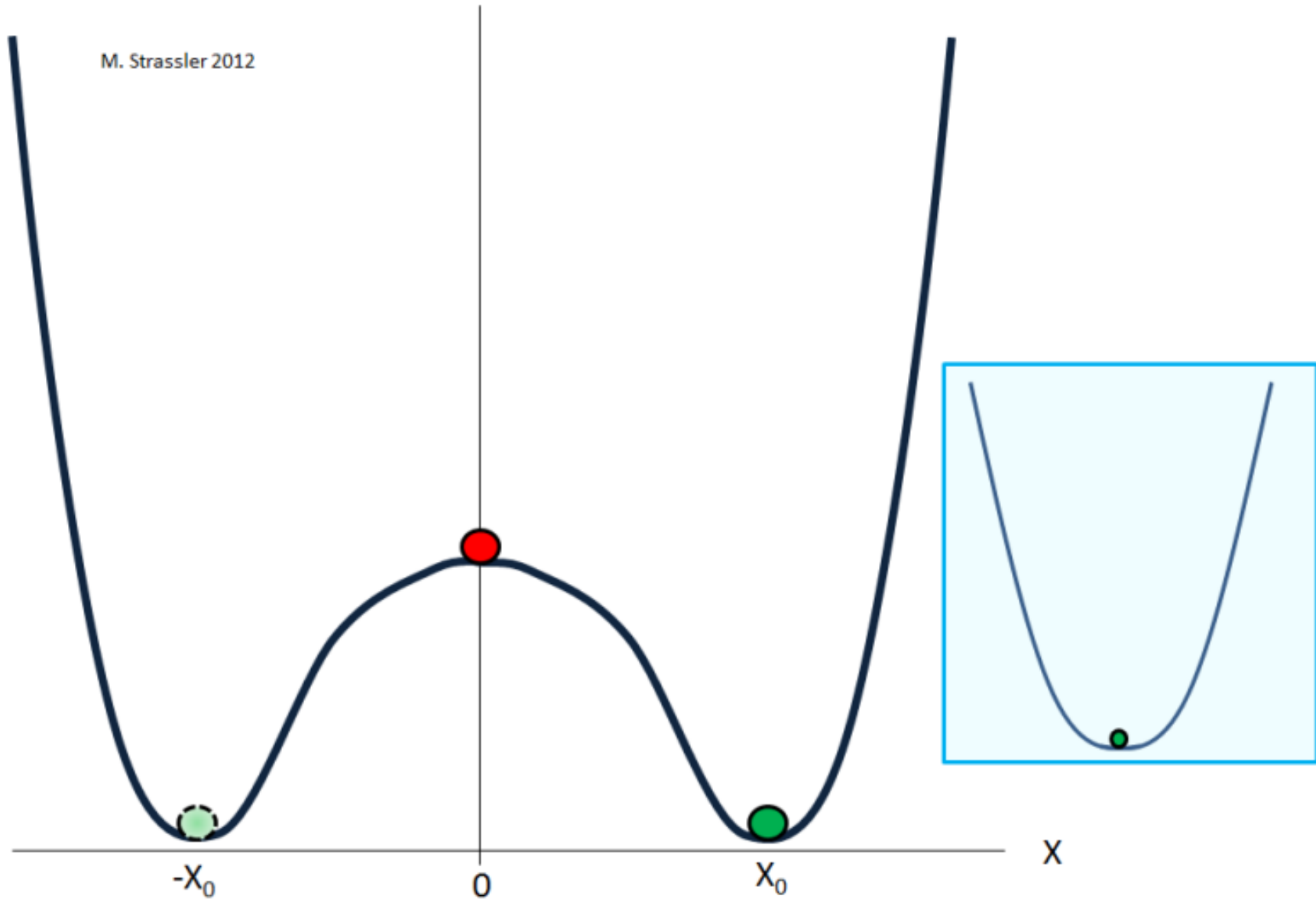
High Temperature



Low Temperature



# Spontaneous Breaking of Electroweak Symmetry



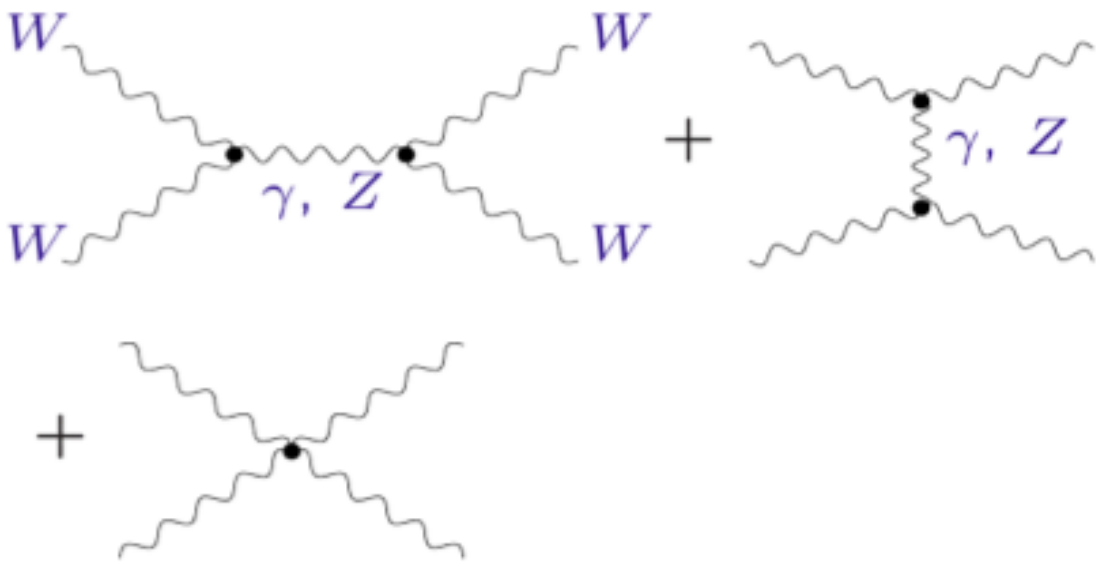
# The Higgs Field

- Unlike the electromagnetic field, this field is “on” everywhere in space
  - The physical vacuum is not “empty”, it is filled with the Higgs field everywhere
- The value of the field picked by Nature determines the physics (and chemistry!) of the Universe we live in
- Fundamental particles are subject to a new kind of force and the strength of this interaction with a given particle determines its mass



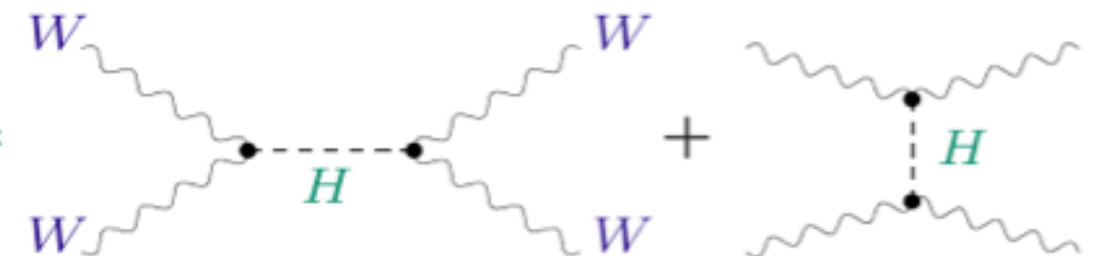
# WW Scattering and the Higgs Boson

WW scattering  
violates unitarity  
above  $\sim 1$  TeV

$$\mathcal{M}_V =$$


The diagram shows two Feynman diagrams for WW scattering. The first diagram is a t-channel exchange of a photon ( $\gamma$ ) or Z boson between two W bosons. The second diagram is a s-channel exchange of a photon ( $\gamma$ ) or Z boson between two W bosons. The labels  $\gamma, Z$  are in purple.

New diagrams  
needed to regulate  
the cross section

$$\mathcal{M}_S =$$


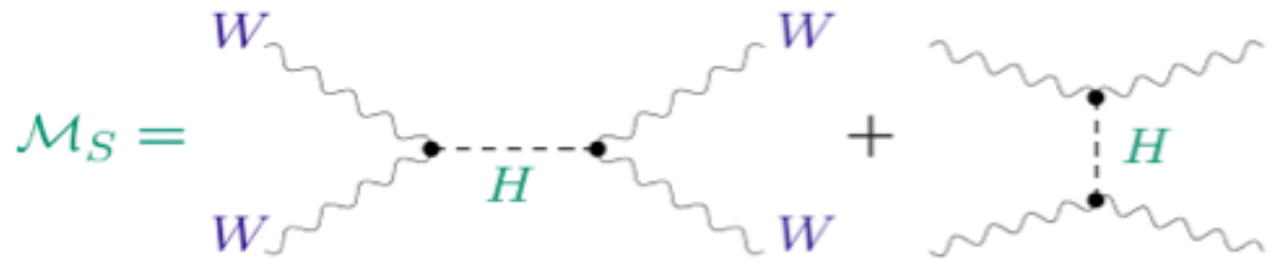
The diagram shows two Feynman diagrams for WW scattering. The first diagram is a t-channel exchange of a Higgs boson ( $H$ ) between two W bosons. The second diagram is a s-channel exchange of a Higgs boson ( $H$ ) between two W bosons. The labels  $H$  are in green.

Adding diagrams  
with a scalar solves  
the problem



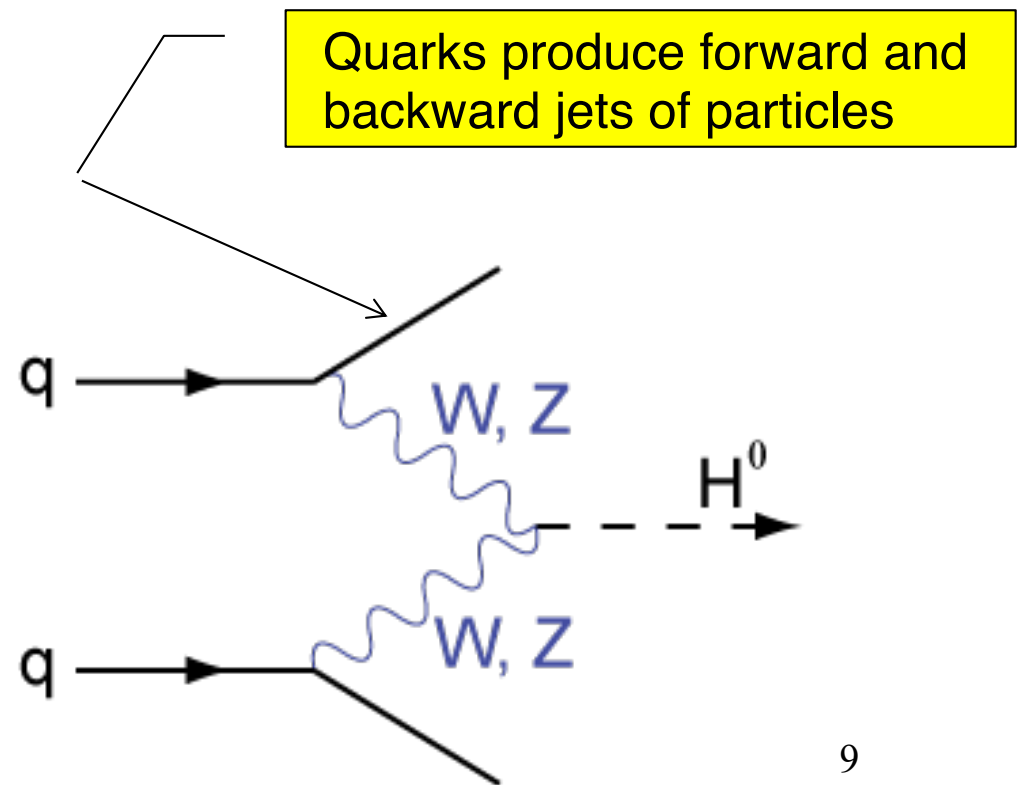
# WW Scattering and the Higgs Boson

WW scattering  
violates unitarity  
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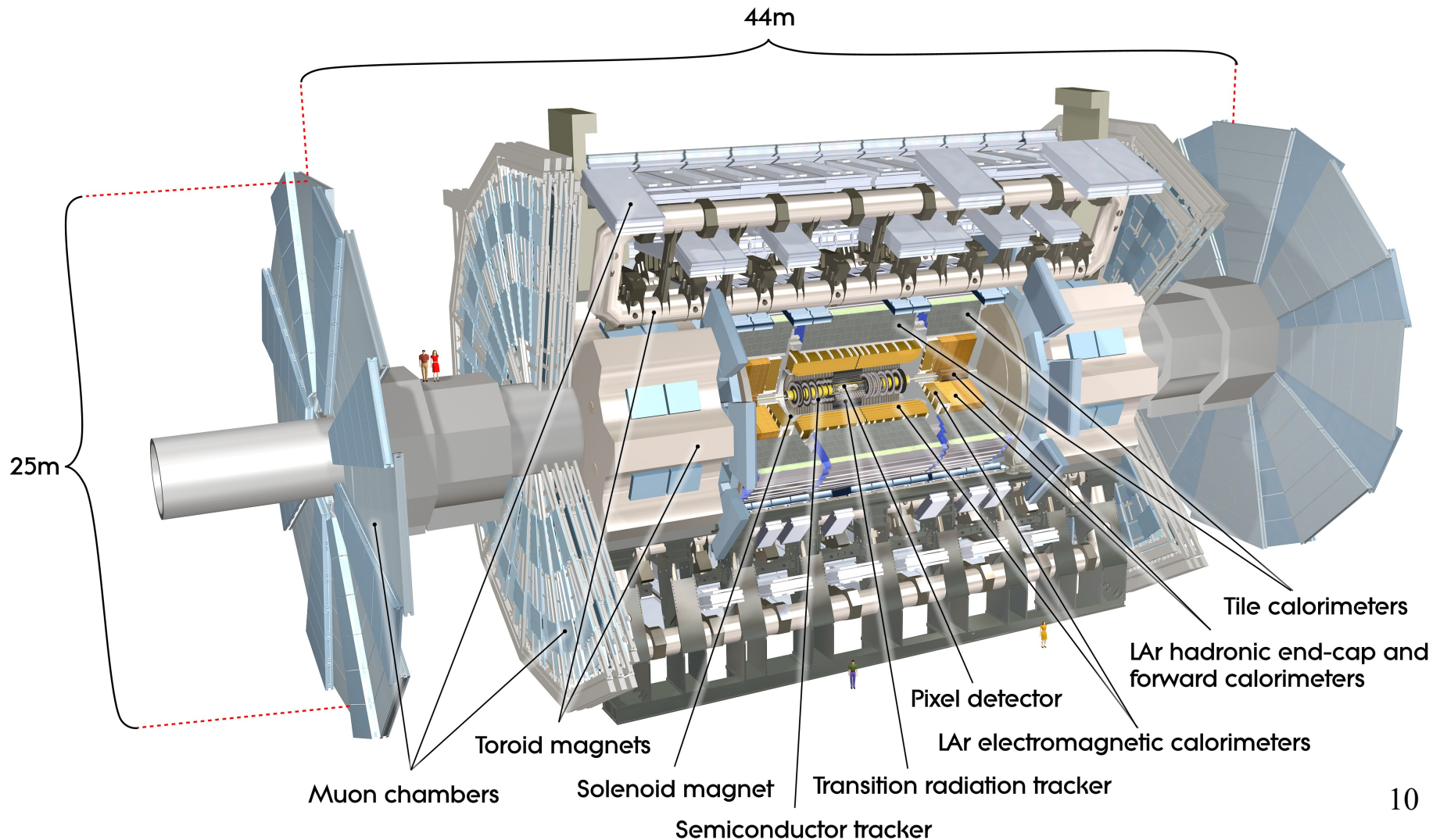
New diagrams  
needed to regulate  
the cross section

Adding Higgs  
diagrams solves the  
problem

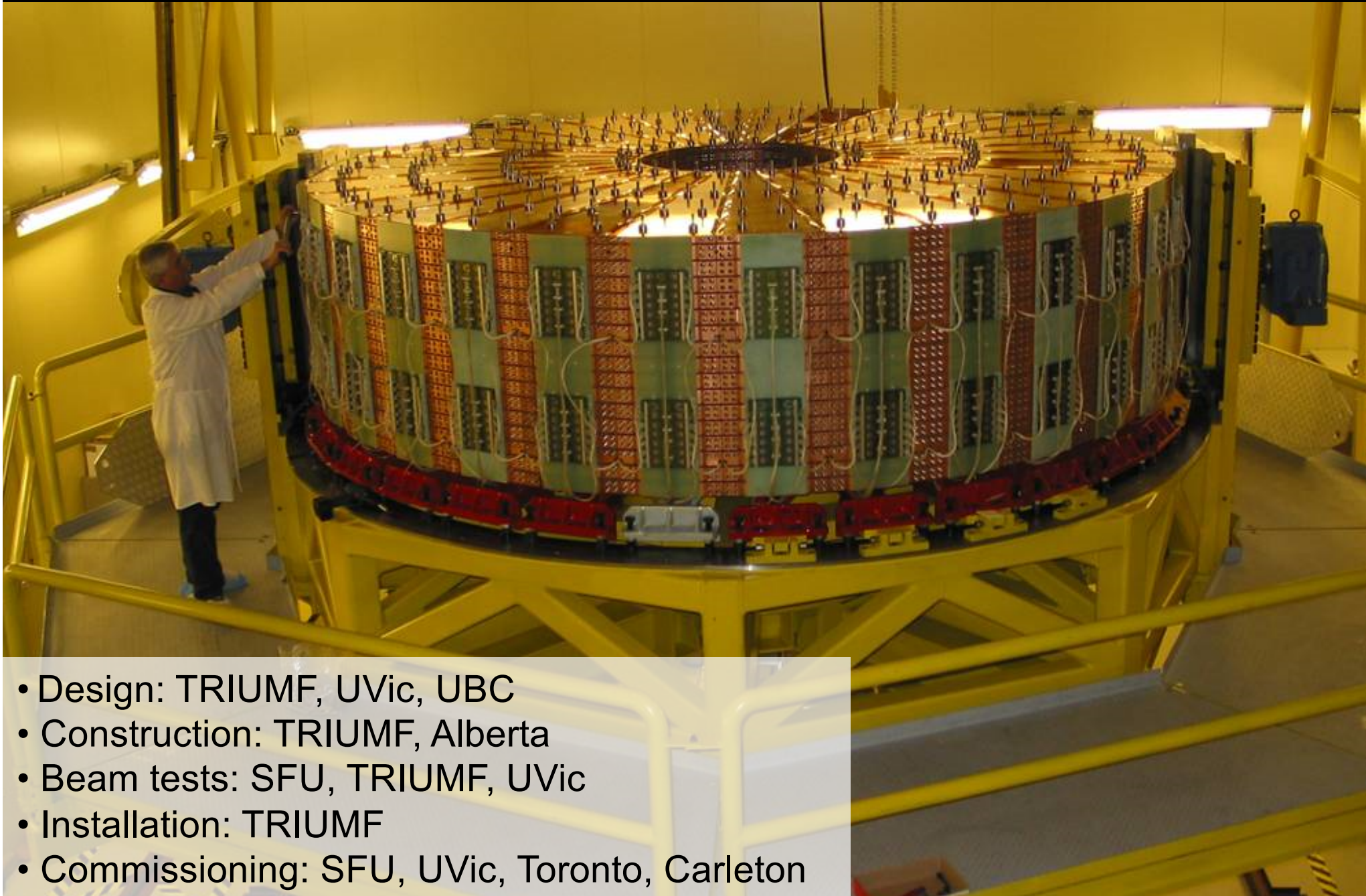


# ATLAS Detector

Designed to find or exclude the SM Higgs over a mass range up to 1 TeV



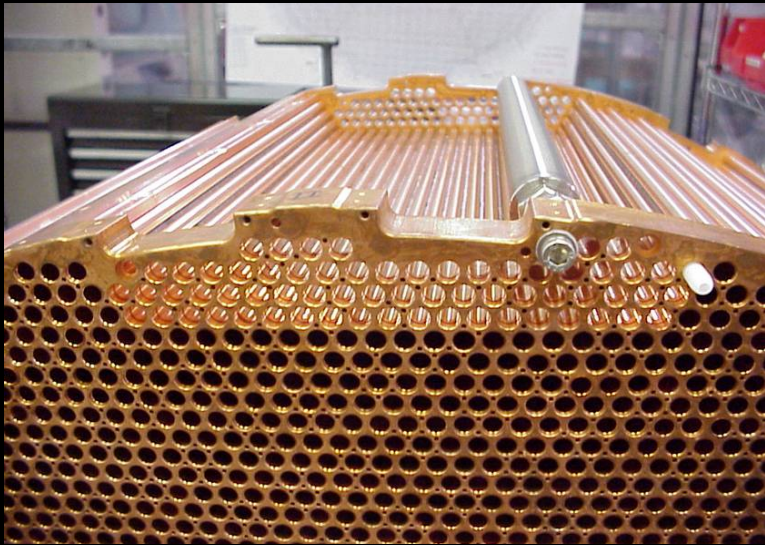
# Endcap Calorimeter



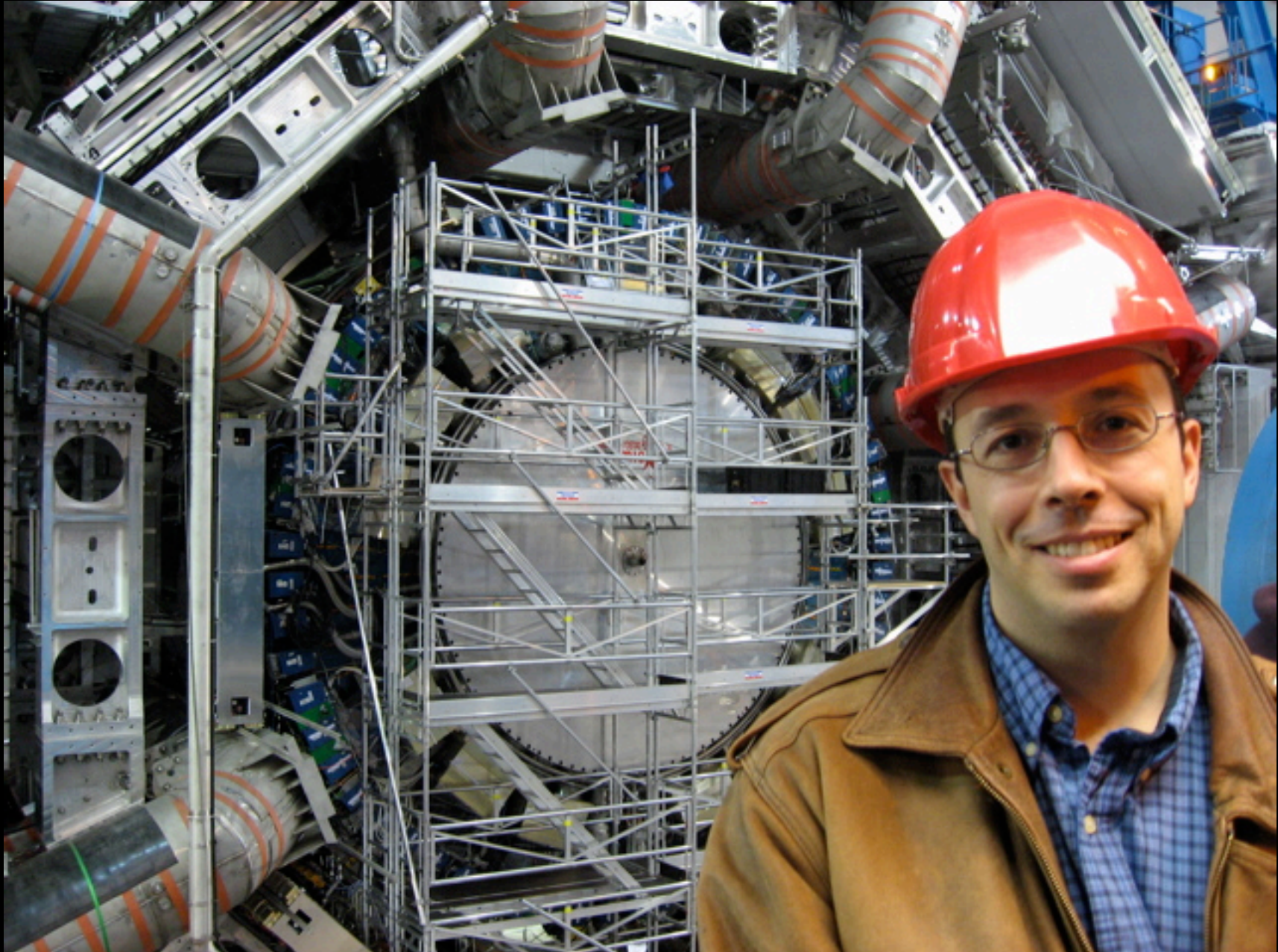
- Design: TRIUMF, UVic, UBC
- Construction: TRIUMF, Alberta
- Beam tests: SFU, TRIUMF, UVic
- Installation: TRIUMF
- Commissioning: SFU, UVic, Toronto, Carleton

# Forward Calorimeter

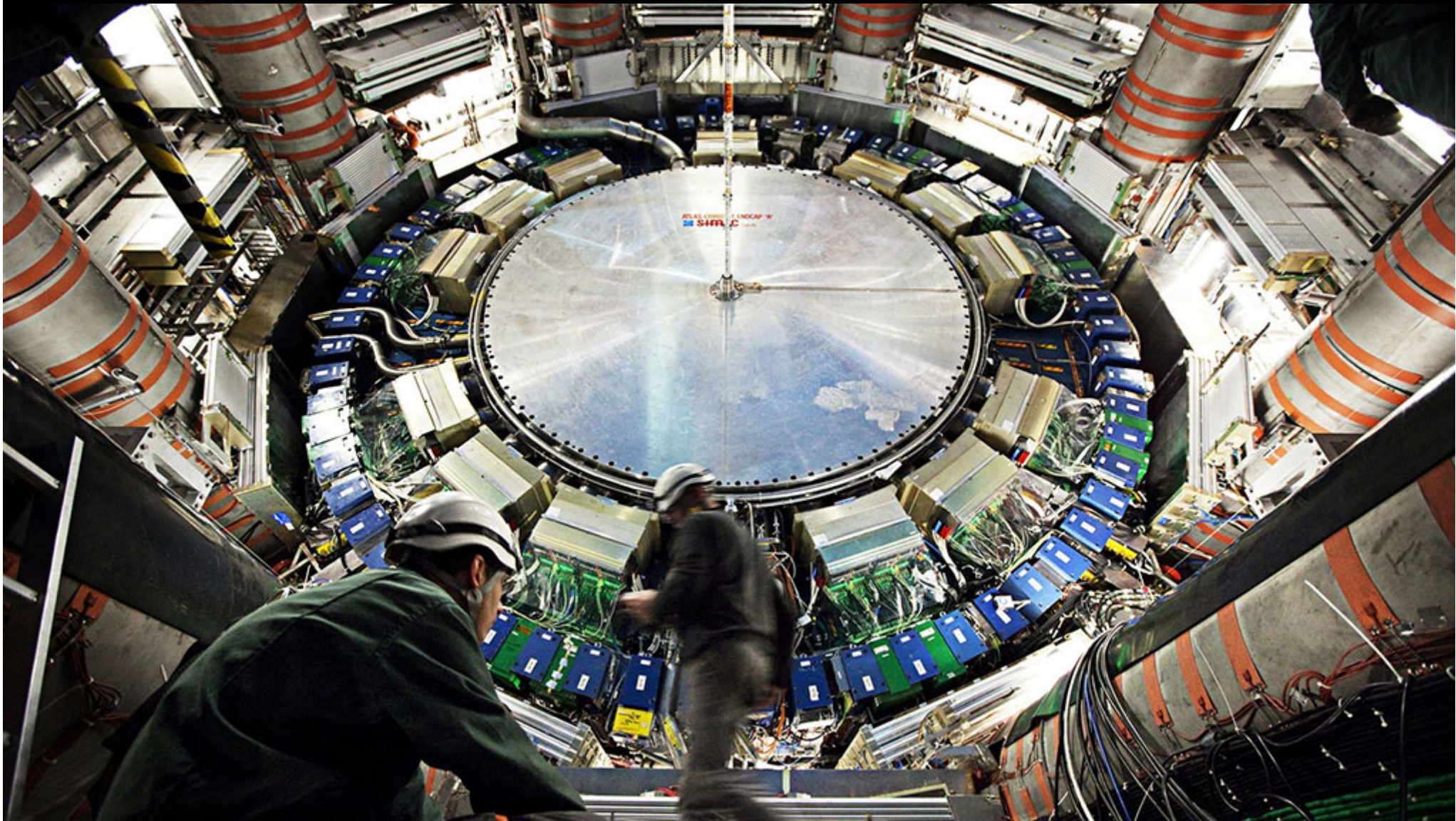
Design, Construction, Beam tests,  
Installation & Commissioning:  
Carleton, Toronto



# ATLAS Detector with Calorimeter

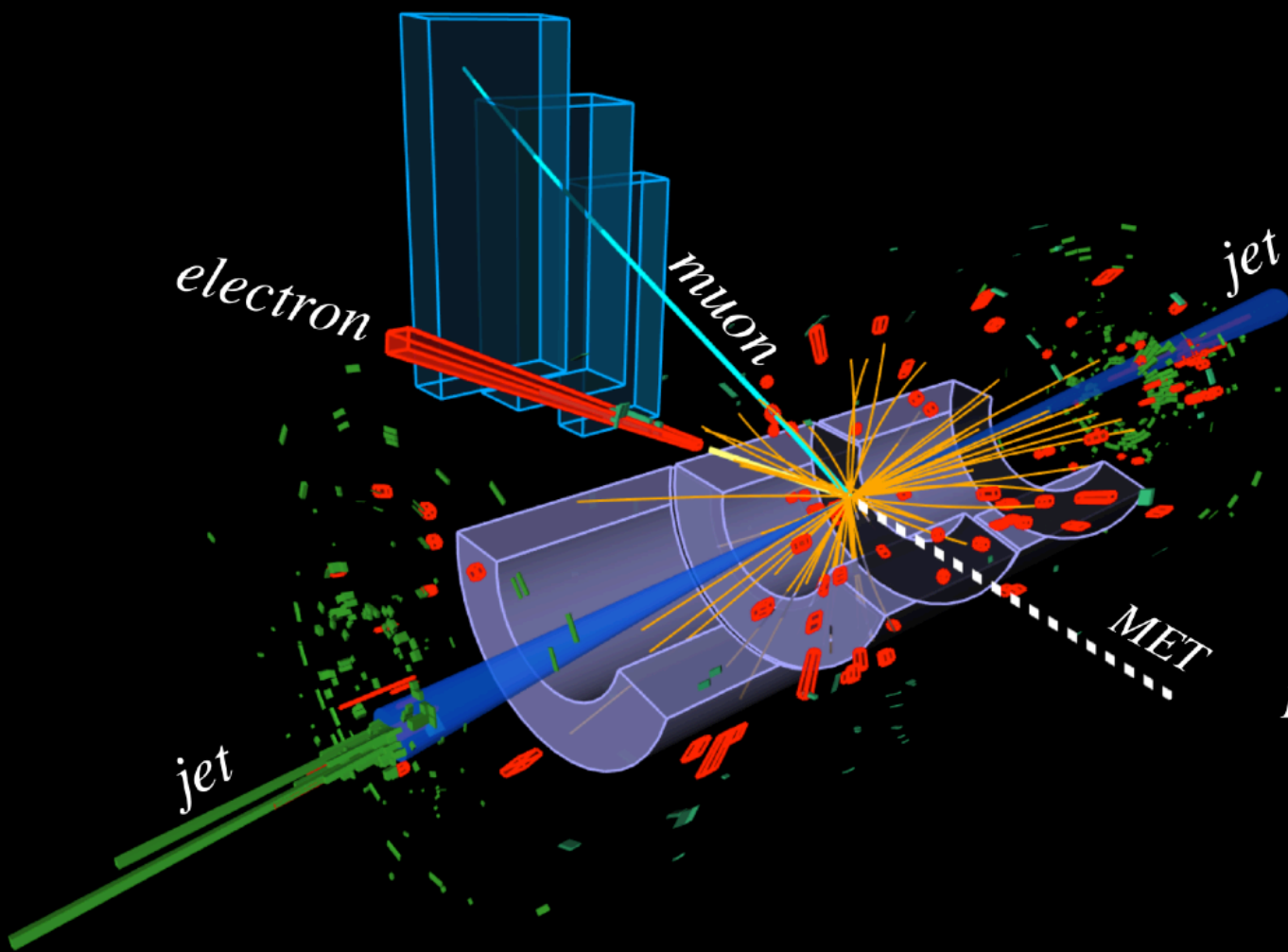


# ATLAS Detector with Calorimeter

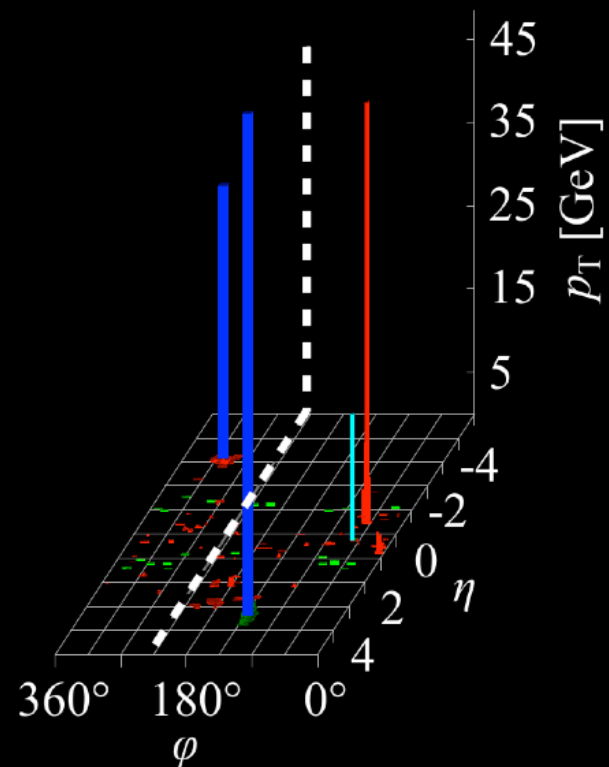


$H \rightarrow WW^* \rightarrow e\nu\mu\nu$  candidate and two jets with VBF topology

*Longitudinal view*



*Projected  $\eta$ - $\phi$  view*



Run 214680, Ev. no. 271333760

Nov. 17, 2012, 07:42:05 CET



<http://atlas.ch>

# The Standard Model

Some important dates

1967:

“A Model of Leptons”  
(Weinberg)

1971-73:

Renormalizability of theory,  
Quantum Chromodynamics  
asymptotic freedom

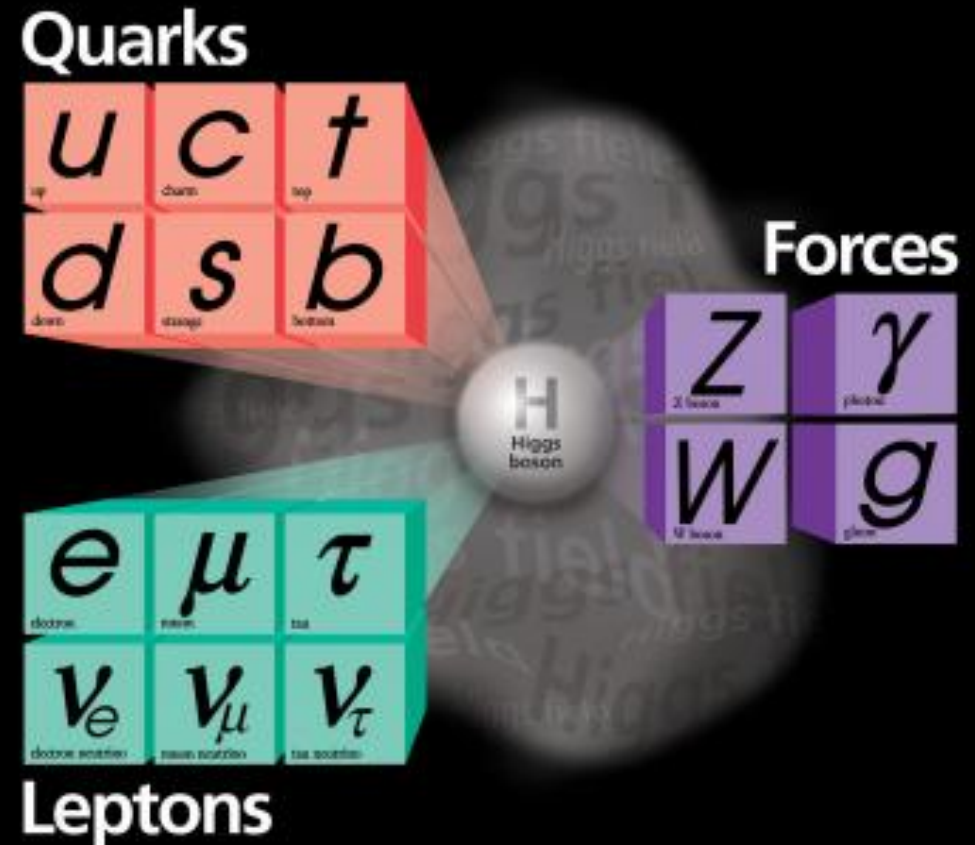
Mid-70s:

Discovery of 3<sup>rd</sup> generation  
fermions

1995:

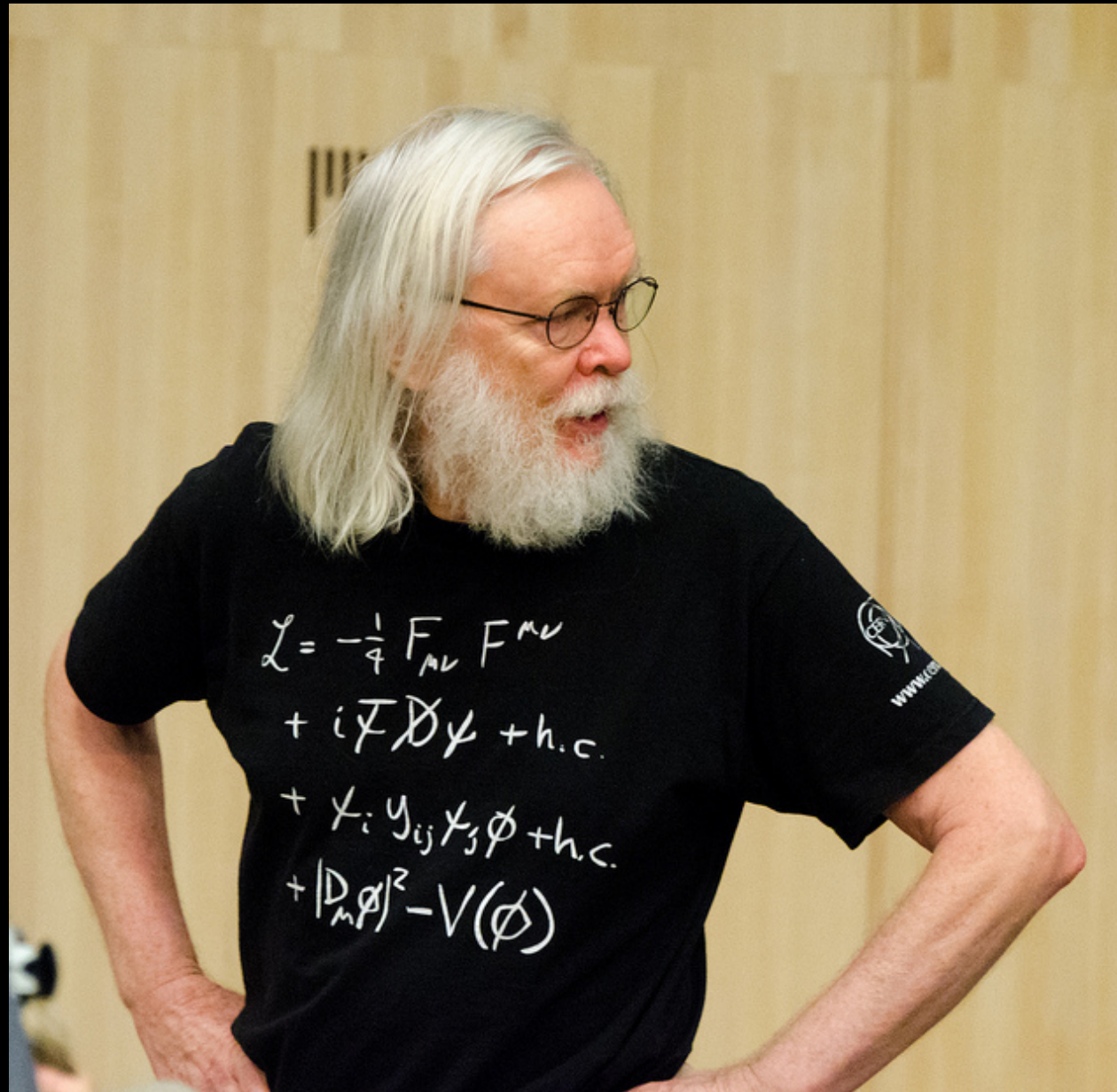
Top quark discovery,  
completed the third fermion  
generation\*

Until 2012 we were still missing a key  
Ingredient of the model: a scalar  
boson



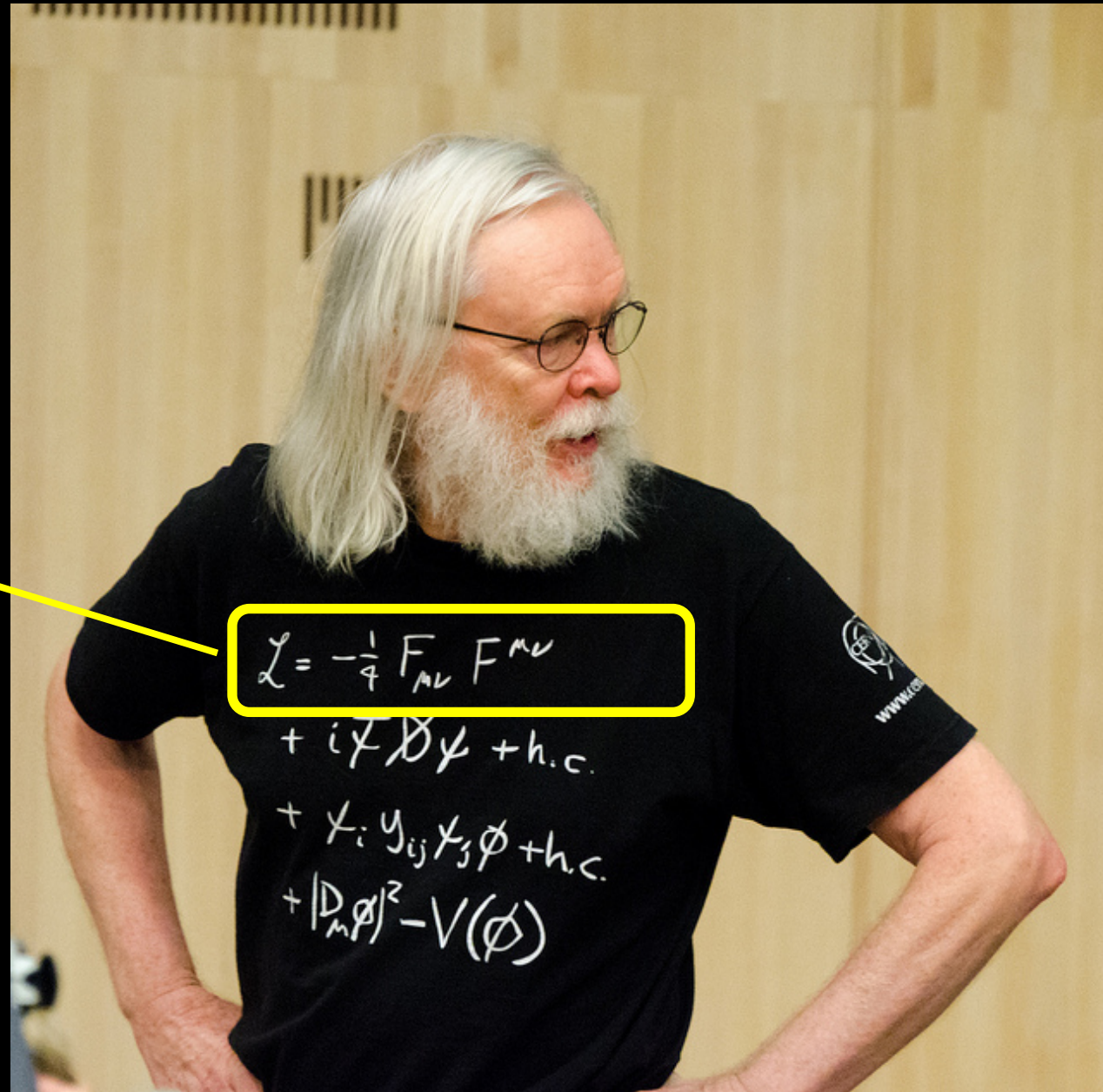


# The Standard Model Lagrangian



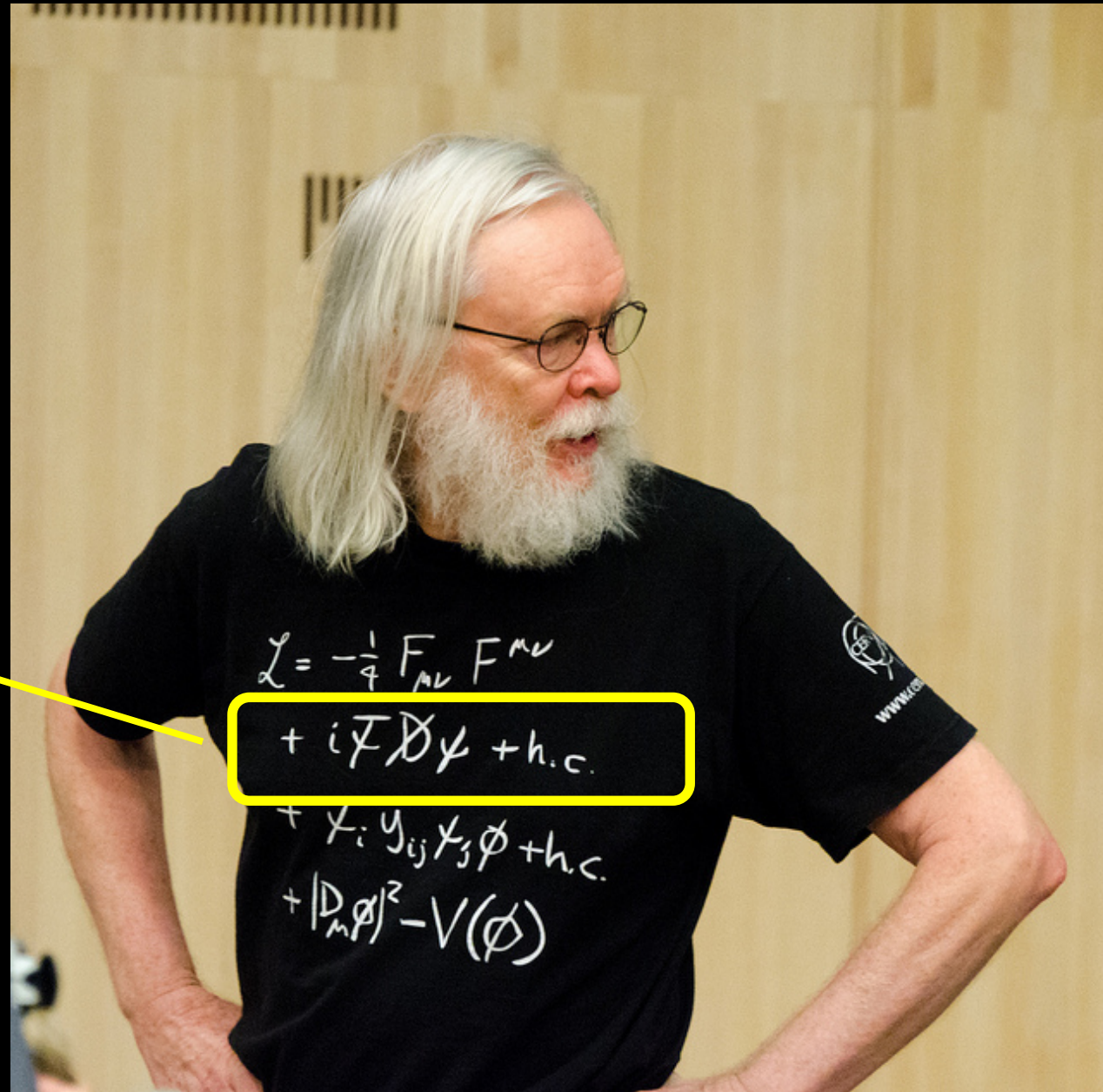
# The Standard Model on a T-Shirt

Kinetic term for  
gauge bosons



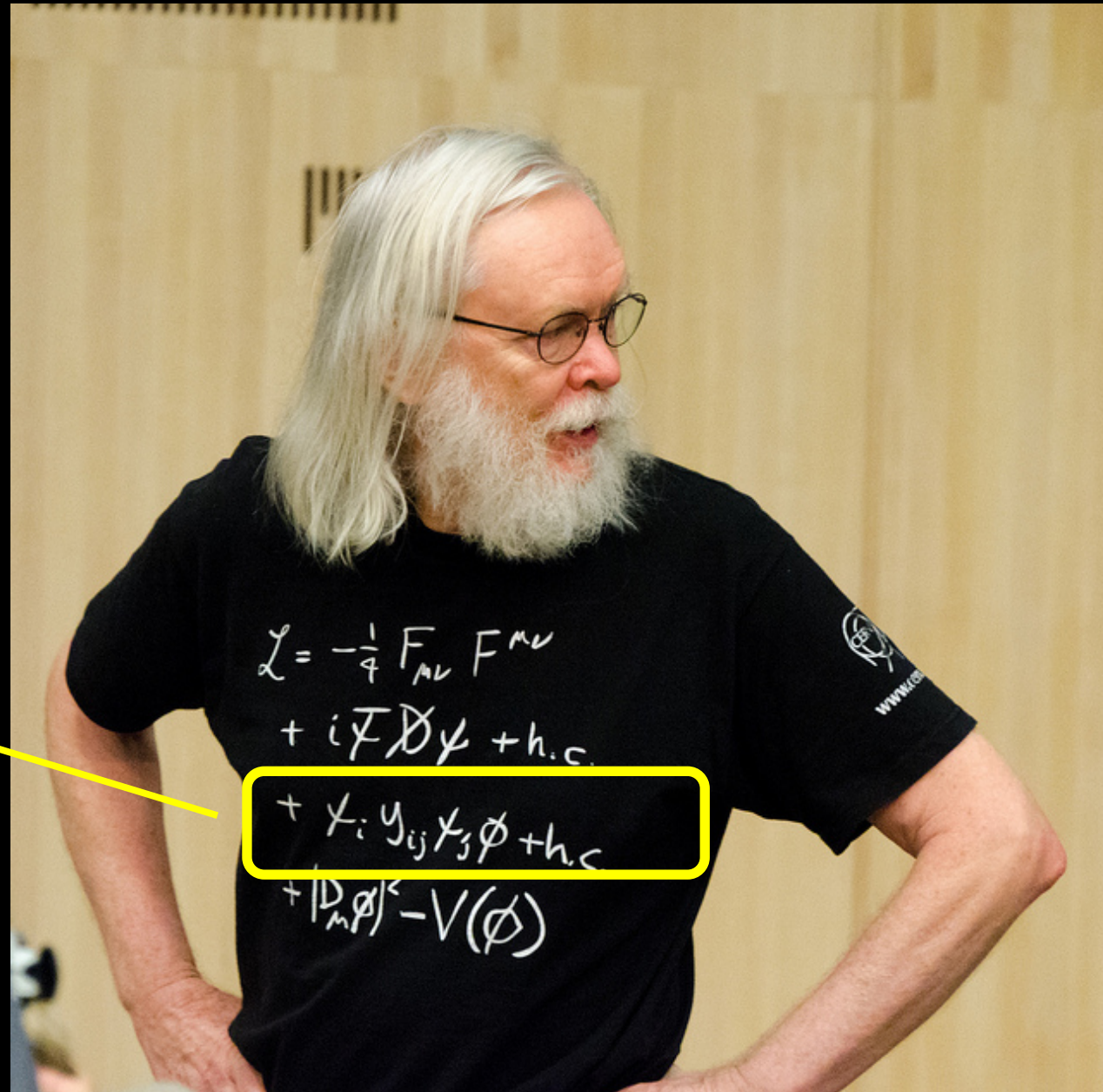
# The Standard Model on a T-Shirt

Kinetic term for fermions



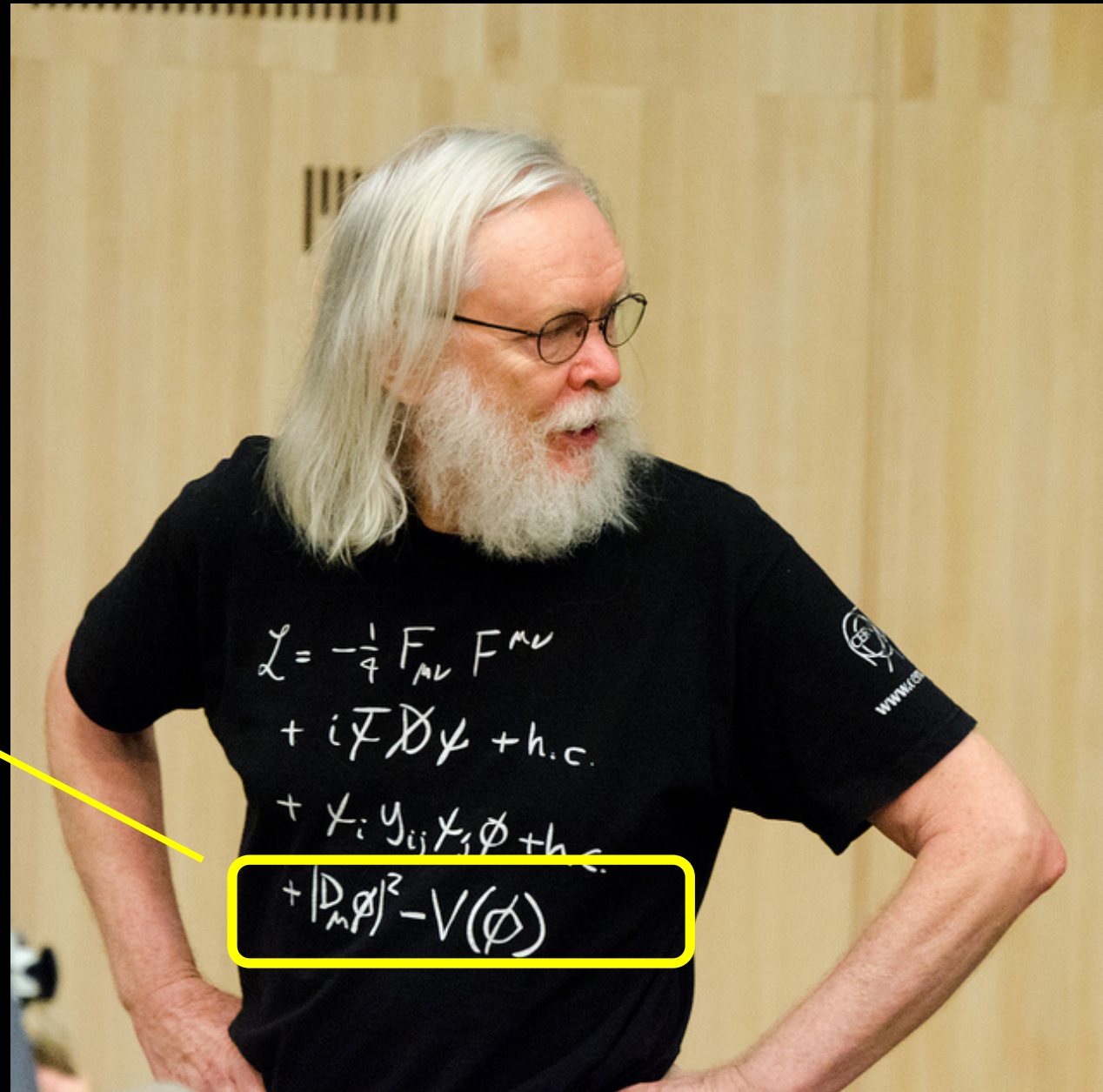
# The Standard Model on a T-Shirt

Interaction of the Higgs with fermions, fermion mass terms



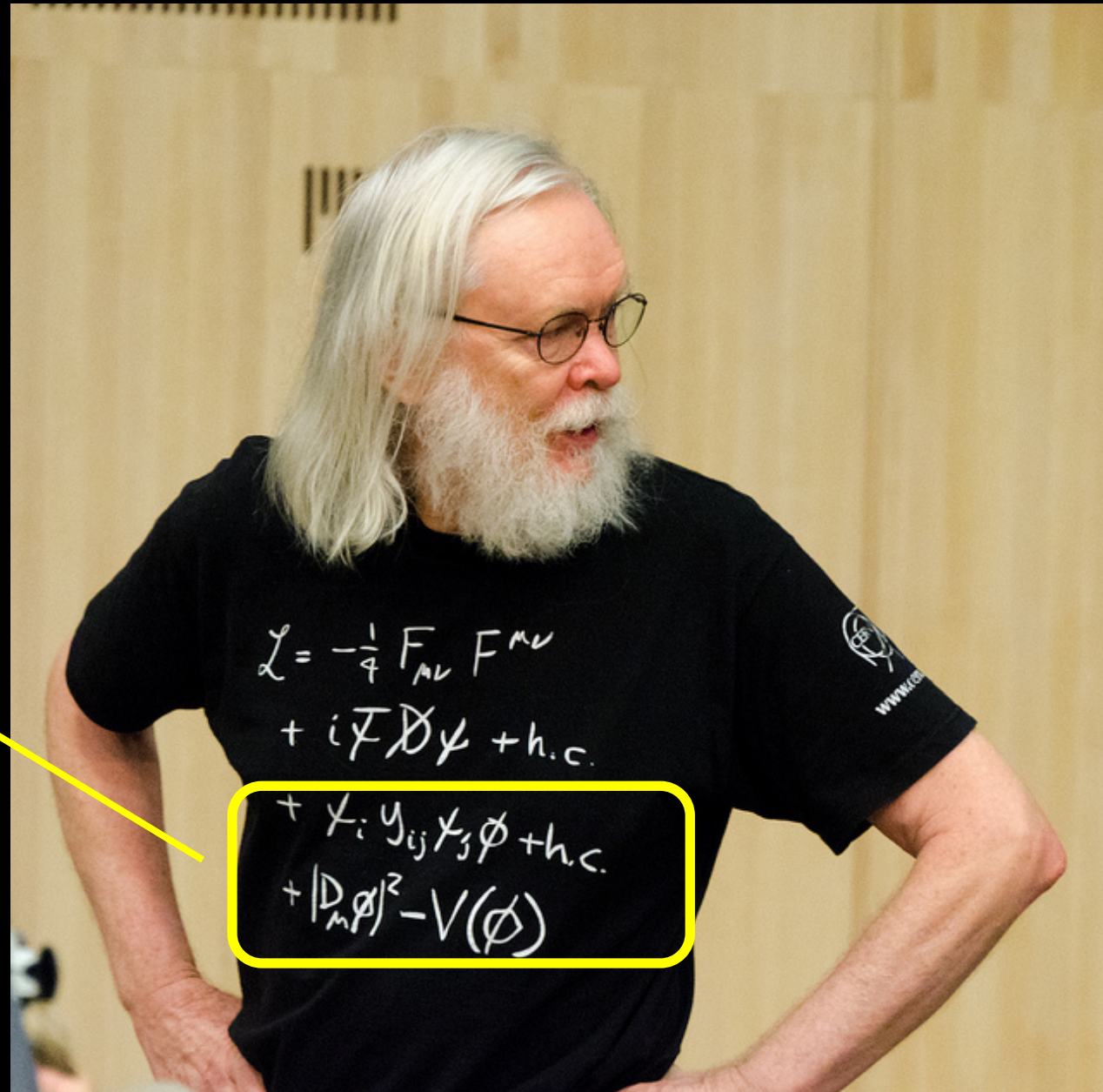
# The Standard Model on a T-Shirt

Interactions of the Higgs with weak gauge bosons (and itself), weak boson mass terms



# The Standard Model on a T-Shirt

Do these terms  
really describe  
Nature?

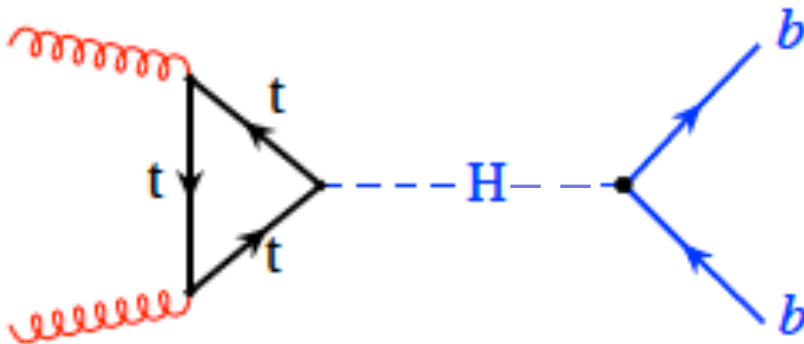
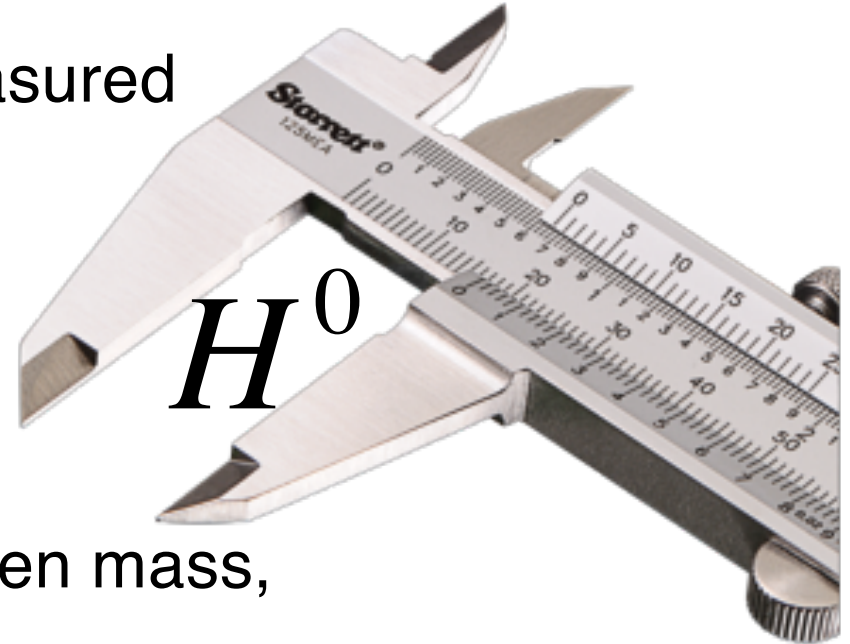


# Testing the Theory...

- How do we go about testing that there is such a thing as a Higgs field that interacts with massive particles?
  - We need to produce excitations of that field i.e. produce the Higgs bosons and measure how they interact with massive particles
- Producing Higgs bosons and demonstrating that what you observe is indeed a Higgs boson is very difficult...
  - The Higgs boson is very heavy, need a lot of energy
  - The Higgs boson is produced very rarely. And, only a small fraction of those produced can be identified as likely Higgs candidates: need many, many collisions
    - Roughly 1 Higgs in every 10,000,000,000 collisions
    - And we typically look for rare decays of the Higgs...

# PROPERTIES OF SM HIGGS BOSON

- Mass: not predicted must be measured
  - Electric charge: 0 (easy!)
  - Spin: 0
  - Coupling to SM particles: at a given mass, predicted for all particles
- Can be extracted from production and decay rates



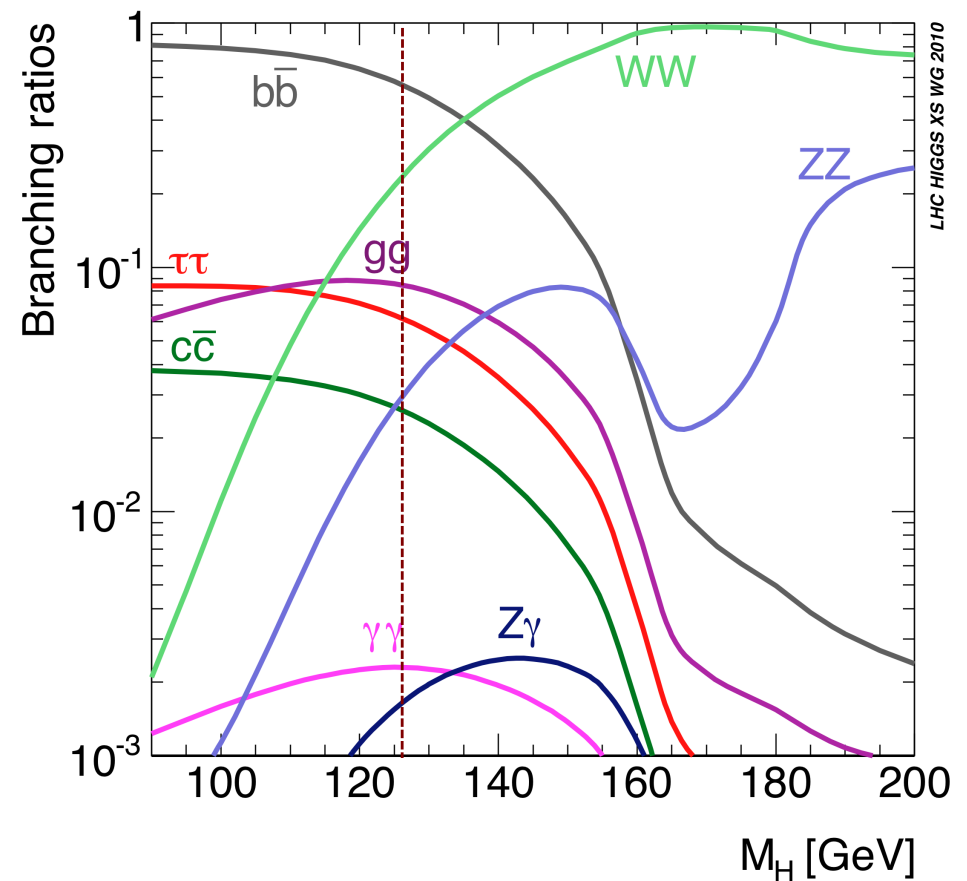
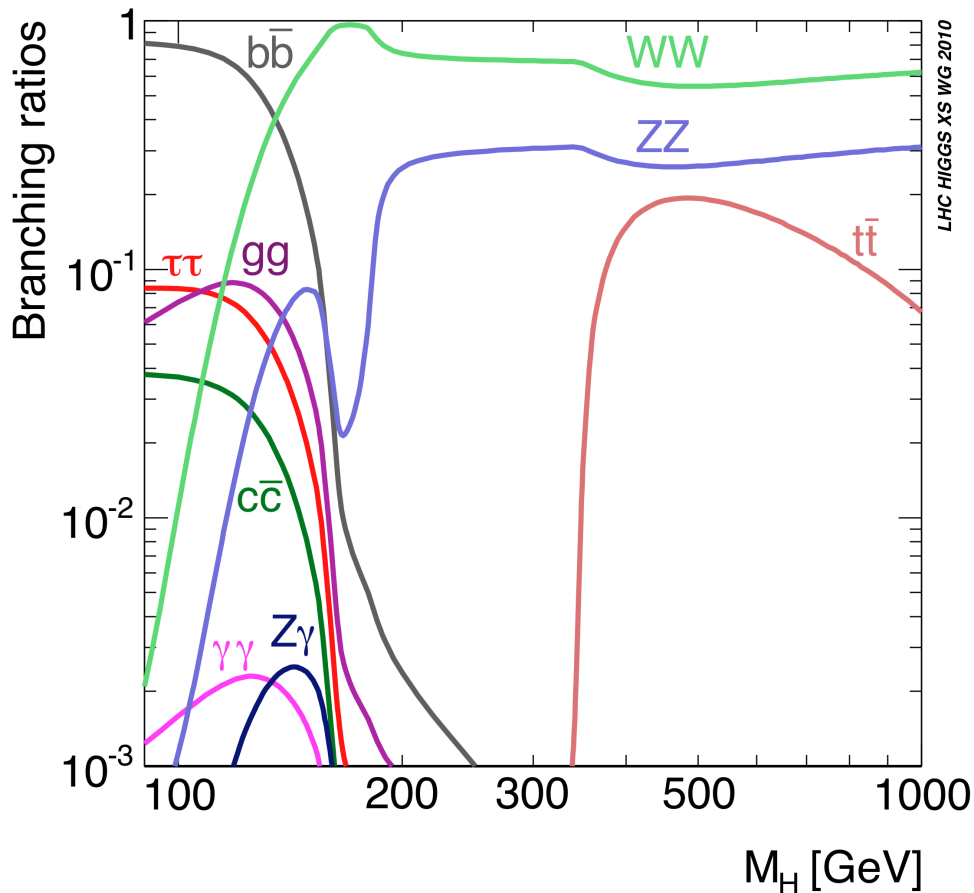
The observed rate of the process on the left depend on the Higgs's coupling to both top quarks and bottom quarks



# Higgs Decays

- Standard Model is a **very predictive theory** with respect to the Higgs boson once we measure the Higgs mass

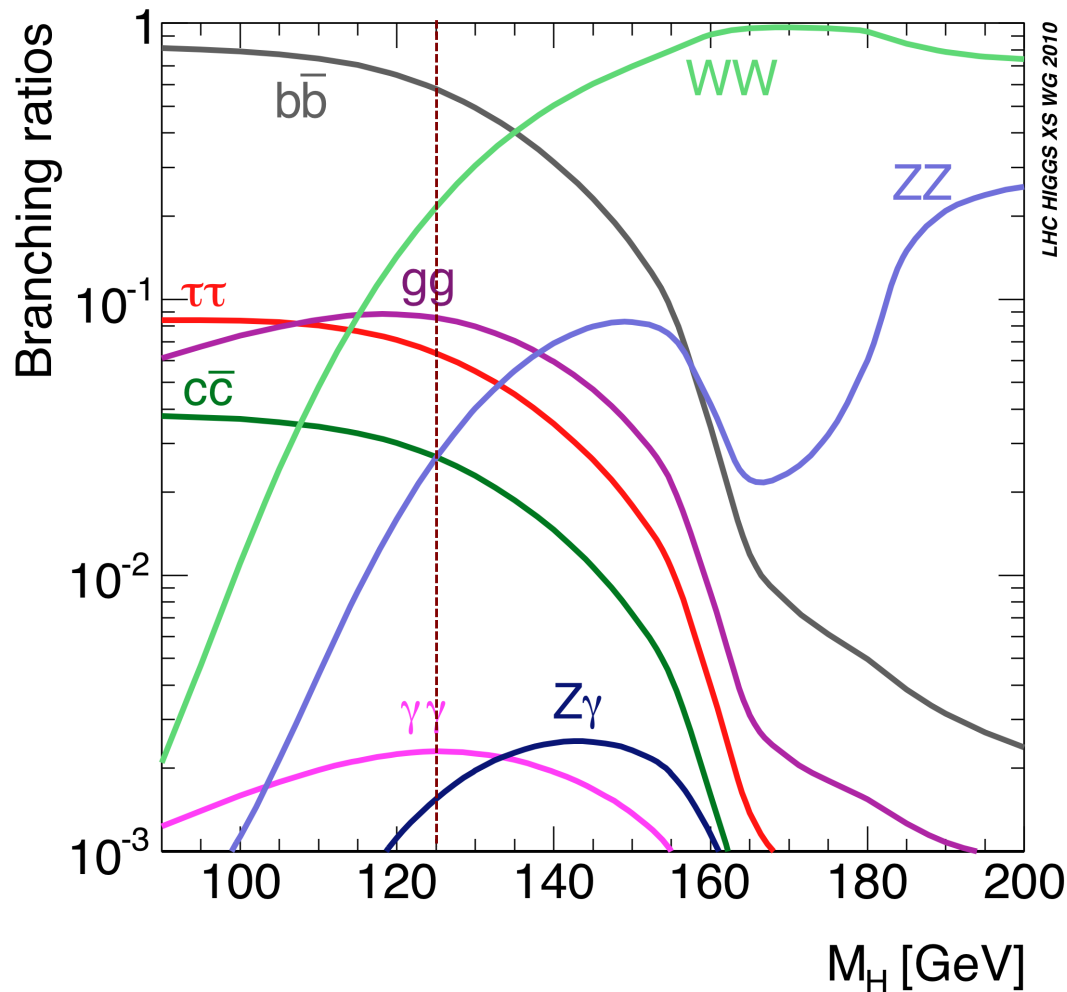
$$M_H^2 = 2v^2 \lambda$$



# HIGGS DECAYS

- At  $m_H = 125$  GeV, many decay channels can be studied

SM Decay Modes  
( $M_H = 125.1$  GeV)



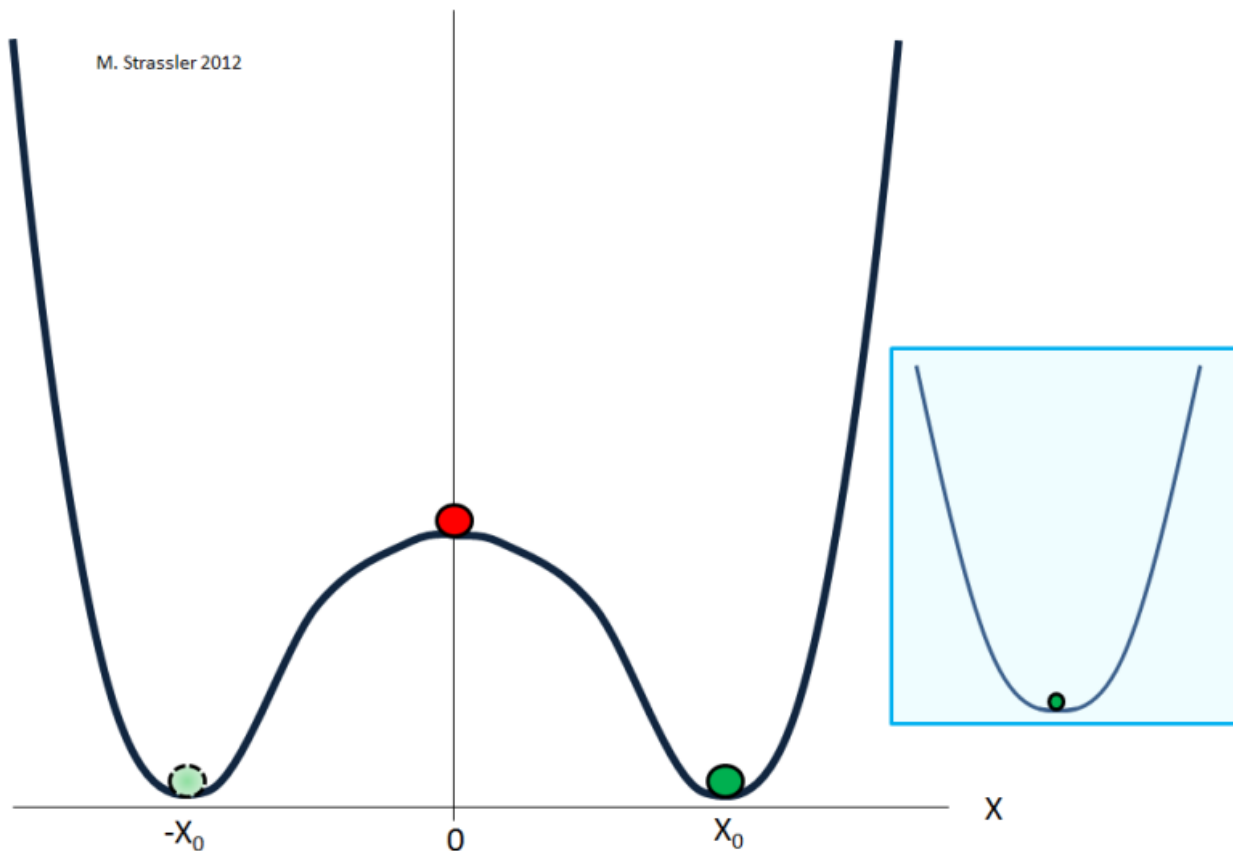
Decay Process	Decay Fraction
$bb$	0.58
$WW$	0.22
$\tau\tau$	0.06
$ZZ$	0.027
$\gamma\gamma$	0.0023
$Z\gamma$	0.0016
$\mu\mu$	0.0002

# Higgs Decays

- Standard Model is a **very predictive theory** with respect to the Higgs boson: the only unknown parameter is the Higgs mass

$$M_H^2 = 2v^2 \lambda$$

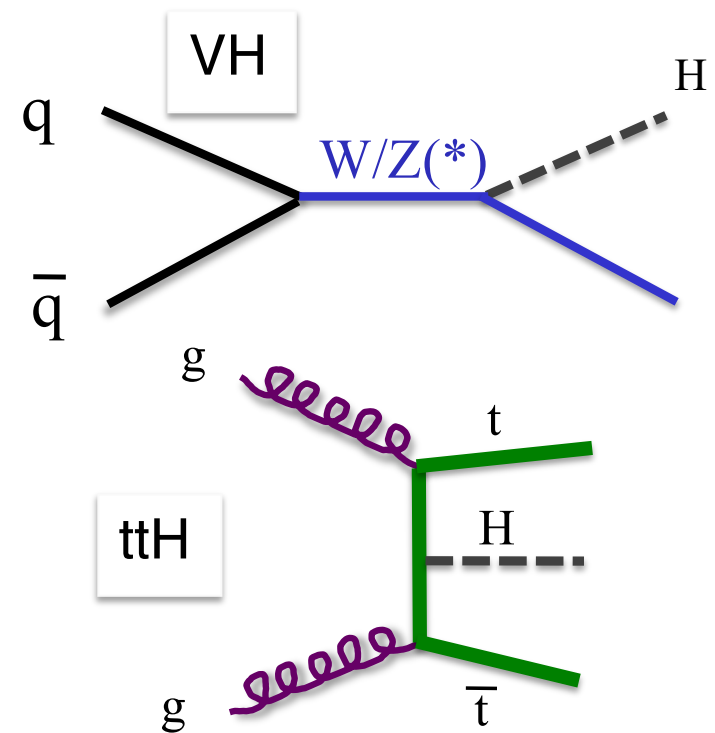
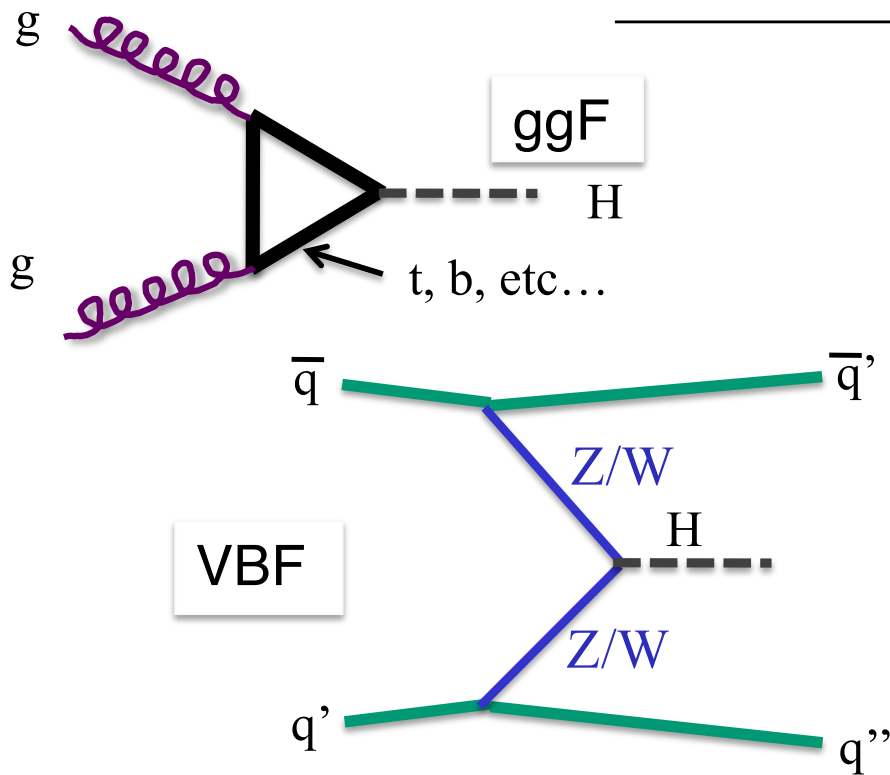
$$V(\Phi) = m^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$



# Higgs Production

Cross sections for  
 $m_H=125$  GeV:

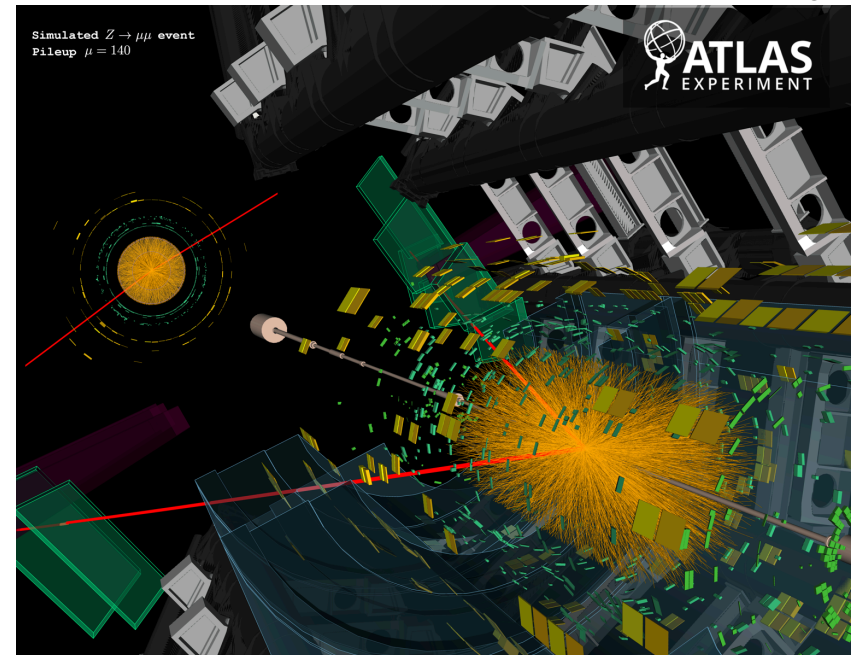
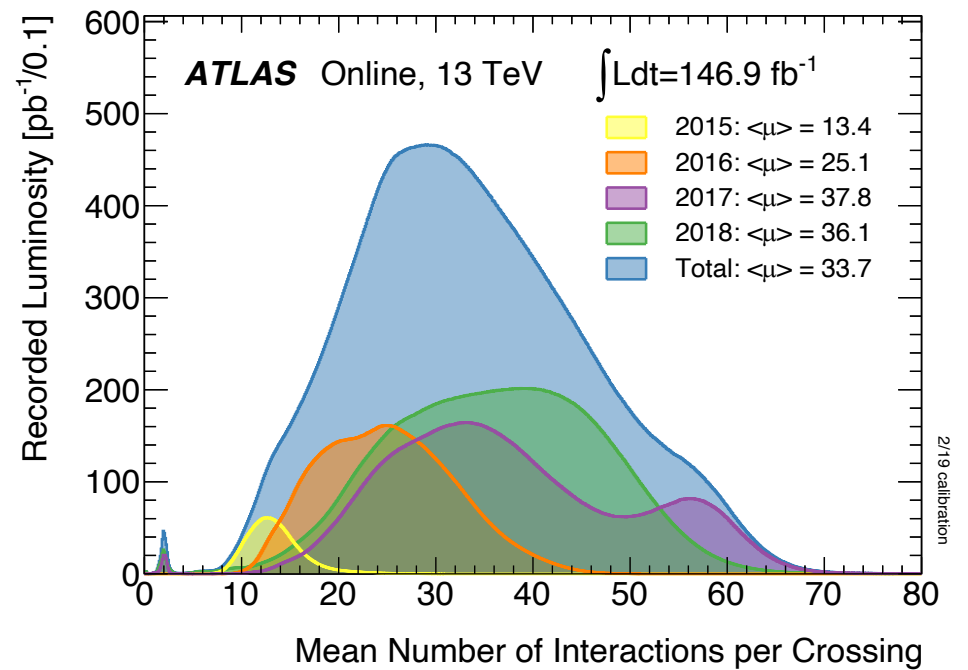
	process	8 TeV	13 TeV
<b>ggF</b>	gluon-gluon fusion	19 pb	44 pb
<b>VBF</b>	vector-boson fusion	1.6 pb	3.7 pb
<b>VH</b>	associated production	1.1 pb	2.2 pb
<b>ttH</b>	associated production	0.13 pb	0.51 pb



# LHC Datasets

## Proton-proton datasets:

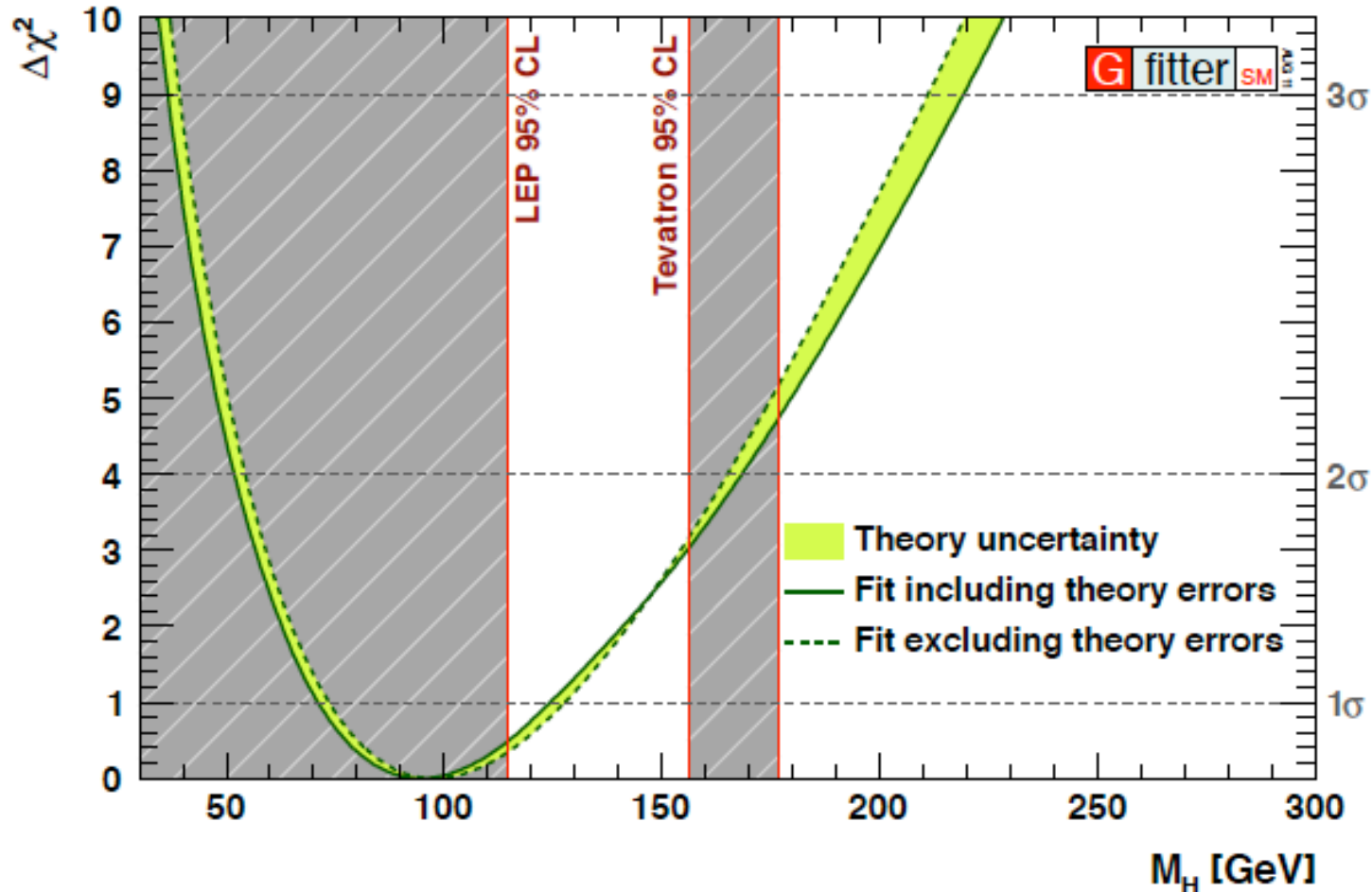
- Run 1  $\sim 25 \text{ fb}^{-1}$  at 7 and 8 TeV
- Run 2  $\sim 140 \text{ fb}^{-1}$  at 13 TeV
  - Large, high quality dataset yielding very impressive results
- Run 3 expected to yield  $\sim 170\text{-}190 \text{ fb}^{-1}$  with CM energy to be determined
  - Next 8 years: Run 2 dataset will be combined with Run 3 data for most measurements
- HL-LHC expected to yield  $3000 \text{ fb}^{-1}$  at 14 TeV: a  $\sim 20$ -fold increase over current dataset
  - A Higgs factory that will produce close to 200M Higgs bosons
  - Challenging collision environment: up to 200 interactions per bunch crossing



# The Past

# Before LHC: where to expect the Higgs?

- Fits to Standard Model data favoured a “light” Higgs Boson
- After 2010, at 95% CL, a 40 GeV window was left for the SM Higgs

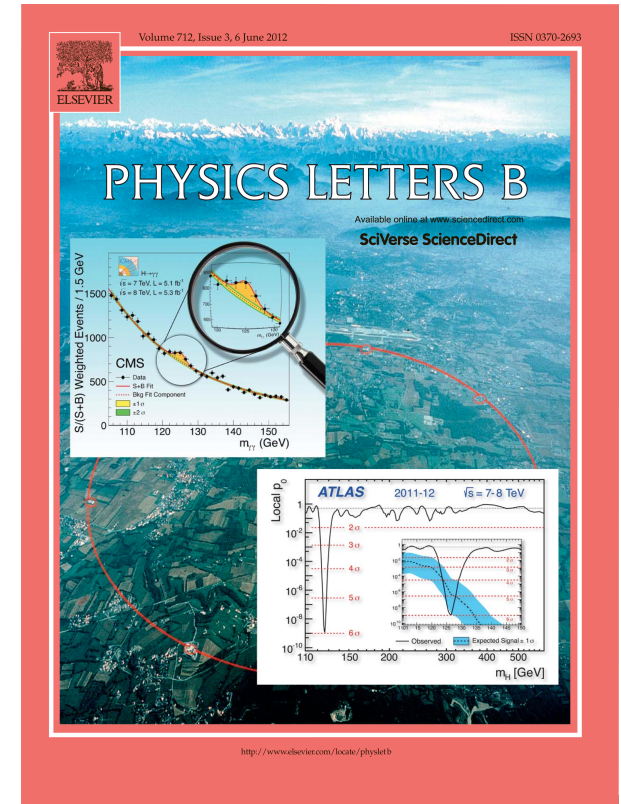


# DISCOVERY

- In April of 2012, the LHC started to produce collisions at a centre of mass of 8 TeV. The 7 TeV run showed some hints of a signal consistent with the SM Higgs boson
- On July 4<sup>th</sup> of 2012, both ATLAS and CMS announced that they had discovered a new particle consistent with what was expected in the SM. Two papers were then submitted at the end of July

- The ATLAS paper was titled:

“Observation of a New Particle in the Search for the Standard Model Higgs Boson”





# RUN 1, AFTER DISCOVERY

- The LHC delivered 4 times the data used for the discovery by the end of 2012. Results using the full dataset presented in March 2013 prompted CERN to declare that the new particle was “a” Higgs boson
- This was followed by two key ATLAS papers on the couplings and the spin of the discovered particle

Physics Letters B 726 (2013) 88–119

Contents lists available at [ScienceDirect](#)



Physics Letters B

[www.elsevier.com/locate/physletb](http://www.elsevier.com/locate/physletb)



Measurements of Higgs boson production and couplings in diboson final states with the ATLAS detector at the LHC<sup>☆</sup>



ATLAS Collaboration<sup>\*</sup>

Physics Letters B 726 (2013) 120–144

Contents lists available at [ScienceDirect](#)



Physics Letters B

[www.elsevier.com/locate/physletb](http://www.elsevier.com/locate/physletb)



Evidence for the spin-0 nature of the Higgs boson using ATLAS data<sup>☆</sup>

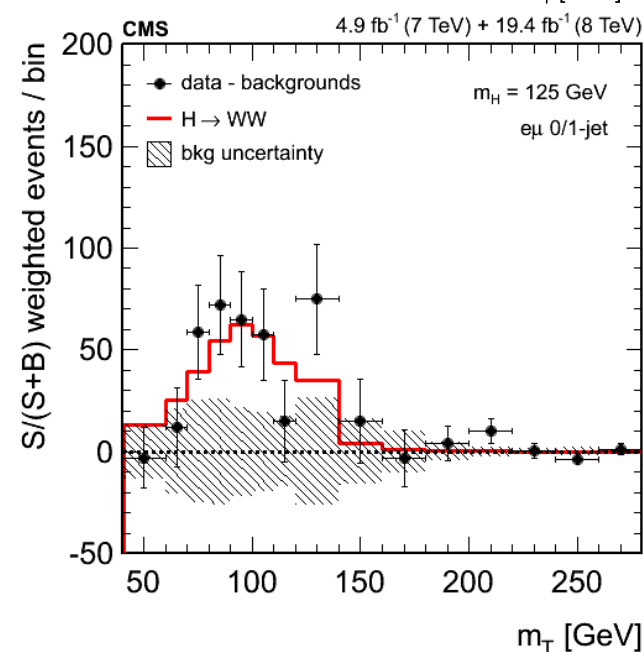
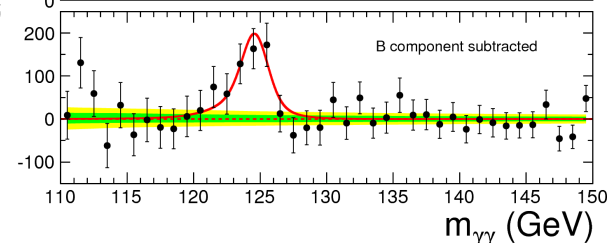
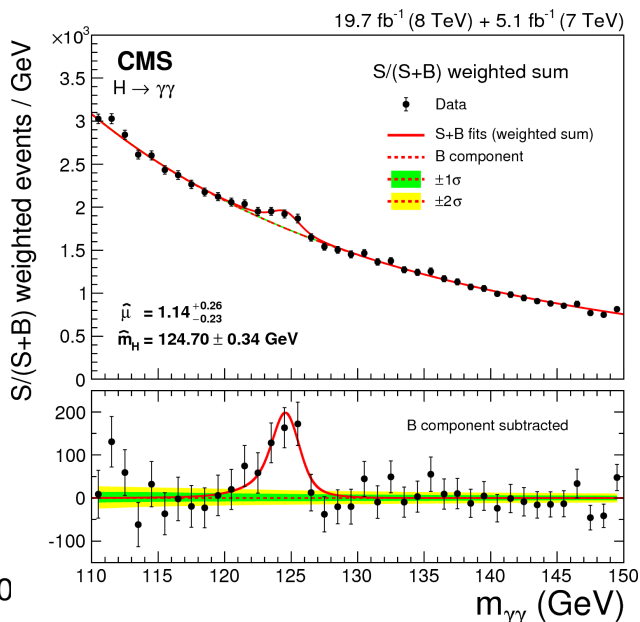
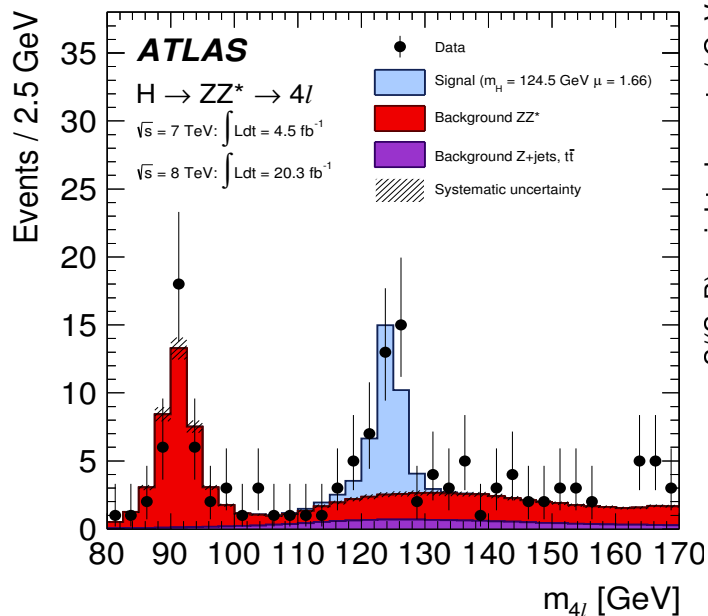
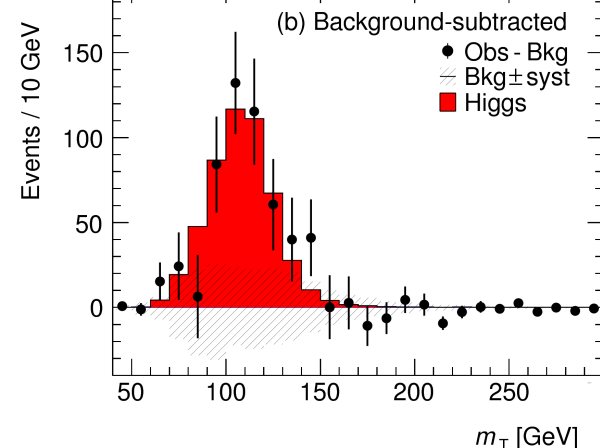
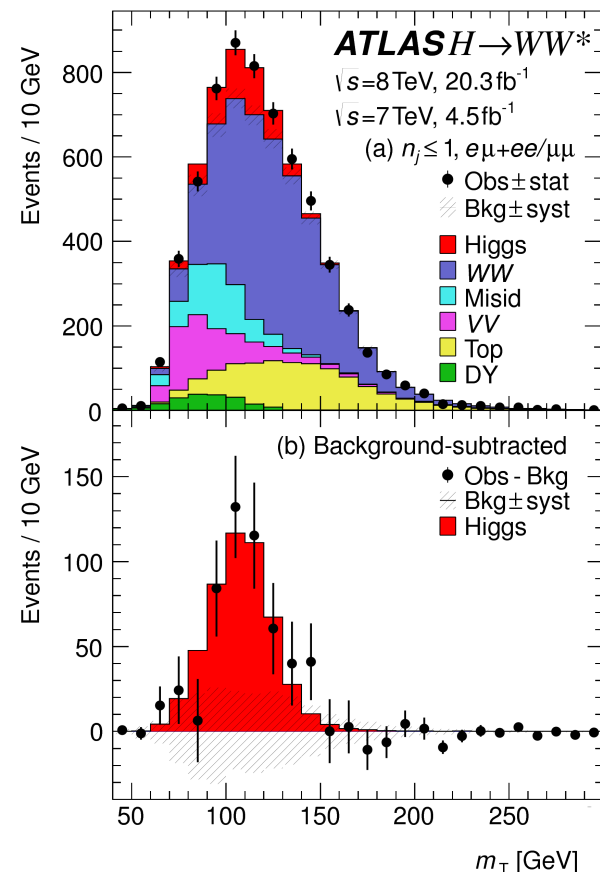
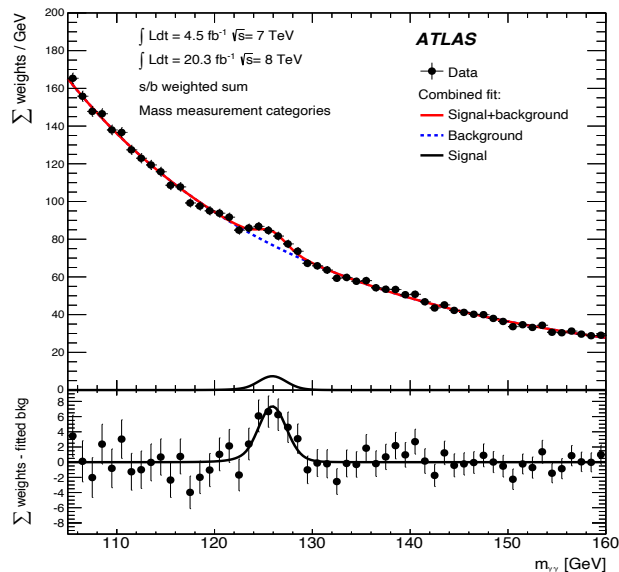
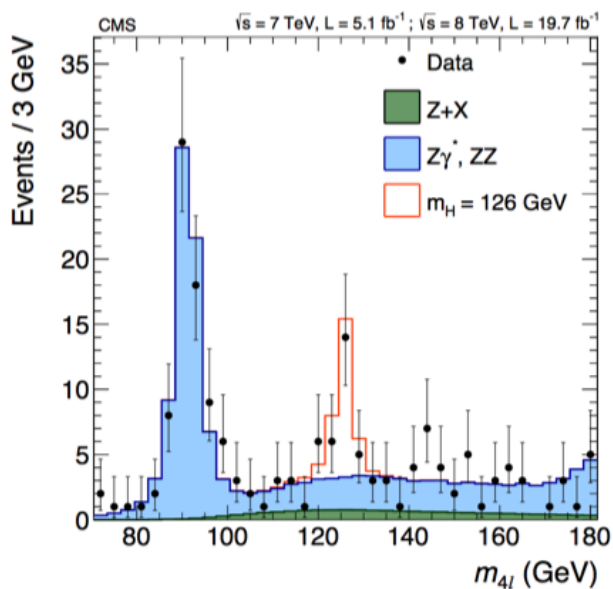


ATLAS Collaboration<sup>\*</sup>



2013 Nobel Prize in physics

# RUN 1: 5 $\sigma$ OBSERVATION IN DECAYS TO BOSONS



# HIGGS BOSON IN PDG

Citation: J. Beringer et al. (Particle Data Group), PR D86, 010001 (2012) and 2013 partial update for the 2014 edition (URL: <http://pdg.lbl.gov>)

## GAUGE AND HIGGS BOSONS

$\gamma$

$$I(J^{PC}) = 0,1(1^{--})$$

Mass  $m < 1 \times 10^{-18}$  eV  
 Charge  $q < 1 \times 10^{-35}$  e  
 Mean life  $\tau = \text{Stable}$

$g$   
or gluon

$$I(J^P) = 0(1^-)$$

Mass  $m = 0$  [a]  
 SU(3) color octet

Z

$$J = 1$$

Charge = 0  
 Mass  $m = 91.1876 \pm 0.0021$  GeV [d]  
 Full width  $\Gamma = 2.4952 \pm 0.0023$  GeV  
 $\Gamma(\ell^+ \ell^-) = 83.984 \pm 0.086$  MeV [b]  
 $\Gamma(\text{invisible}) = 499.0 \pm 1.5$  MeV [e]  
 $\Gamma(\text{hadrons}) = 1744.4 \pm 2.0$  MeV  
 $\Gamma(\mu^+ \mu^-) / \Gamma(e^+ e^-) = 1.0009 \pm 0.0028$   
 $\Gamma(\tau^+ \tau^-) / \Gamma(e^+ e^-) = 1.0019 \pm 0.0032$  [f]

W

$$J = 1$$

Charge =  $\pm 1$  e  
 Mass  $m = 80.385 \pm 0.015$  GeV  
 $m_Z - m_W = 10.4 \pm 1.6$  GeV  
 $m_{W^+} - m_{W^-} = -0.2 \pm 0.6$  GeV  
 Full width  $\Gamma = 2.085 \pm 0.042$  GeV  
 $\langle N_{\pi^\pm} \rangle = 15.70 \pm 0.35$   
 $\langle N_{K^\pm} \rangle = 2.20 \pm 0.19$   
 $\langle N_p \rangle = 0.92 \pm 0.14$   
 $\langle N_{\text{charged}} \rangle = 19.39 \pm 0.08$

# H<sup>0</sup>

## Higgs Bosons — H<sup>0</sup> and H<sup>±</sup>

H<sup>0</sup> Mass  $m = 125.9 \pm 0.4$  GeV

H<sup>0</sup> signal strengths in different channels [n]

Combined Final States =  $1.07 \pm 0.26$  (S = 1.4)

WW\* Final State =  $0.88 \pm 0.33$  (S = 1.1)

ZZ\* Final State =  $0.89^{+0.30}_{-0.25}$

$\gamma\gamma$  Final State =  $1.65 \pm 0.33$

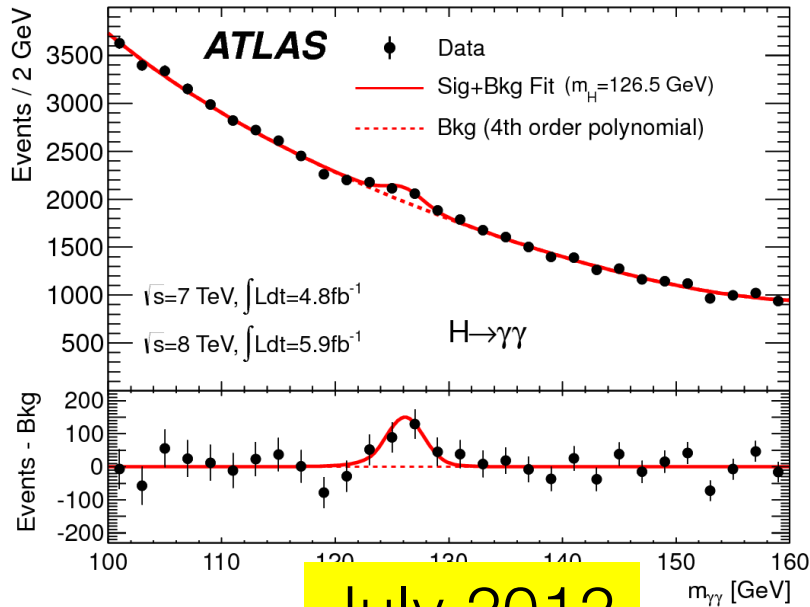
b $\bar{b}$  Final State =  $0.5^{+0.8}_{-0.7}$

$\tau^+ \tau^-$  Final State =  $0.1 \pm 0.7$

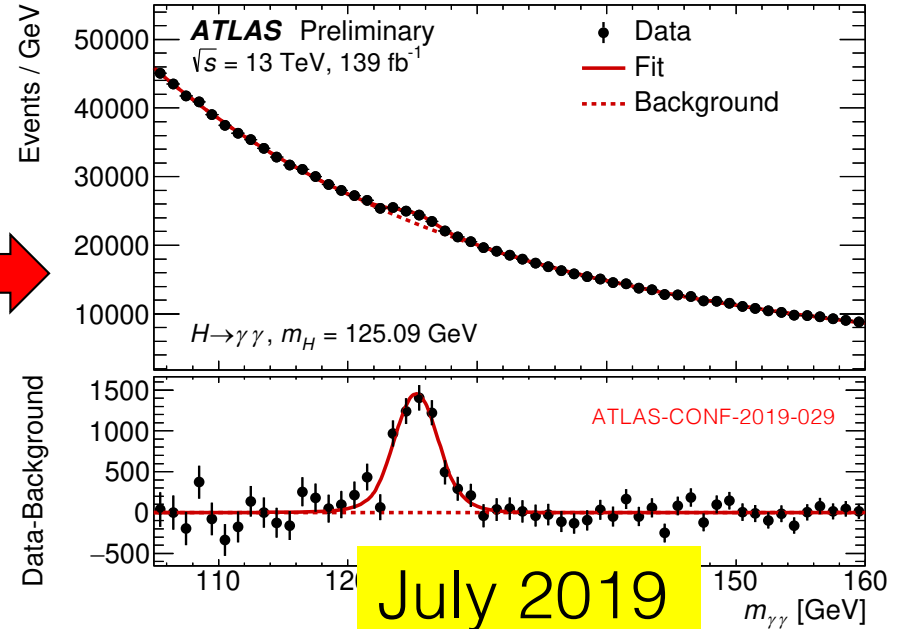
# The Present

# The Higgs Boson: then and now

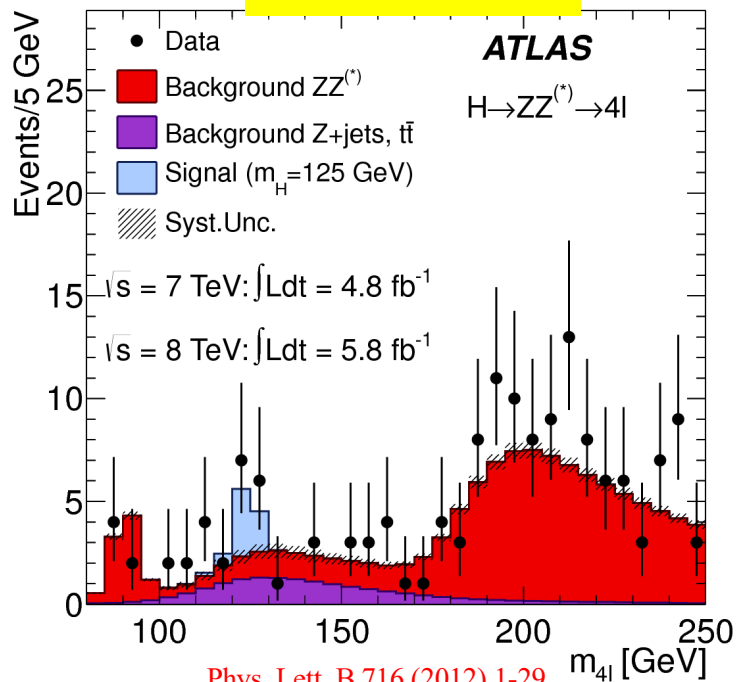
Full Run-2



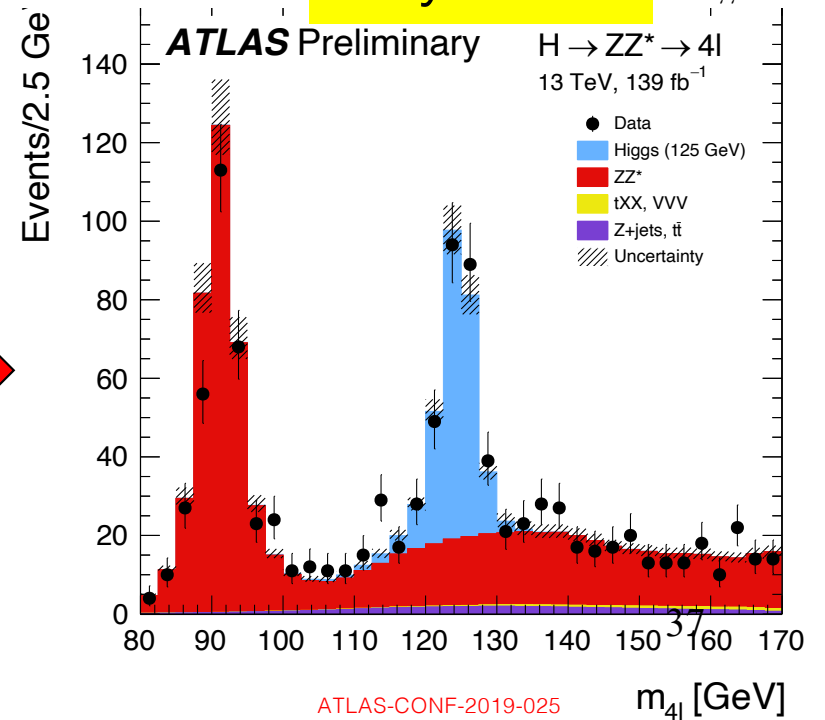
July 2012



July 2019

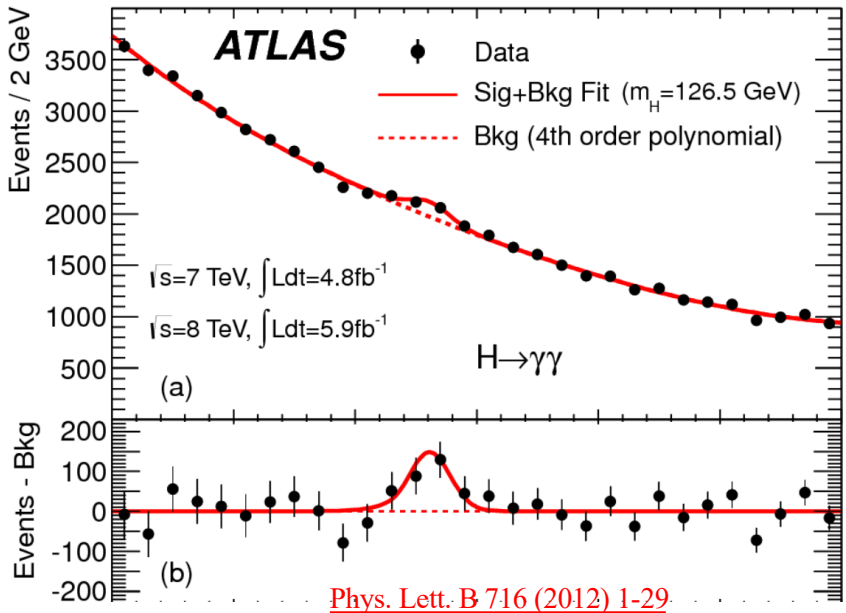


Phys. Lett. B 716 (2012) 1-29

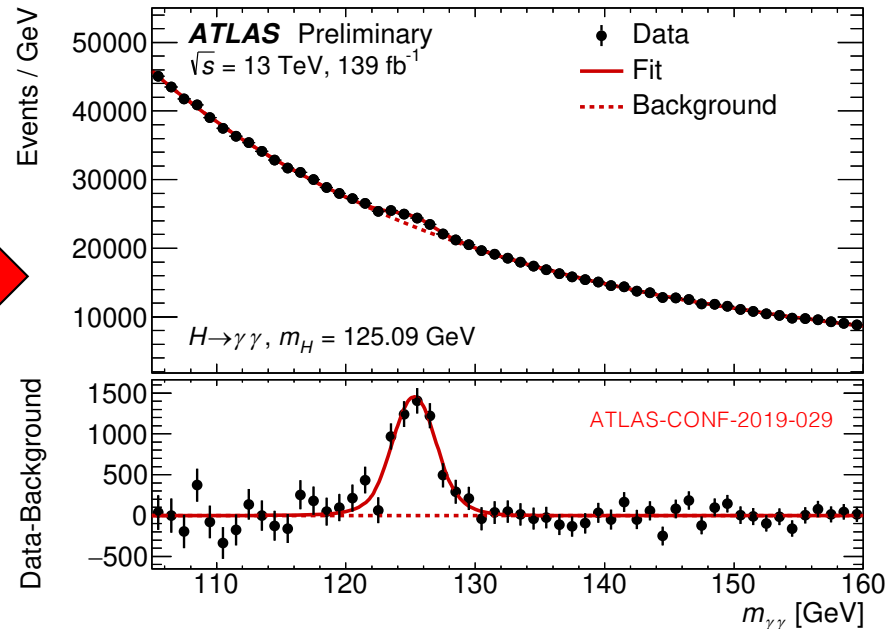


# The Higgs Boson: then and now

Full Run-2

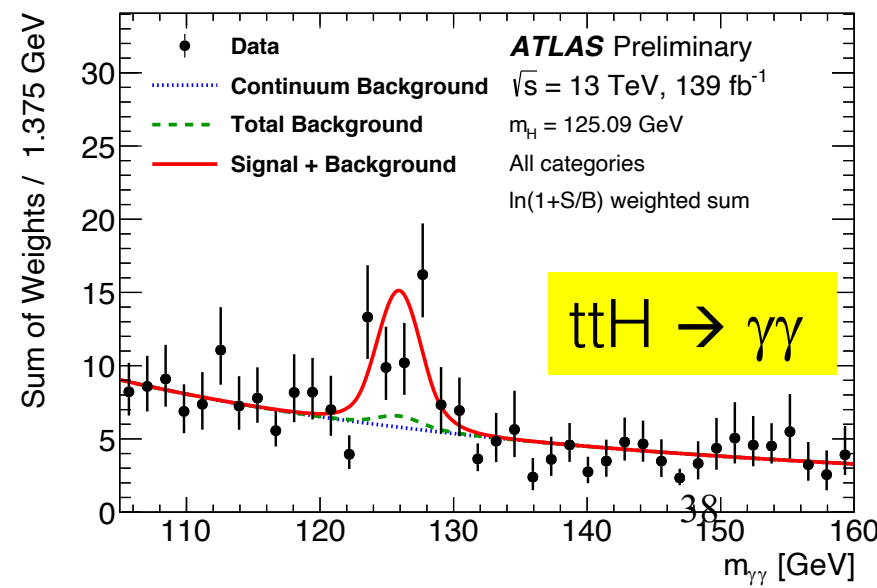


July 2012



July 2019

Observation of  $H \rightarrow \gamma\gamma$  in rare  $t\bar{t}H$  production channel in 2019  
 ATLAS-CONF-2019-004



$t\bar{t}H \rightarrow \gamma\gamma$

# Main Production and Decays

Analyses performed by ATLAS targeting specific production and decay modes:

	WW	ZZ	$\gamma\gamma$	bb	$\tau\tau$	$\mu\mu$
ggH	X	X	X		X	X
VBF	X	X	X	X	X	X
WH	X	X	X	X	X	X
ZH	X	X	X	X	X	X
ttH	X	X	X	X	X	X

**As of 2020, all the production and decay modes above except  $\mu\mu$  have been observed with a significance above 5 standard deviations**

**The observed rates consistent with Standard Model predictions**

# Higgs Boson Physics: where we are now

- Large sample of  $\sim 8\text{M}$  Higgs bosons (per experiment) produced in Run 2 allows for precision tests of the Higgs sector of the SM:

Channel	Produced	Selected	Mass resolution
$H \rightarrow \gamma\gamma$	18,200	6,440	1–2%
$H \rightarrow ZZ^*$	210,000	( $\rightarrow 4\ell$ ) 210	1–2%
$H \rightarrow WW^*$	1,680,000	( $\rightarrow 2\ell 2\nu$ ) 5,880	20%
$H \rightarrow \tau\tau$	490,000	2,380	15%
$H \rightarrow bb$	4,480,000	9,240	10%

## Major progress in last few years:

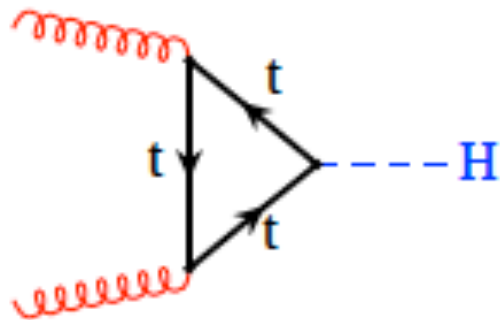
- Observation of  $H \rightarrow bb$  decay
- Observation of  $ttH$  production
- Evidence of  $H \rightarrow \mu\mu$  decay

## At the end of Run 2:

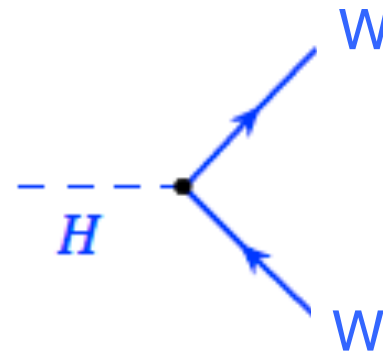
- Mass measurement precision  $\sim 0.1\%$
- All major production and decay modes have been observed. Other targets for future runs:
  - $Z\gamma$ ,  $\mu\mu$  decay modes
  - $tH$  production mode (if SM, very difficult...)



# COUPLINGS FRAMEWORK



$\mu_p^i$



$\mu_{BR}^j$

- Signal strength “ $\mu$ ” reflects the ratio of observed rates to predicted rates
- “ $\kappa$  framework”: use all measurements of signal strength parameters ( $\mu_p^i$ ,  $\mu_{BR}^j$ ) to extract modifiers to the SM couplings:

- Decay:  $\Gamma_i = \kappa_i^2 \Gamma_i^{SM}$

- Production:  $\sigma_i = \kappa_i^2 \sigma_i^{SM}$

- Width:  $\Gamma_H = \sum_i \kappa_i^2 \Gamma_i^{SM}$

## Assumptions:

- Only one Higgs
- SM production and decay kinematics
  - Tensor structure is that of SM
  - 0+ scalar
- Narrow resonance

# Where we are now, before start of Run 3

The 125 GeV Higgs boson has been shown to be very compatible with the Standard Model:

- Recent ATLAS and CMS combined signal strengths:

ATLAS-CONF-2019-037  $\mu = 1.06 \pm 0.07$

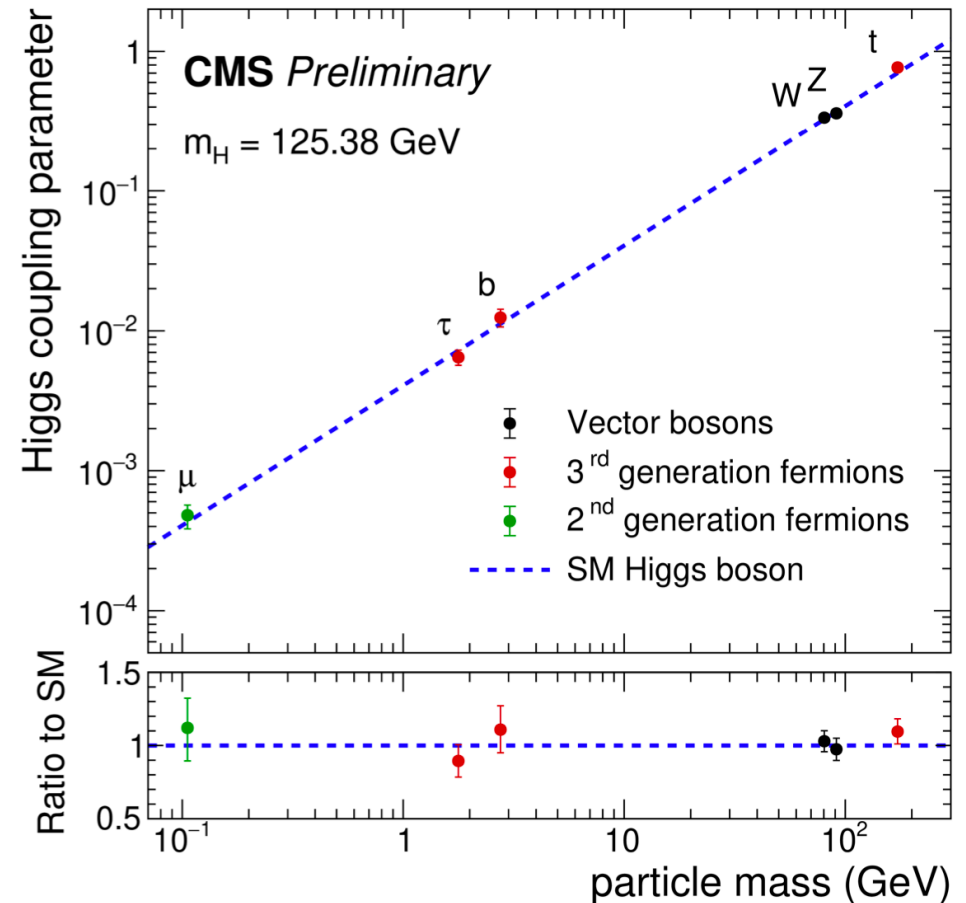
CMS-PAS-HIG-19-005  $\mu = 1.02^{+0.07}_{-0.06}$

- Yukawa couplings observed for top, bottom, and tau fermions with non-universal pattern predicted by the SM

Parameter	Result
$\kappa_Z$	$1.02 \pm 0.06$
$\kappa_W$	$1.05 \pm 0.06$
$\kappa_b$	$0.98^{+0.14}_{-0.13}$
$\kappa_t$	$0.96 \pm 0.08$
$\kappa_\tau$	$1.06^{+0.15}_{-0.14}$
$\kappa_\mu$	$1.12^{+0.26}_{-0.32}$

CMS-PAS-HIG-19-006

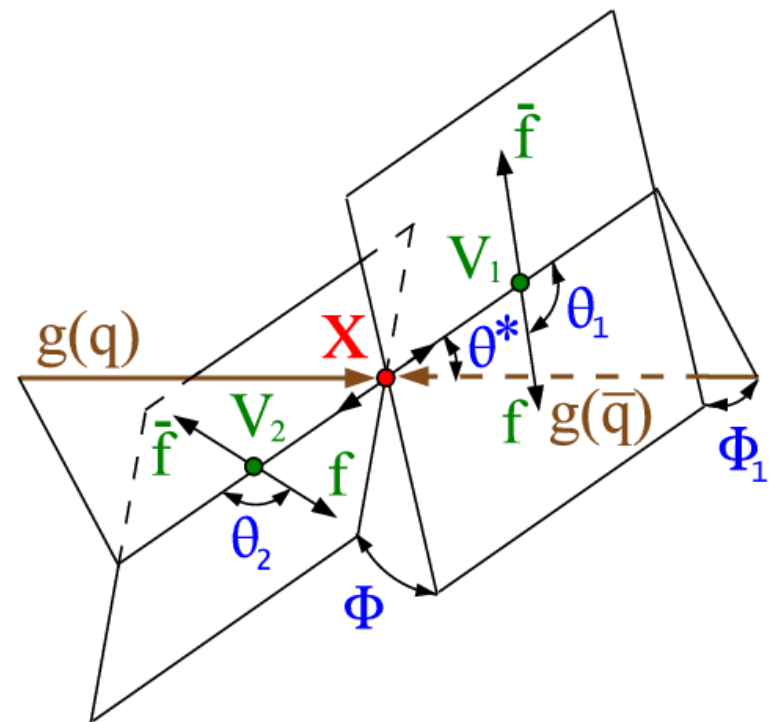
35.9-137 fb<sup>-1</sup> (13 TeV)



# Where we are now, before start of Run 3

- Spin 0 confirmed during Run 1 with pure CP-odd state easily excluded in  $H \rightarrow ZZ$
- No hints yet of CP-odd coupling component to SM fermions
  - analyses using  $t\bar{t}H$  production by ATLAS and CMS exclude pure CP-odd top coupling at more than  $3\sigma$  C.L.
  - Analysis excludes pure CP-odd  $\tau$  coupling at more than  $3\sigma$  C.L.

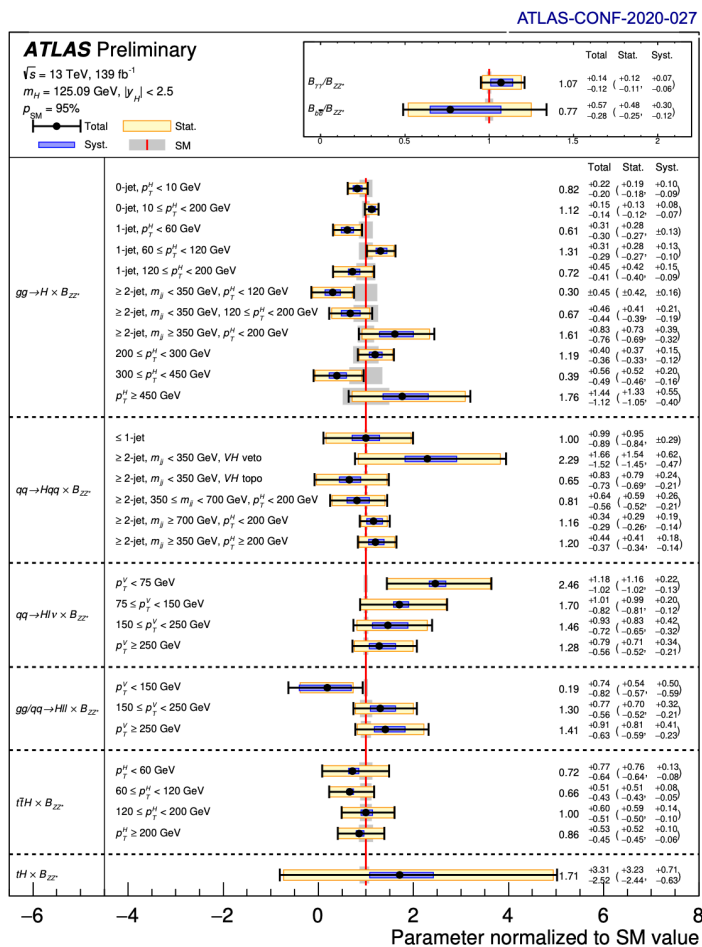
ZZ: full kinematic information available for spin/CP determination



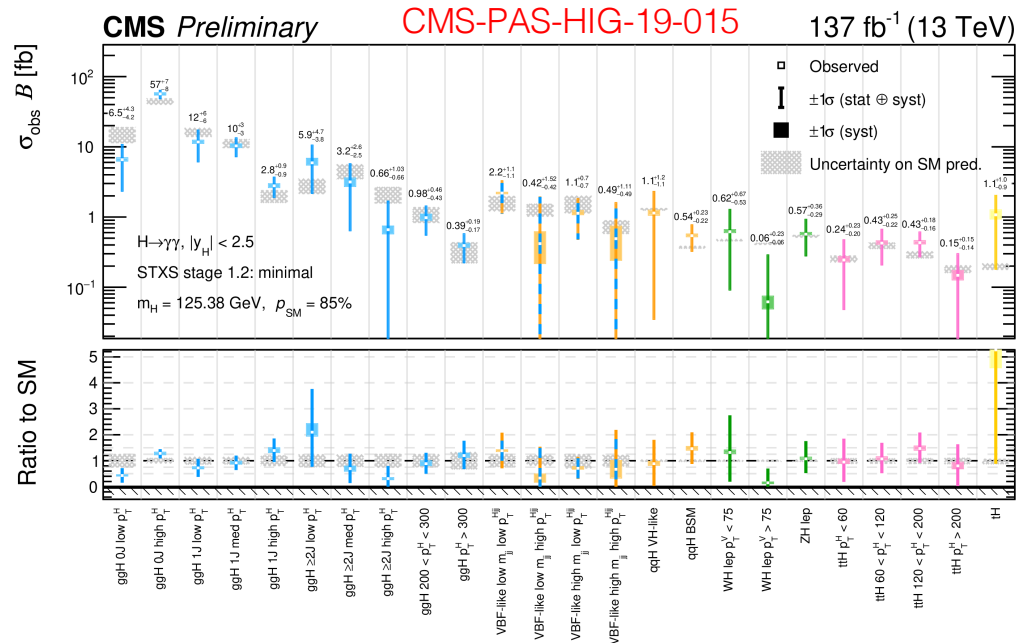
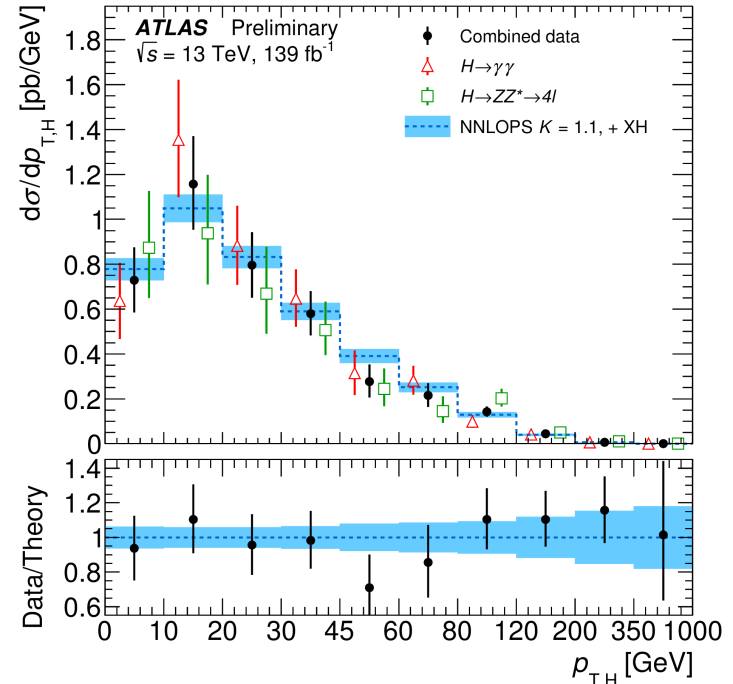
# Where we are now: comprehensive kinematic studies

After observation, vast programme of kinematic measurements was launched

- Differential cross sections
- Measurements by production mode in various kinematic regions (STXS)



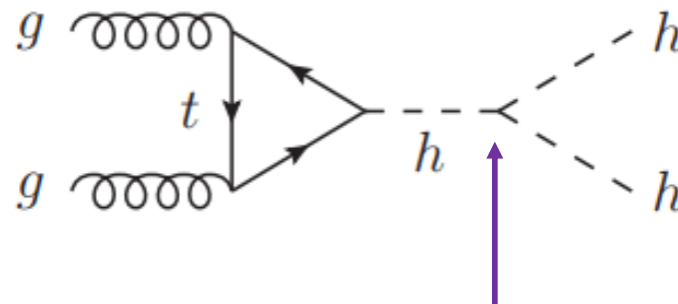
ATLAS-CONF-2020-026



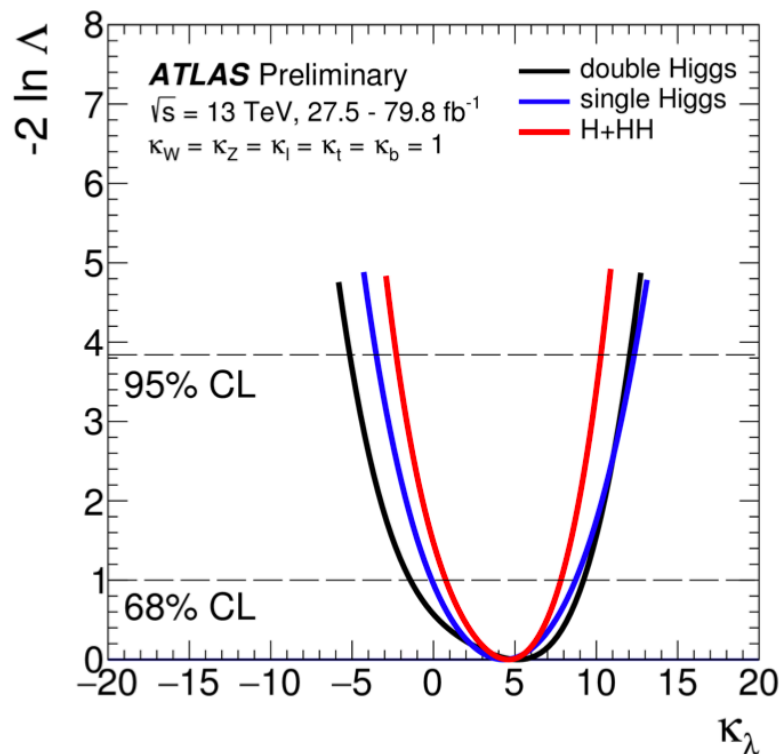
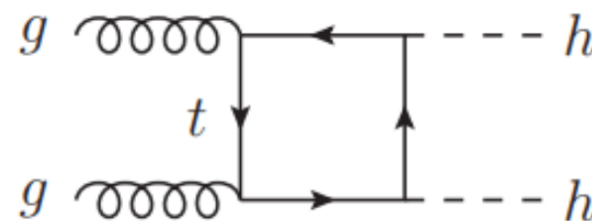
# Where we are now: towards the self-coupling

Making progress towards testing the shape of the Higgs potential through the Higgs self-coupling ( $\lambda_3$ )

Sensitivity to SM coupling will require HL-LHC but much progress can be made in coming years (more later...)



$$V(\Phi) = m^2\Phi^\dagger\Phi + \lambda(\Phi^\dagger\Phi)^2$$

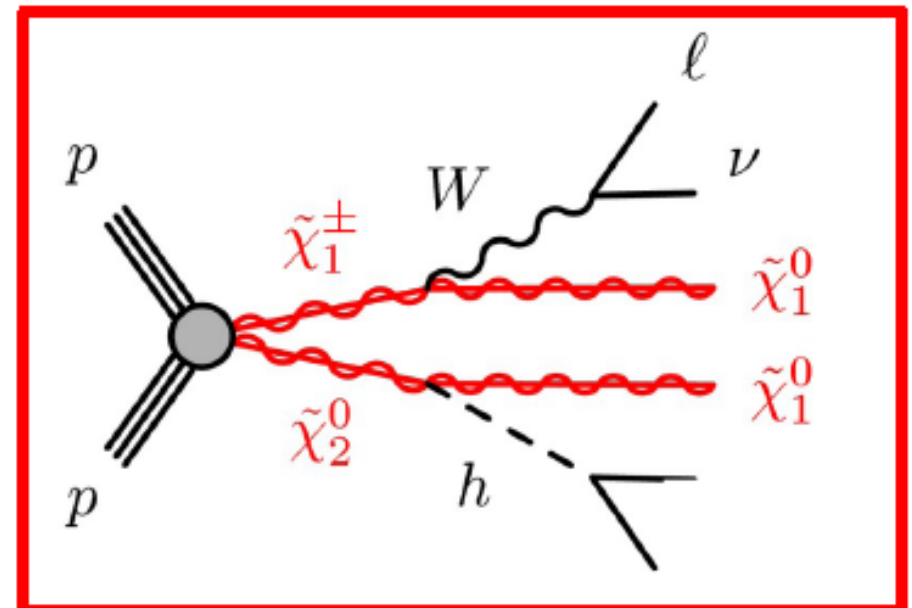
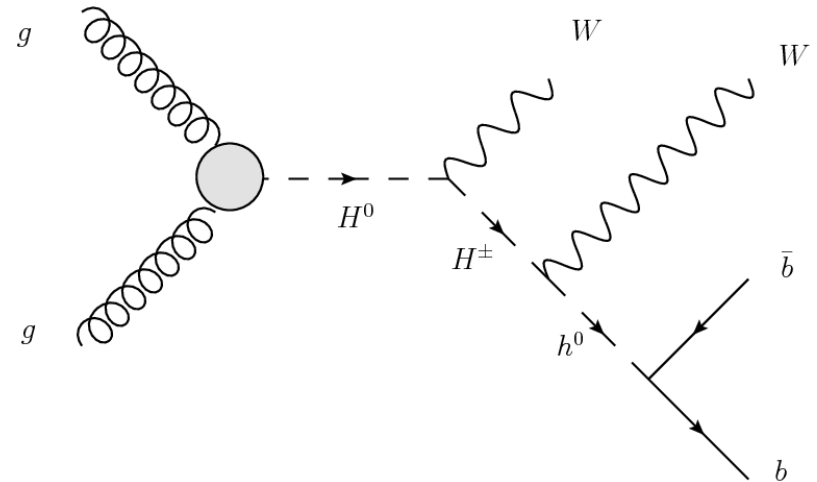


Model	$\kappa_\lambda^{+1\sigma}$ $\kappa_\lambda^{-1\sigma}$	$\kappa_\lambda$ [95% CL]	
$\kappa_\lambda$ -only	$4.6^{+3.2}_{-3.8}$	$[-2.3, 10.3]$	obs.
	$1.0^{+7.3}_{-3.8}$	$[-5.1, 11.2]$	exp.
Generic	$5.5^{+3.5}_{-5.2}$	$[-3.7, 11.5]$	obs.
	$1.0^{+7.6}_{-4.5}$	$[-6.2, 11.6]$	exp.

# BSM HIGGS PHYSICS

Different strategies to look for physics beyond the SM in the Higgs sector

- Search for other Higgs bosons
  - Extensive search program looking for new scalars in many final states
- Search for exotic decays of the 125 Higgs
  - Rich phenomenology that includes invisible decays, displaced decays
- Use the 125 GeV Higgs as a tool to find new physics
  - Higgs in decays
  - Use measurements to constrain BSM scenarios



# BSM HIGGS SEARCHES

A non-exhaustive list...

Neutral Heavy Higgs to Fermions	$H/A \rightarrow (b)\tau\tau$ (LL, LH, HH)
	$H/A \rightarrow (b)\mu\mu$
	$H/A \rightarrow (b)bb$
	$H/A \rightarrow tt$

Neutral Heavy Higgs to Bosons	$H \rightarrow \gamma\gamma$
	$H \rightarrow ZZ \rightarrow 4l$
	$H \rightarrow ZZ \rightarrow ll\nu\nu$
	$H \rightarrow ZZ \rightarrow llqq$
	$H \rightarrow ZZ \rightarrow \nu\nu qq$
	$H \rightarrow WW \rightarrow l\nu l\nu$
	$H \rightarrow WW \rightarrow l\nu qq$

Neutral Heavy Higgs to Bosons, including light Higgs	$(H \rightarrow) hh \rightarrow \gamma\gamma bb$
	$(H \rightarrow) hh \rightarrow 4b$
	$(H \rightarrow) hh \rightarrow bb\tau\tau$
	$(H \rightarrow) hh \rightarrow VV\gamma\gamma \rightarrow 4j\gamma\gamma$
	$(H \rightarrow) hh \rightarrow WW\gamma\gamma \rightarrow l\nu qq\gamma\gamma$
	$A \rightarrow Zh \rightarrow ll\tau\tau$ (LL, LH, HH)
	$A \rightarrow Zh \rightarrow (ll/\nu\nu)bb$

Heavy and light Charged Higgs	$H^\pm \rightarrow \tau\nu + jets$
	$H^\pm \rightarrow tb$ (resolved)
	$H^\pm \rightarrow tb$ s-chan (had, L+j)
	$H^\pm \rightarrow \tau\nu + lep(s)$
	$H^\pm \rightarrow \mu\nu$
	$H^\pm \rightarrow cs$
	$H^\pm \rightarrow cb$
	- AW
	$H^\pm \rightarrow Wh$ (WH, WA)
	$H^\pm \rightarrow W\gamma$
	$H^\pm \rightarrow tb$ (boosted)
	$H^\pm \rightarrow WZ \rightarrow tb$ ( $l\nu qq, qqll$ )
	$H^\pm$

LFV / FCNC / rare decays	$H \rightarrow \tau\mu, \tau e$
	$H \rightarrow e\mu$
	$H \rightarrow J/\psi\gamma, Y\gamma$
	$H \rightarrow ZJ/\psi, ZY$
	$H \rightarrow \phi\gamma$
	$t \rightarrow cH$ (various)

Exotics decays with MET, Dark-sector Inspired	mono H ( $\rightarrow \gamma\gamma + MET$ )
	mono H ( $\rightarrow bb + MET$ )
	mono H ( $\rightarrow 4l + MET$ )
	$H \rightarrow \gamma\gamma dark$
	$ZH \rightarrow (ll)INV$
	VBF $H \rightarrow INV$
	$VH \rightarrow (jj)INV$
$ttH \rightarrow INV$ (various)	
<b>ggF <math>H \rightarrow INV</math> (monojet).</b>	

Exotics decays with no MET, Dark-sector / NMSSM Inspired	$H \rightarrow ZdarkZ(dark) \rightarrow 4l$
	$h \rightarrow 2a \rightarrow \mu\mu\mu\mu$
	$h \rightarrow Za \rightarrow ll\mu\mu$
	$a \rightarrow \mu\mu$
	$h \rightarrow 2a \rightarrow 4\gamma$ (multiphoton)
	$h \rightarrow 2a \rightarrow bb\mu\mu$
	$h \rightarrow 2a \rightarrow bb\tau\tau$
	$(bb)a \rightarrow (bb)\tau\tau \rightarrow (bb)e\mu$
	$h \rightarrow 2a \rightarrow 4\tau$
$H^\pm \rightarrow aW$	

# BSM HIGGS SEARCHES

A non-exhaustive list... Many of the searches below were performed in Run 1

Neutral Heavy Higgs to Fermions	$H/A \rightarrow (b)\tau\tau$ (LL, LH, HH)		$H \rightarrow \tau\nu + \text{jets}$	Exotics	mono H ( $\rightarrow \gamma\gamma + \text{MET}$ )
	$H/A \rightarrow (b)\mu\mu$		$H \rightarrow tb$ (resolved)		mono H ( $\rightarrow bb + \text{MET}$ )
	$H/A \rightarrow (b)bb$		$H \rightarrow tb$ s-chan (had, L+j)		mono H ( $\rightarrow 4l + \text{MET}$ )
	$H/A \rightarrow tt$		$H \rightarrow \text{out}(e)$		$H \rightarrow \text{dark}$
Neutral Heavy Higgs to Bosons	$H \rightarrow \gamma\gamma$				$(ll)INV$
	$H \rightarrow ZZ$				$H \rightarrow INV$
	$H \rightarrow ZZ$				$(jj)INV$
	$H \rightarrow ZZ$				$INV$ (various)
	$H \rightarrow ZZ$				$H \rightarrow INV$ (monojet).
	$H \rightarrow WW$				$\text{dark}Z(\text{dark}) \rightarrow 4l$
	$H \rightarrow WW$				$a \rightarrow \mu\mu\mu\mu$
Neutral Heavy Higgs to Bosons, including light Higgs	$(H \rightarrow) hh \rightarrow \gamma\gamma bb$	LFV / FCNC / rare decays	$H^{++}$	Exotics decays with no MET, Dark-sector / NMSSM Inspired	$h \rightarrow Za \rightarrow ll\mu\mu$
	$(H \rightarrow) hh \rightarrow 4b$		$H \rightarrow \tau\mu, \tau e$		$a \rightarrow \mu\mu$
	$(H \rightarrow) hh \rightarrow bb\tau\tau$		$H \rightarrow e\mu$		$h \rightarrow 2a \rightarrow 4\gamma$ (multiphoton)
	$(H \rightarrow) hh \rightarrow VV\gamma\gamma \rightarrow 4j\gamma\gamma$		$H \rightarrow J/\psi\gamma, Y\gamma$		$h \rightarrow 2a \rightarrow bb\mu\mu$
	$(H \rightarrow) hh \rightarrow WW\gamma\gamma \rightarrow lvqq\gamma\gamma$		$H \rightarrow ZJ/\psi, ZY$		$h \rightarrow 2a \rightarrow bb\tau\tau$
	$A \rightarrow Zh \rightarrow ll\tau\tau$ (LL, LH, HH)		$H \rightarrow \phi\gamma$		$(bb)a \rightarrow (bb)\tau\tau \rightarrow (bb)e\mu$
	$A \rightarrow Zh \rightarrow (ll\nu\nu)bb$		$t \rightarrow cH$ (various)		$h \rightarrow 2a \rightarrow 4\tau$
					$H \rightarrow aW$

Executive Summary:  
No significant excess yet...



# PRELIMINARY CONCLUSIONS

- The properties of the Higgs boson we have found are consistent with SM predictions (within current uncertainties)
  - The SM is a very good effective theory that is consistent with observations at high energy colliders so far
- But
  - The SM is incomplete and faces theoretical problems when we extend it to very high energies
  - SM makes many predictions and is therefore easy to falsify: deviations from the SM may yet appear with more precise measurements
- The road ahead:
  - Continue the characterization program of the discovered Higgs boson
  - Search for Higgs boson physics beyond the Standard Model

# **The Future**

# Run 3 Schedule

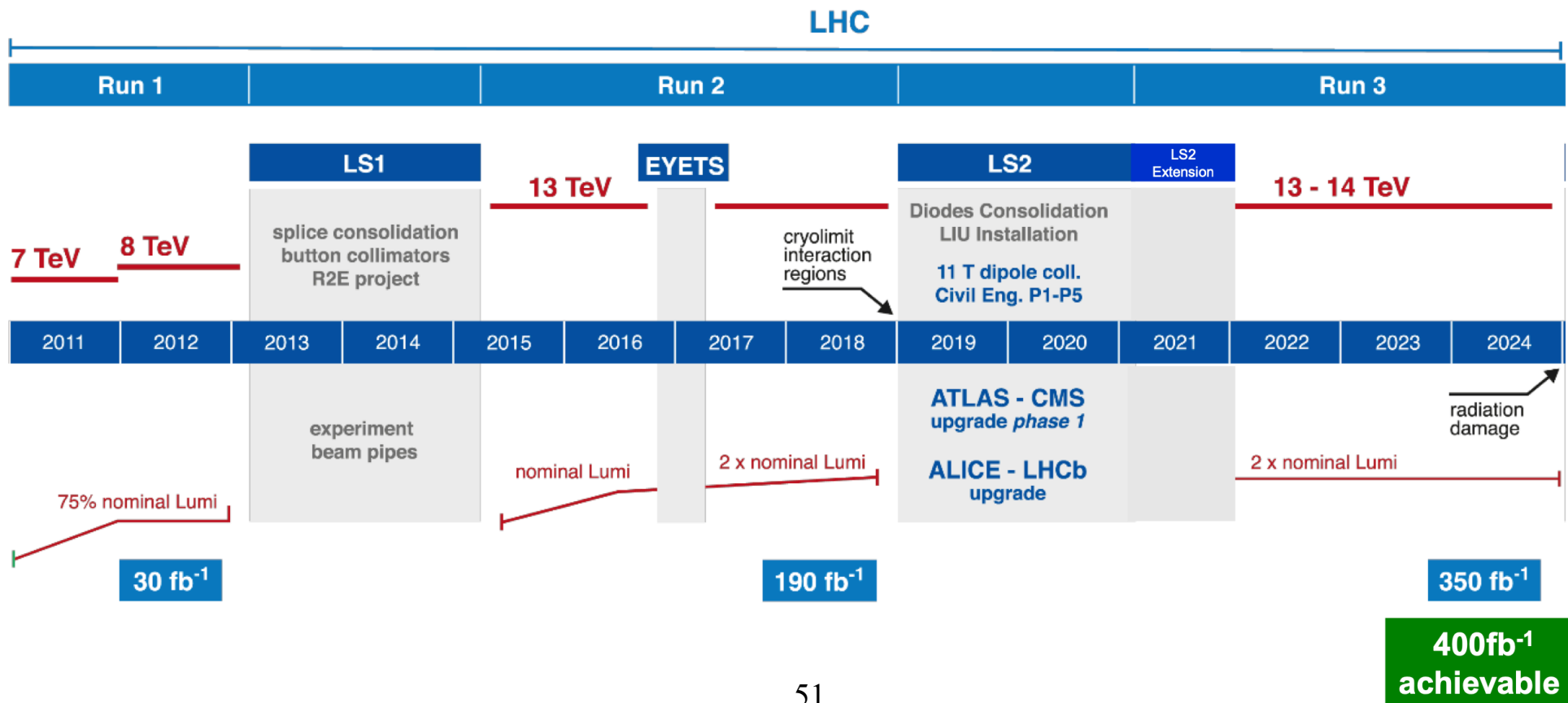
Start of Run 3 has been delayed due to the pandemic

- A 3-year Run
- Plan to run at 2 times nominal luminosity

Slide from F. Bordry

## Run 3 outlook

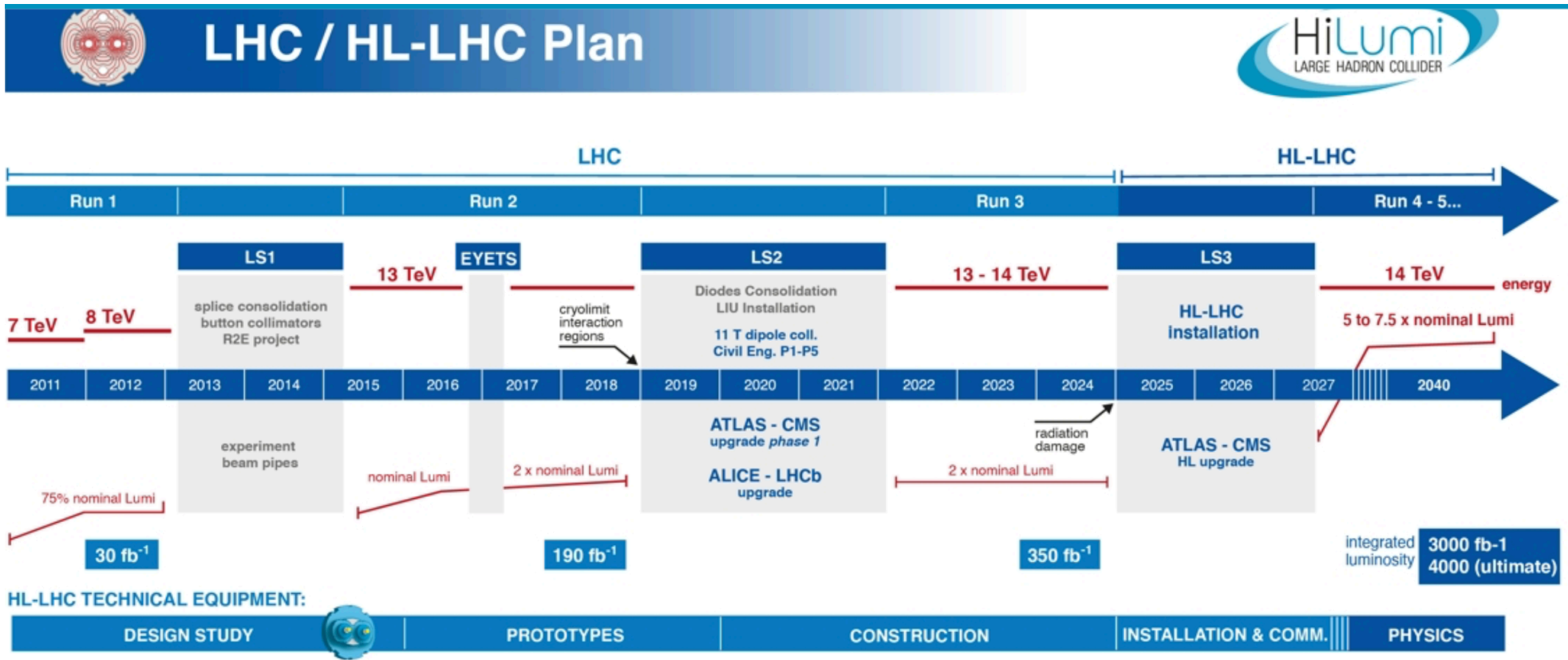
**Goal:**  $\Sigma(\text{Run1} + \text{Run2} + \text{Run 3}) > 350 \text{ fb}^{-1}$



# HL-LHC Schedule

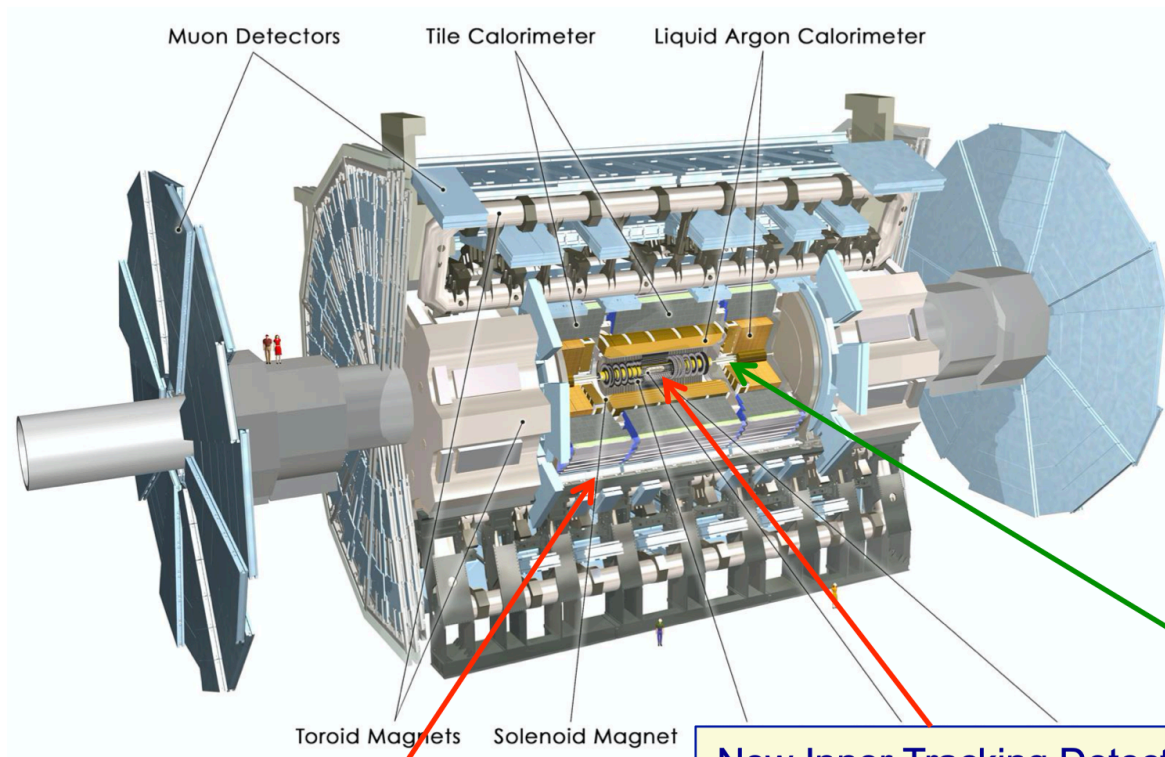
Start of HL-LHC planned for 2027, physics ~2028

- Run at nominal energy of 14 TeV
- Plan to run at 5-7.5 times nominal luminosity



# ATLAS Phase-II Upgrade

From K. Jakobs



## Upgraded Trigger and Data Acquisition System:

- L0: 1 MHz
- Improved High-Level Trigger

## Electronics Upgrade :

- LAr Calorimeter
- Tile Calorimeter
- Muon system

New Inner Tracking Detector  
(all silicon tracker, up to  $|\eta| = 4$ )

New muon chambers  
in the inner barrel region

High granularity timing detector  
(forward region)

Approved by CERN Research Board (16<sup>th</sup> Sept.)

# Run 3: some considerations

- Last 8 years: enormous progress
  - Large increases in integrated luminosity vs time
  - LHC collision energy increase
  - Major improvements in performance and analysis techniques
- Next 8 years: analysis of Run 2 and Run 3 data
  - ~2-2.5 increase in integrated luminosity
  - Modest increase in energy
  - Significant progress possible on many fronts

# Run 3: challenges and opportunities

- Systematic uncertainties are becoming more dominant for many results

	th.	stat.	syst.	bbb				
• CMS H tautau ( $\mu$ )	$0.85^{+0.12}_{-0.11}$	$+0.08$ $-0.07$	$+0.06$ $-0.06$	$+0.07$ $-0.06$ $+0.04$ $-0.03$				
• ATLAS H bb ( $\mu$ )	$1.01 \pm 0.12$	(stat.)	$+0.16$ $-0.15$	(syst.)				
• ATLAS VBF ( $\sigma$ )	$0.85 \pm 0.10$	(stat.)	$+0.08$ $-0.07$	(exp syst.)	$+0.13$ $-0.10$	(sig. theo.)	$+0.07$ $-0.06$	(bkg. theo.) pb
• CMS global $\mu$ fit:	$1.02 \pm 0.04$	(th)	$\pm 0.04$	(exp)	$\pm 0.04$	(stat)		

- Need to reduce systematics (**Theory**, **Modelling**, PDFs), and improve our treatment of systematics

# HL-LHC Higgs Physics

Physics potential of the HL-LHC recently re-appraised and documented in a series of Yellow Reports

Yellow reports were submitted to the European Strategy Group at the end of 2018

Higgs Yellow Report:

**CERN-LPCC-2018-04**

- A huge effort: 400 authors, 340 pages

Can't do justice to the work in this document: I encourage you to take a look and see the full extent and promise of HL-LHC Higgs physics

arXiv:1902.00134v1 [hep-ph] 31 Jan 2019



FERMILAB-PUB-19-074-PPD-T

CERN-LPCC-2018-04  
February 4, 2019

## Higgs Physics at the HL-LHC and HE-LHC

Report from Working Group 2 on the Physics of the HL-LHC, and Perspectives at the HE-LHC

### Convenors:

*M. Cepeda*<sup>1,2</sup>, *S. Gori*<sup>3</sup>, *P. Ilten*<sup>4</sup>, *M. Kado*<sup>5,6,7</sup>, *F. Riva*<sup>8</sup>

### Contributors:

*R. Abdul Khalek*<sup>9,10</sup>, *A. Aboubrahim*<sup>11</sup>, *J. Alimena*<sup>12</sup>, *S. Alioli*<sup>13,13</sup>, *A. Alves*<sup>14</sup>, *C. Asawatangtrakuldee*<sup>15</sup>, *A. Azatov*<sup>16,17</sup>, *P. Azzi*<sup>18</sup>, *S. Bailey*<sup>19</sup>, *S. Banerjee*<sup>20</sup>, *E.L. Barberio*<sup>21</sup>, *D. Barducci*<sup>16</sup>, *G. Barone*<sup>22</sup>, *M. Bauer*<sup>20</sup>, *C. Bautista*<sup>23</sup>, *P. Bechtle*<sup>24</sup>, *K. Becker*<sup>25</sup>, *A. Benaglia*<sup>26</sup>, *M. Bengala*<sup>27</sup>, *N. Berger*<sup>28</sup>, *C. Bertella*<sup>29</sup>, *A. Bethani*<sup>30</sup>, *A. Betti*<sup>30</sup>, *A. Biekotter*<sup>31</sup>, *F. Bishara*<sup>15</sup>, *D. Bloch*<sup>32</sup>, *P. Bokan*<sup>33</sup>, *O. Bondur*<sup>34</sup>, *M. Bowin*<sup>6</sup>, *L. Borgonovi*<sup>35,36</sup>, *M. Borsato*<sup>37</sup>, *S. Boselli*<sup>38</sup>, *S. Braibant-Giacomelli*<sup>35,36</sup>, *G. Buchalla*<sup>39</sup>, *L. Cadamuro*<sup>40</sup>, *C. Caillol*<sup>41</sup>, *A. Calandri*<sup>42,43</sup>, *A. Calderon Tazon*<sup>44</sup>, *J.M. Campbell*<sup>45</sup>, *F. Caola*<sup>20</sup>, *M. Capozzi*<sup>46</sup>, *M. Carena*<sup>45,47</sup>, *C.M. Carloni Calame*<sup>48</sup>, *A. Carmona*<sup>49</sup>, *E. Carquin*<sup>50</sup>, *A. Carvalho Antunes De Oliveira*<sup>51</sup>, *A. Castaneda Hernandez*<sup>52</sup>, *O. Cata*<sup>53</sup>, *A. Celis*<sup>54</sup>, *A. Cerri*<sup>55</sup>, *F. Cerutti*<sup>56,57</sup>, *G.S. Chahal*<sup>58,20</sup>, *A. Chakraborty*<sup>59</sup>, *G. Chaudhary*<sup>60</sup>, *X. Chen*<sup>61</sup>, *A.S. Chisholm*<sup>1,4</sup>, *R. Contino*<sup>62</sup>, *A.J. Costa*<sup>27</sup>, *R. Covarelli*<sup>63,64</sup>, *N. Craig*<sup>65</sup>, *D. Curtin*<sup>66</sup>, *L. D'Eramo*<sup>67</sup>, *M. Dührssen*<sup>1</sup>, *N.P. Dang*<sup>68</sup>, *P. Das*<sup>69</sup>, *S. Dawson*<sup>22</sup>, *O.A. De Aguiar Francisco*<sup>1</sup>, *J. de Blas*<sup>70,18</sup>, *S. De Curtis*<sup>71</sup>, *N. De Filippis*<sup>72,73</sup>, *H. De la Torre*<sup>74</sup>, *L. de Lima*<sup>75</sup>, *A. De Wit*<sup>15</sup>, *C. Delaere*<sup>34</sup>, *M. Delcourt*<sup>34</sup>, *M. Delmastro*<sup>28</sup>, *S. Demers*<sup>76</sup>, *N. Dev*<sup>77</sup>, *R. Di Nardo*<sup>78</sup>, *S. Di Vita*<sup>79</sup>, *S. Dildick*<sup>80</sup>, *L.A.F. do Prado*<sup>81,23</sup>, *M. Donadelli*<sup>82</sup>, *D. Du*<sup>83</sup>, *G. Durieux*<sup>84,15</sup>, *O. Eberhardt*<sup>85</sup>, *K. El Morabit*<sup>86</sup>, *J. Elias-Miro*<sup>1</sup>, *J. Ellis*<sup>87,51,1</sup>, *C. Englert*<sup>88</sup>, *S. Falke*<sup>28</sup>, *M. Farina*<sup>89</sup>, *A. Ferrari*<sup>90</sup>, *A. Ferrogli*<sup>91</sup>, *M.C.N. Fiolhais*<sup>92</sup>, *M. Flechl*<sup>93</sup>, *S. Folgueras*<sup>94</sup>, *E. Fontanesi*<sup>35,36</sup>, *P. Francavilla*<sup>67,95</sup>, *R. Franceschini*<sup>96,97</sup>, *R. Frederix*<sup>98</sup>, *S. Frixione*<sup>99</sup>, *G. Gómez-Ceballos*<sup>100</sup>, *A. Gabrielli*<sup>56,57</sup>, *S. Gadatsch*<sup>1</sup>, *M. Gallinaro*<sup>27</sup>, *A. Gandrakota*<sup>101</sup>, *J. Gao*<sup>102</sup>, *F.M. Garay Walls*<sup>103</sup>, *T. Gehrmann*<sup>61</sup>, *Y. Gershtein*<sup>101</sup>, *T. Ghosh*<sup>104</sup>, *A. Gilbert*<sup>1</sup>, *R. Gleis*<sup>105</sup>, *E.W.N. Glover*<sup>20</sup>, *R. Gomez-Ambrosio*<sup>20</sup>, *R. Gonçalo*<sup>27</sup>, *D. Gonçalves*<sup>106</sup>, *M. Gorbahn*<sup>107</sup>, *E. Gouveia*<sup>27</sup>, *M. Gouzevitch*<sup>108</sup>, *P. Govoni*<sup>26,13</sup>, *M. Grazzini*<sup>61</sup>, *B. Greenberg*<sup>101</sup>, *K. Grimm*<sup>109</sup>, *A.V. Gritsan*<sup>110</sup>, *A. Grohsjean*<sup>15</sup>, *C. Grojean*<sup>15</sup>, *J. Gu*<sup>111</sup>, *R. Gugel*<sup>25</sup>, *R.S. Gupta*<sup>20</sup>, *C.B. Gwilliam*<sup>112</sup>, *S. Höche*<sup>113</sup>, *M. Haacke*<sup>103</sup>, *Y. Haddad*<sup>58</sup>, *U. Haisch*<sup>46</sup>, *G.N. Hamity*<sup>114</sup>, *T. Han*<sup>106</sup>, *L.A. Harland-Lang*<sup>19</sup>, *R. Harnik*<sup>45</sup>, *S. Heinemeyer*<sup>44,115</sup>, *G. Heinrich*<sup>46</sup>, *B. Henning*<sup>8</sup>, *V. Hirschi*<sup>43</sup>, *K. Hoepfner*<sup>116</sup>, *J.M. Hogan*<sup>117,118</sup>, *S. Homiller*<sup>119,22</sup>, *Y. Huang*<sup>120</sup>, *A. Huss*<sup>1</sup>, *S. Jézéquel*<sup>28</sup>, *S. Jain*<sup>69</sup>, *S.P. Jones*<sup>1</sup>, *K. Köneke*<sup>25</sup>, *J. Kalinowski*<sup>121</sup>, *J.F. Kamenik*<sup>122,123</sup>, *M. Kaplan*<sup>100</sup>, *A. Karlberg*<sup>61</sup>, *M. Kaur*<sup>60</sup>, *P. Keicher*<sup>86</sup>, *M. Kerner*<sup>61</sup>, *A. Khanov*<sup>124</sup>, *J. Kieseler*<sup>1</sup>, *J.H. Kim*<sup>125</sup>, *M. Kim*<sup>126</sup>, *T. Klijnsma*<sup>43</sup>, *F. Kling*<sup>127</sup>, *M. Klute*<sup>100</sup>, *J.R. Komaragiri*<sup>128</sup>, *K. Kong*<sup>125</sup>, *J. Kozaczuk*<sup>129</sup>, *P. Kozow*<sup>121</sup>, *C. Krause*<sup>45</sup>, *S. Lai*<sup>33</sup>, *J. Langford*<sup>38</sup>, *B. Le*<sup>21</sup>, *L. Lechner*<sup>93</sup>, *W.A. Leigh*<sup>130</sup>, *K.J.C. Leney*<sup>131</sup>, *T. Lenz*<sup>24</sup>, *C.-Q. Li*<sup>132</sup>, *H. Li*<sup>83</sup>, *Q. Li*<sup>133</sup>, *S. Liebler*<sup>134</sup>, *J. Lindert*<sup>20</sup>, *D. Liu*<sup>135</sup>, *J. Liu*<sup>136</sup>, *Y. Liu*<sup>137</sup>, *Z. Liu*<sup>38,45</sup>, *D. Lombardo*<sup>8</sup>, *A. Long*<sup>139</sup>, *K. Long*<sup>41</sup>, *I. Low*<sup>140,135</sup>, *G. Luisoni*<sup>46</sup>, *L.L. Ma*<sup>83</sup>, *A.-M. Magnan*<sup>58</sup>, *D. Majumder*<sup>125</sup>, *A. Malinauskas*<sup>19</sup>, *F. Maltoni*<sup>141</sup>, *M.L. Mangano*<sup>1</sup>, *G. Marchioni*<sup>27,67</sup>, *A.C. Marini*<sup>100</sup>, *A. Martin*<sup>77</sup>, *S. Marzani*<sup>142,99</sup>, *A. Massironi*<sup>1</sup>, *K.T. Matchev*<sup>40,143</sup>, *R.D. Matheus*<sup>23</sup>, *K. Mazumdar*<sup>69</sup>, *J. Mazzielli*<sup>61</sup>, *A.E. Mcdougall*<sup>21</sup>, *P. Meade*<sup>119</sup>, *P. Meridiani*<sup>6</sup>, *A.B. Meyer*<sup>15</sup>, *E. Michielin*<sup>18</sup>, *P. Milenovic*<sup>1,144</sup>, *V. Milosevic*<sup>58</sup>, *K. Mimasu*<sup>141</sup>, *B. Mistlberger*<sup>145</sup>, *M. Mlynarikova*<sup>146</sup>, *M. Mondragon*<sup>147</sup>, *P.F. Monni*<sup>1</sup>, *G. Montagna*<sup>148,48</sup>, *F. Monti*<sup>26,13</sup>, *M. Moreno Llacer*<sup>1</sup>, *A. Mueck*<sup>149</sup>, *P.C. Muñio*<sup>27</sup>, *C. Murphy*<sup>150</sup>, *W.J.*



# HL-LHC Higgs Physics

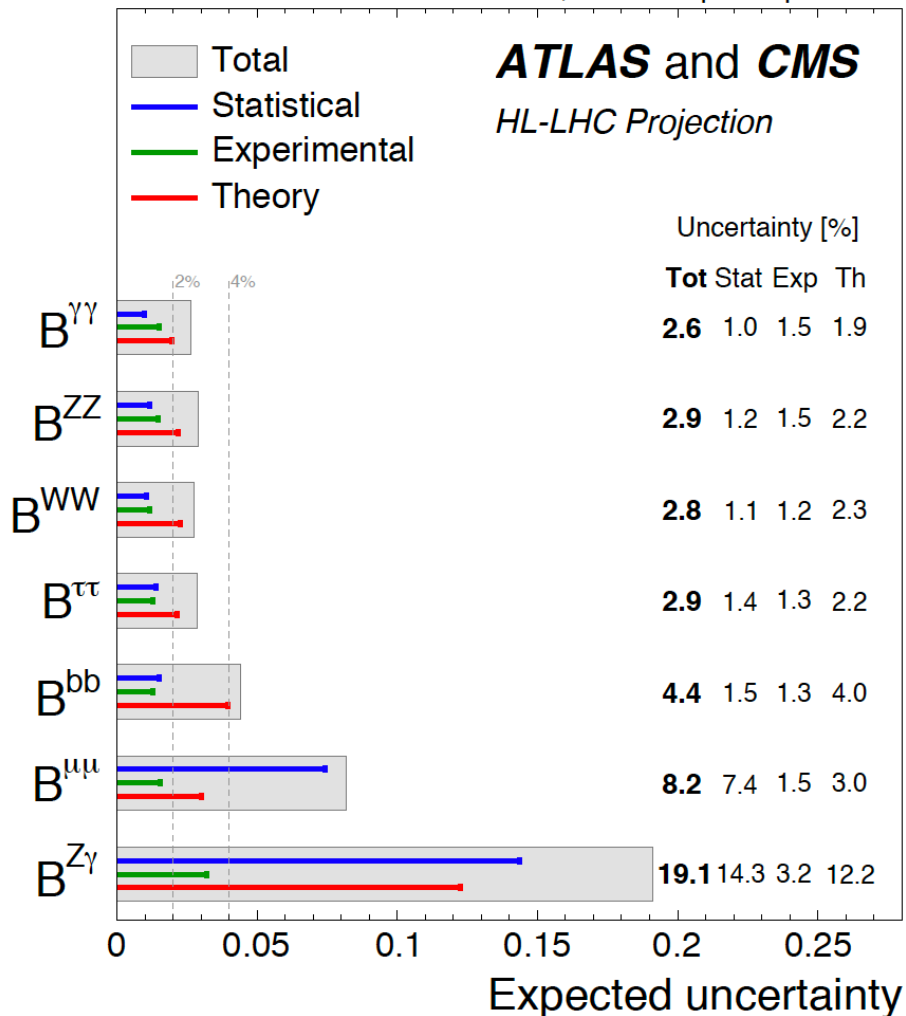
Two uncertainty scenarios were considered for the projections:

- S1 - Conservative, based on the current Run2 systematic uncertainties (including theory)
- S2 - Ultimate, based on estimates of ultimate performance for experimental uncertainties, and applying a factor of 1/2 for theoretical uncertainties

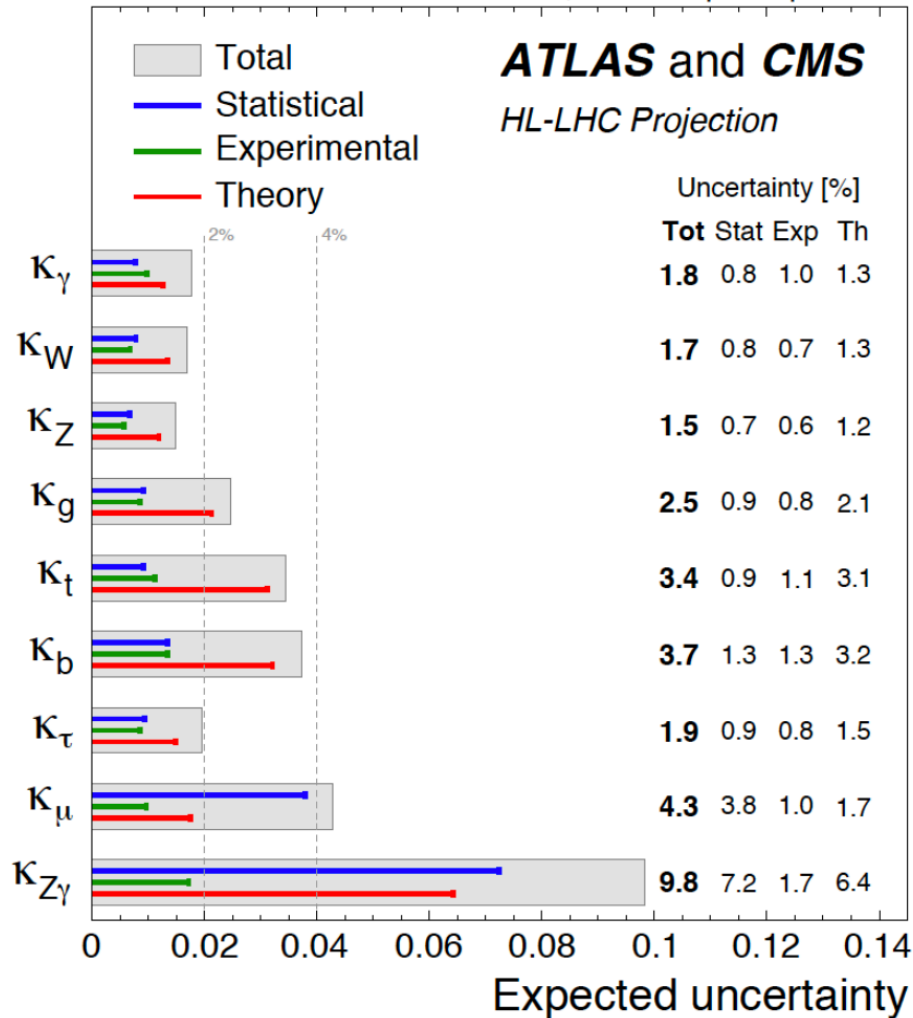
Source	Component	Run 2 uncertainty	Projection minimum uncertainty
Muon ID		1–2%	0.5%
Electron ID		1–2%	0.5%
Photon ID		0.5–2%	0.25–1%
Hadronic tau ID		6%	2.5%
Jet energy scale	Absolute	0.5%	0.1–0.2%
	Relative	0.1–3%	0.1–0.5%
	Pileup	0–2%	Same as Run 2
	Method and sample	0.5–5%	No limit
	Jet flavour	1.5%	0.75%
	Time stability	0.2%	No limit
	Jet energy res.		Varies with $p_T$ and $\eta$
MET scale		Varies with analysis selection	Half of Run 2
b-Tagging	b-/c-jets (syst.)	Varies with $p_T$ and $\eta$	Same as Run 2
	light mis-tag (syst.)	Varies with $p_T$ and $\eta$	Same as Run 2
	b-/c-jets (stat.)	Varies with $p_T$ and $\eta$	No limit
	light mis-tag (stat.)	Varies with $p_T$ and $\eta$	No limit
Integrated lumi.		2.5%	1%

# HL-LHC: Branching Ratios and Cross Sections

$\sqrt{s} = 14 \text{ TeV}$ ,  $3000 \text{ fb}^{-1}$  per experiment



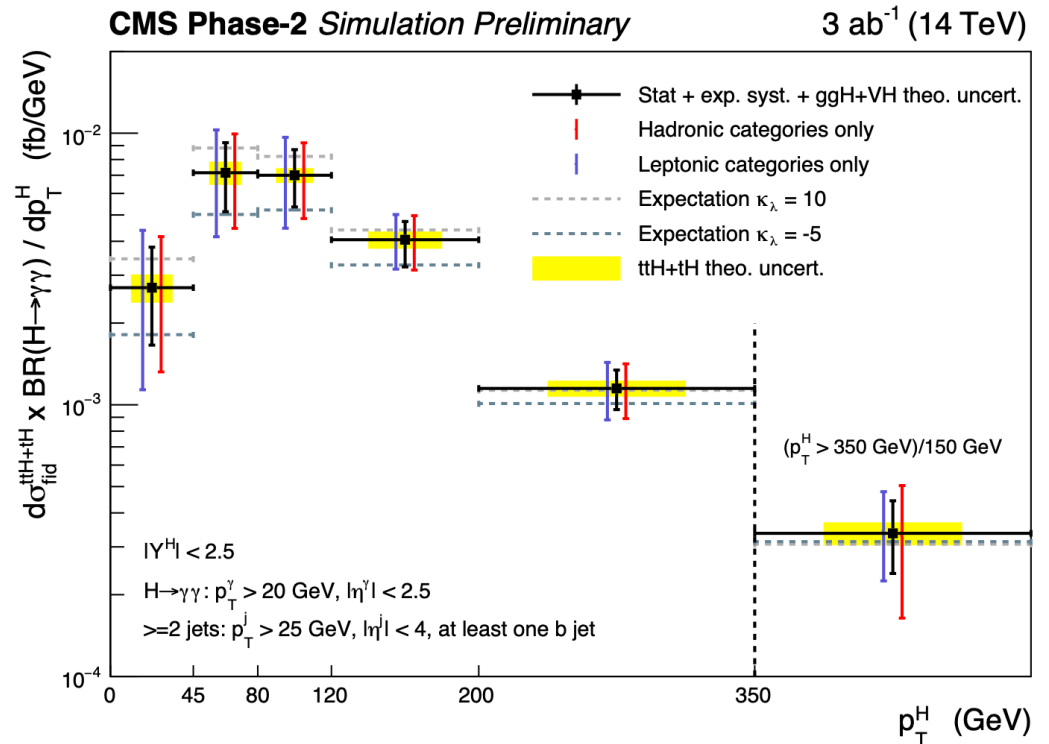
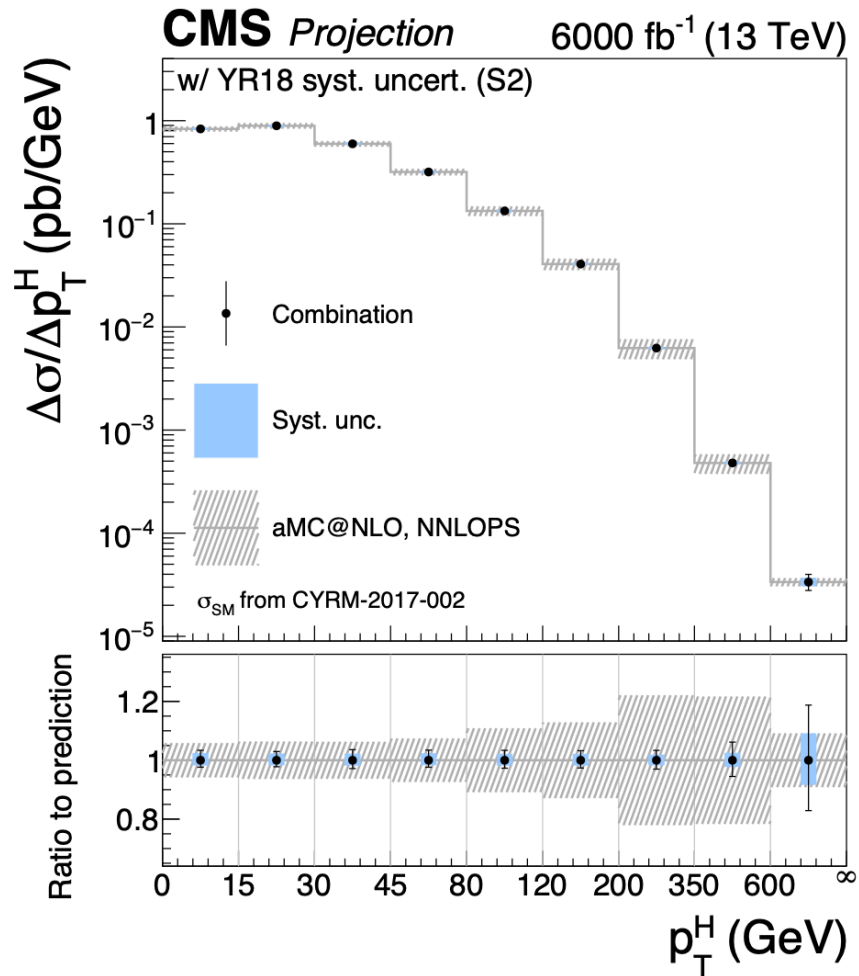
$\sqrt{s} = 14 \text{ TeV}$ ,  $3000 \text{ fb}^{-1}$  per experiment



Combination of ATLAS and CMS for systematic uncertainty scenario 2

- Theory uncertainty remains the largest component for most measurements

# HL-LHC: Differential Cross Sections



CMS-PAS-FTR-18-020

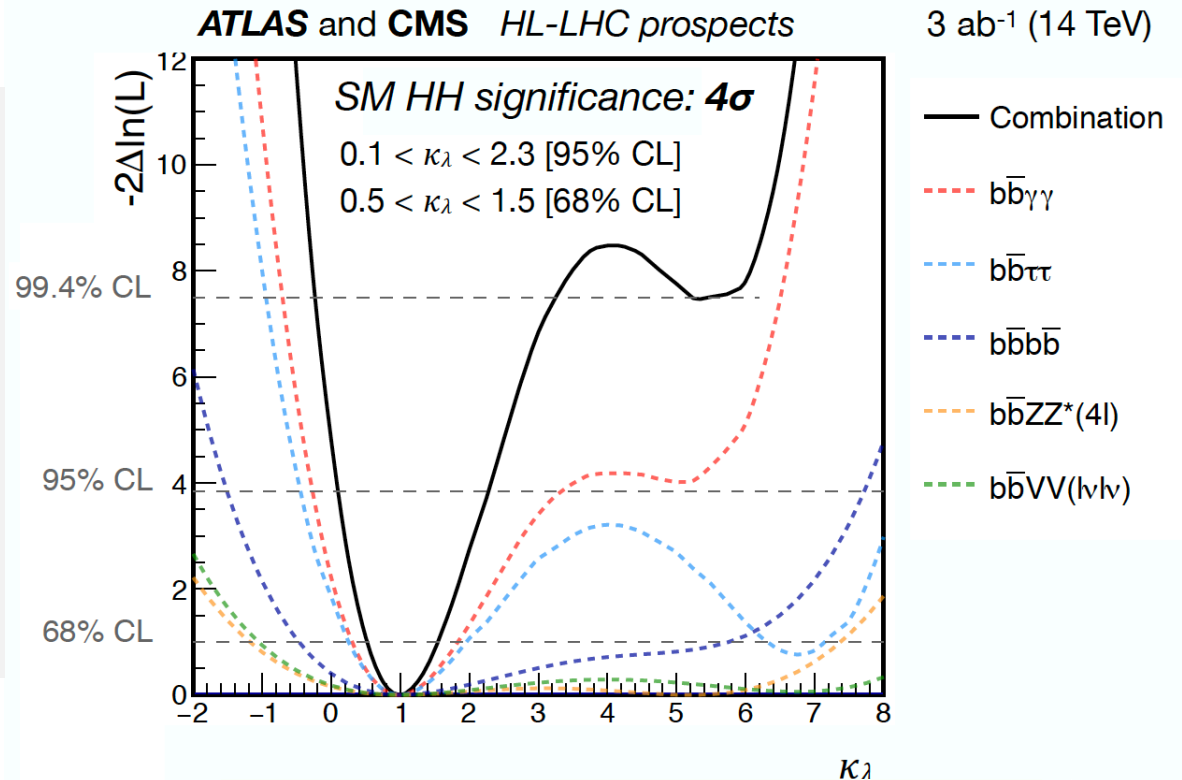
Left plot: results for CMS for systematic uncertainty scenario 2

- Theory uncertainty dominates in all bins except  $p_T > 600$  GeV

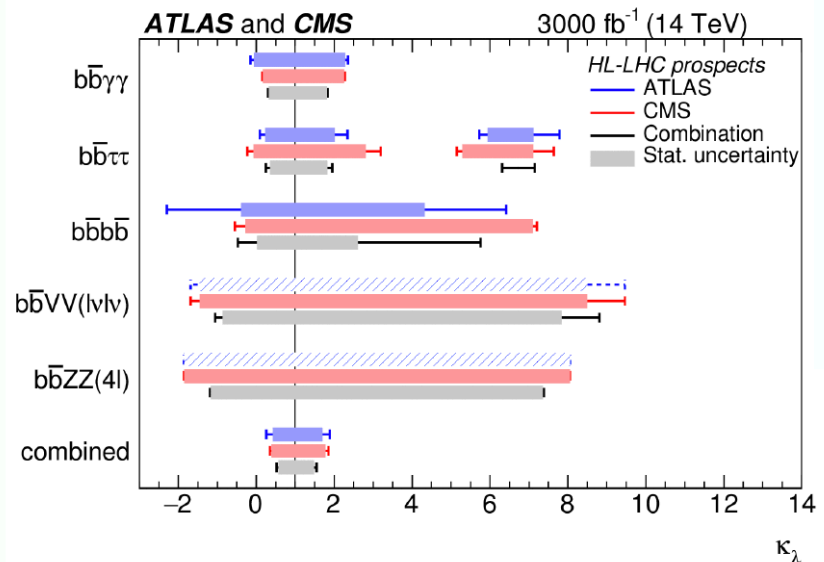
Right plot:  $p_T^H$  distribution for ttH production using the  $H \rightarrow \gamma\gamma$  decay

# HL-LHC: Higgs Self-Coupling

- Significance of HH signal at the  $4\sigma$  level (both exp.)
- Uncertainty on  $\kappa_\lambda$  of 50%
  - 2<sup>nd</sup> minimum excluded at > 99% CL
- Note that HH observation analysis and  $\kappa_\lambda$  analysis require different optimizations



	Statistical-only		Statistical + Systematic	
	ATLAS	CMS	ATLAS	CMS
$HH \rightarrow b\bar{b}b\bar{b}$	1.4	1.2	0.61	0.95
$HH \rightarrow b\bar{b}\tau\tau$	2.5	1.6	2.1	1.4
$HH \rightarrow b\bar{b}\gamma\gamma$	2.1	1.8	2.0	1.8
$HH \rightarrow b\bar{b}VV(l\nu\nu)$	-	0.59	-	0.56
$HH \rightarrow b\bar{b}ZZ(4l)$	-	0.37	-	0.37
combined	3.5	2.8	3.0	2.6
	Combined 4.5		Combined 4.0	



# Some Concluding Thoughts

After Run 2, there is still a long way to go before we reach the summit

Slower dataset doubling time and the fact that systematics are already larger than statistical uncertainties in many analyses means that we will need to **work hard and innovate**:

- True for both experimental and theory communities
- Studies in Higgs HL-LHC Yellow Report imply that we will need major advances in theory to reach the “summit” i.e. exploit the full potential of HL-LHC



Unclear if/when new machines will provide results competitive with most LHC results:

- This challenging work that will be done in coming years could have a very, very long shelf life – worth the effort!

# Summary

- The theory that describes fundamental particles and forces (the Standard Model) predicted the existence of a new particle: the Higgs boson
- More than 40 years after it was postulated, the particle was finally observed in July 2012 by two scientific collaborations at CERN
- The field associated with this particle gives mass to all massive particles. This field is “on” everywhere. It fills the physical vacuum
- This discovery has important implications on cosmology and our understanding of the very early Universe
- 8 years after the discovery, we have entered the era of “precision Higgs physics”. Up until now, the Standard Model describes the data very well
- In the years to come, we will increase the size of our dataset by more than a factor of 10. The Standard Model will face much more stringent tests

# **Backup Slides**