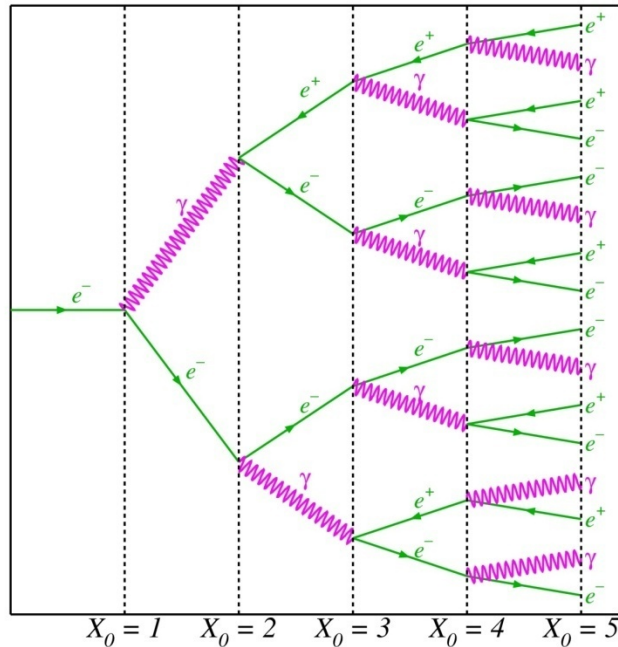


# Electromagnetic Showers



incident  $\gamma$   $E_0$   
 $\chi_0 \downarrow$

Converts to 2 electrons  $E_e = \frac{E_0}{2}$

$\sim 2\chi_0 \downarrow$

Each electrons will have emitted a brems photon of energy  $E_\gamma = \frac{E_0}{4}$

Now have 4 particles energy  $\frac{E_0}{4}$

Number of particles after  $t \cdot \chi_0$   $N = 2^t$

Average energy  $E(t) = \frac{E_0}{2^t}$

At the critical energy  $E_t(\max) = \frac{E_0}{2^{t_{\max}}} = E_c$

*assume this is max depth*

$$N_{\max} \approx \frac{E_0}{E_c} \quad t_{\max} = \frac{\ln(E_0/E_c)}{\ln 2}$$

$$t_{\max} \sim \ln E_0$$

shower grows as  $\ln E$

$$N_{\max} \propto E_0$$

linear energy measurement

$$\sigma_E \sim \sqrt{N} \sim \sqrt{E}$$

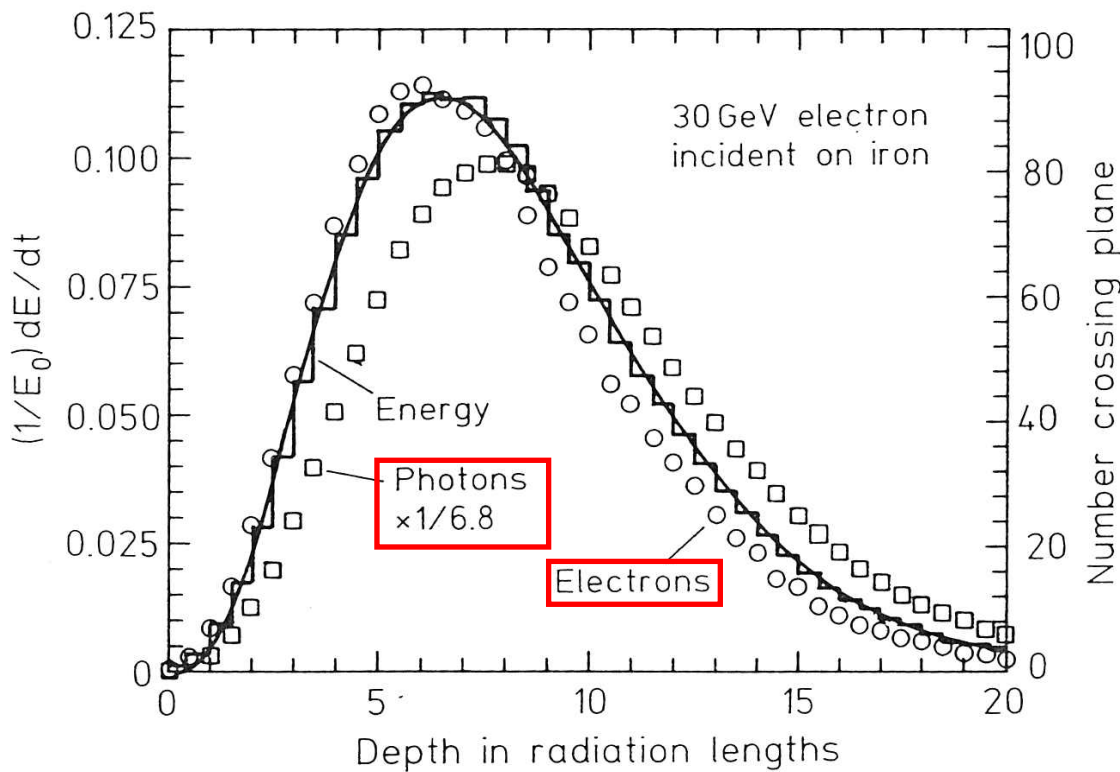
resolution improves with energy

# Monte Carlo Simulation of an EM Shower

$$\frac{dE}{dt} = E_0 \frac{b(bt)^{a-1} e^{-bt}}{\Gamma(a)}$$

$\Rightarrow t_{\max} = \frac{a-1}{b} = 1.0(\ln y + C_i)$

$t_{\max} \sim \ln E_0$   
 $y = E/E_c$   
 $C_i = -0.5(e), +0.5(\gamma)$



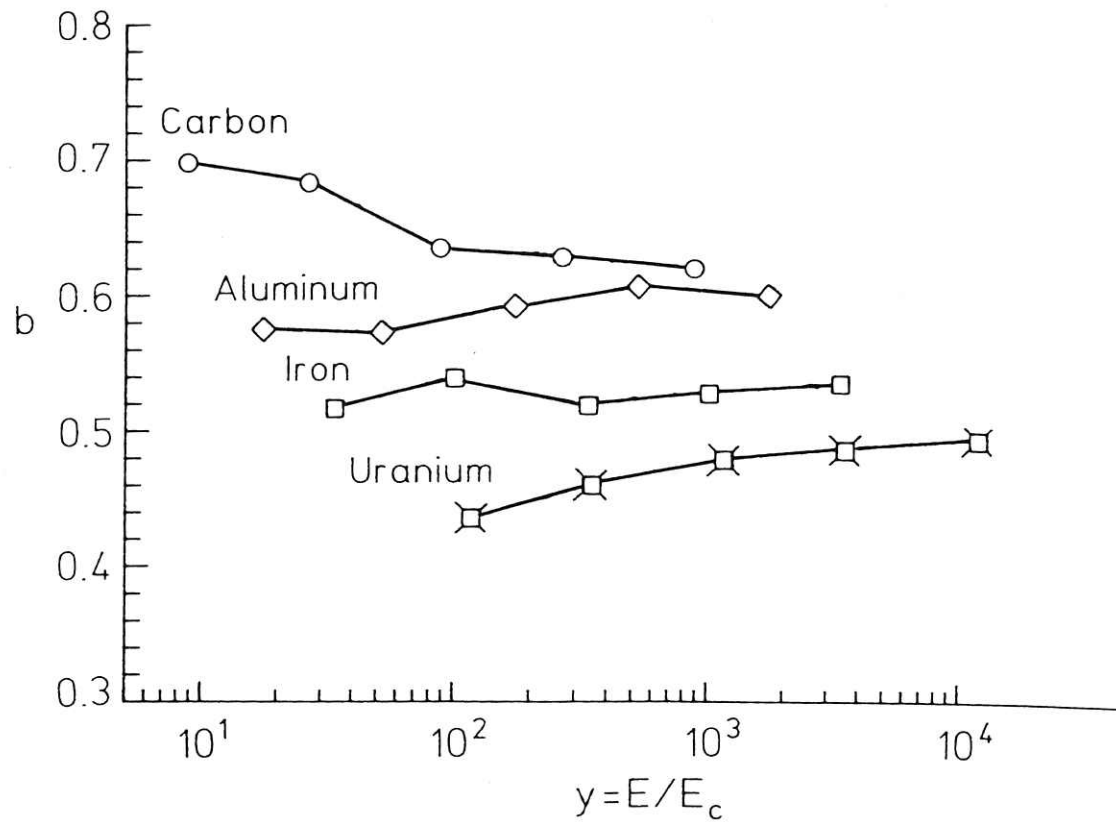
Calc  $t_{\max} = (\ln y + C_i)$

Put  $b=0.5$

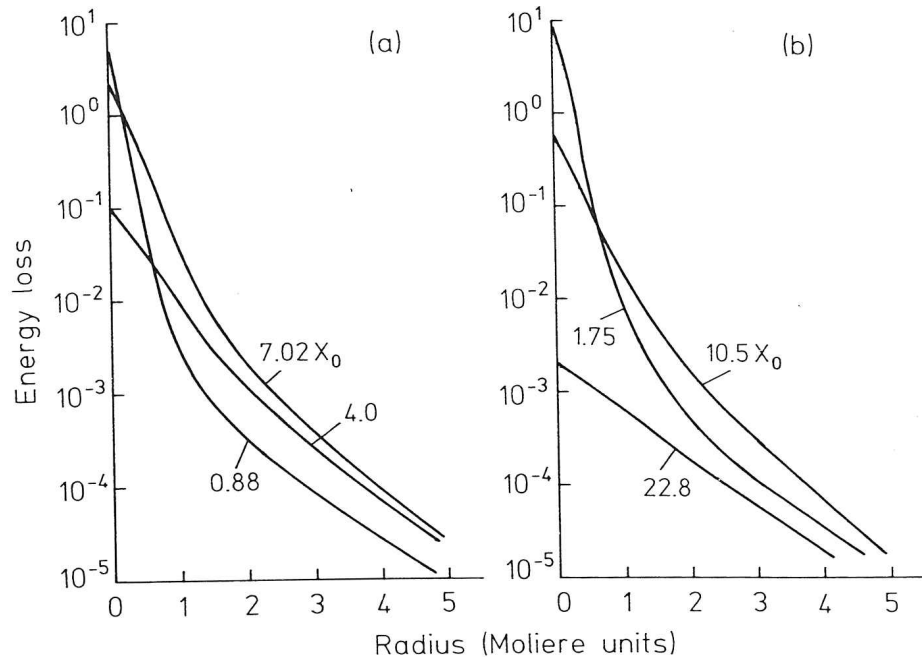
Get  $a$  from  $t_{\max} = \frac{a-1}{b}$

$\frac{dE}{dt}$

“b” is sort of material independent



# Transverse Shower Profile



- Shower Broadens as it develops

- Pair
- Brems
- Compton
- Multiple Coulomb

- Shower Broadens as it develops

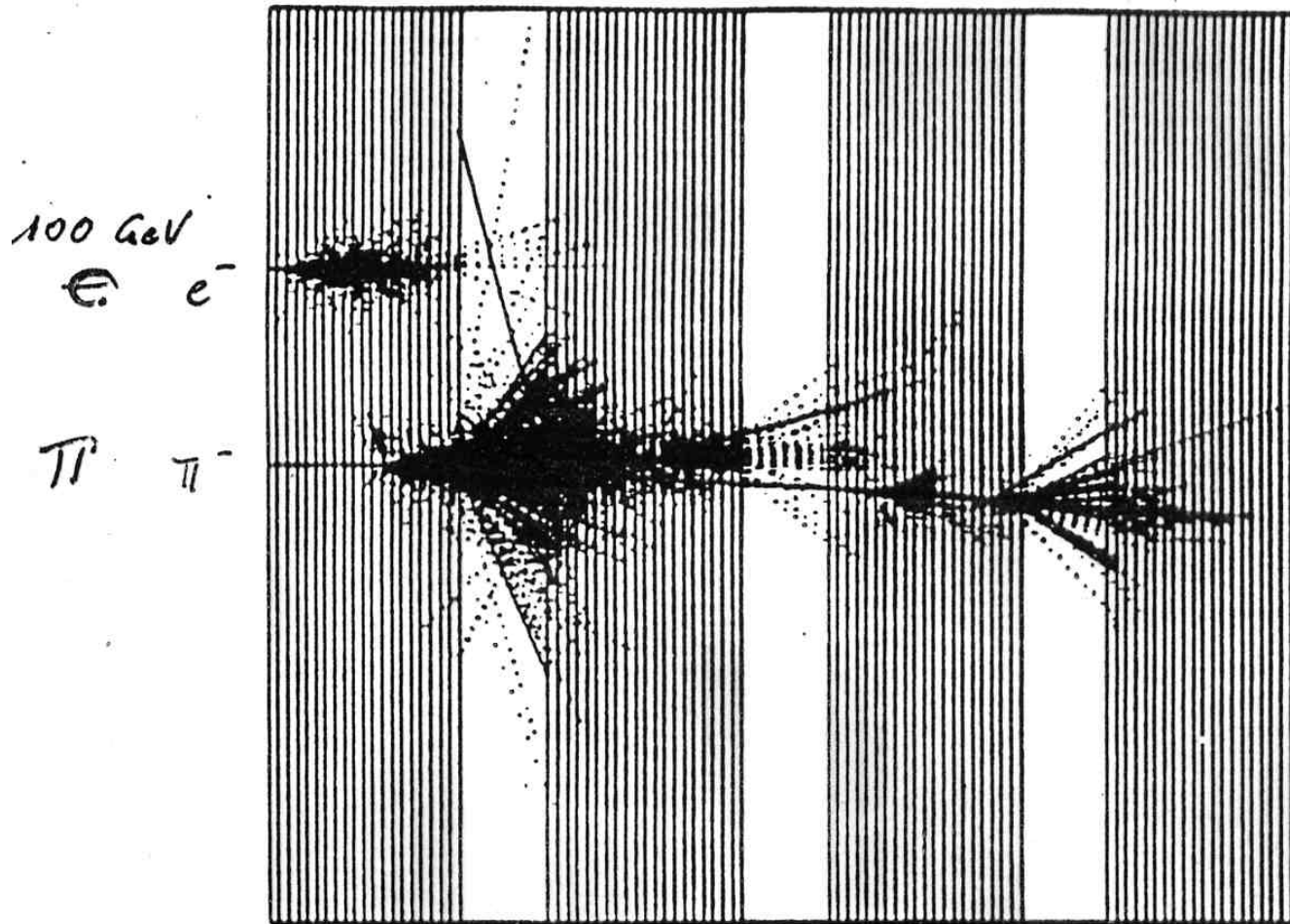
- dense central core
- spreading with depth

- Molière Radius

$$R_M = \chi_0 \frac{E_S}{E_C} \quad E_S = m_e c^2 \sqrt{\frac{4\pi}{\alpha}} = 21.2 \text{ MeV}$$

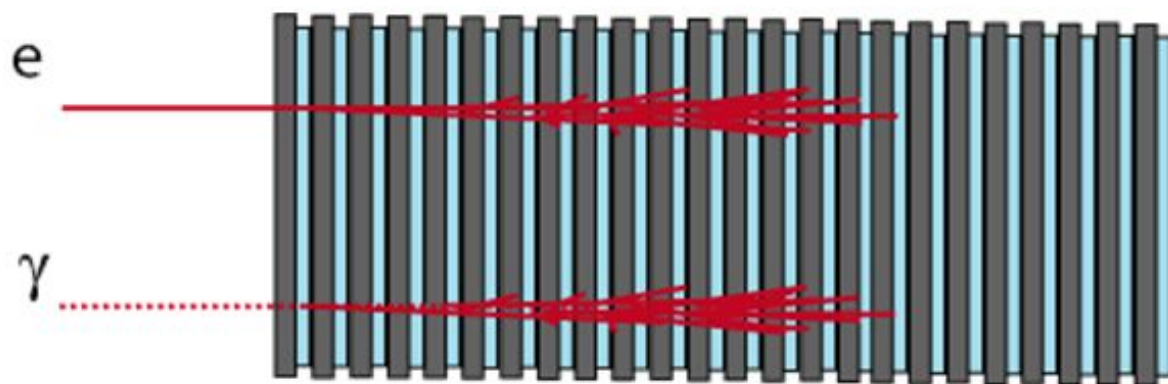
- Like radiation length, Molière radius scales for different materials
- In terms of Molière radius, shower width is roughly independent of material - 90% of energy in  $2 \times R_M$

# Comparison of Hadronic & Electromagnetic Showers



# Electromagnetic Calorimeter Types

- “lead-scintillator sandwich” calorimeter



Energy resolutions:

$$\square E/E \sim 20\%/\sqrt{E}$$

- exotic crystals (BGO, PbW, ...)



$$\square E/E \sim 1\%/\sqrt{E}$$

- liquid argon calorimeter

$$\square E/E \sim \frac{13\%}{\sqrt{E}}$$

$$18\%/\sqrt{E}$$

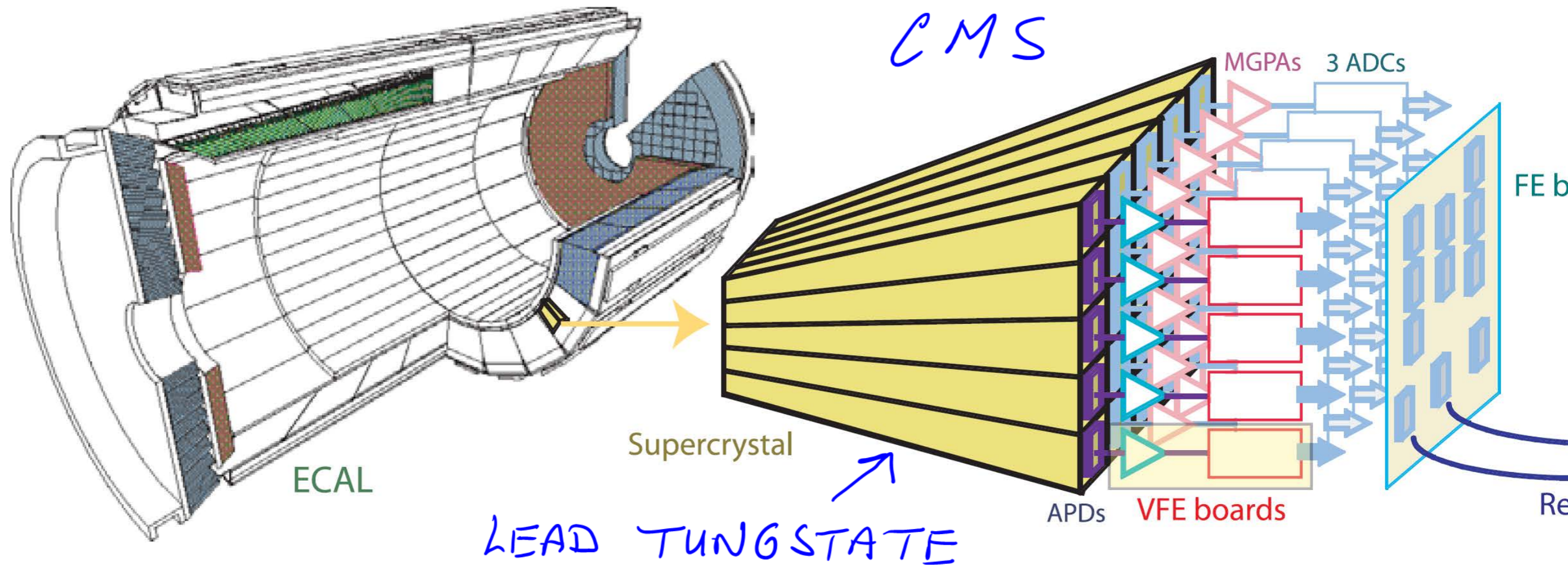


Fig. 2. Design of the CMS Supercrystal detector.

# First Prototype

- Designed between January and April 1990
- Build between April and July
- Exposed to test beam in July-August, using the cryostat and FE electronics of Helios expt
- Demonstrated the concept was sound, although the electronics was not yet fast enough

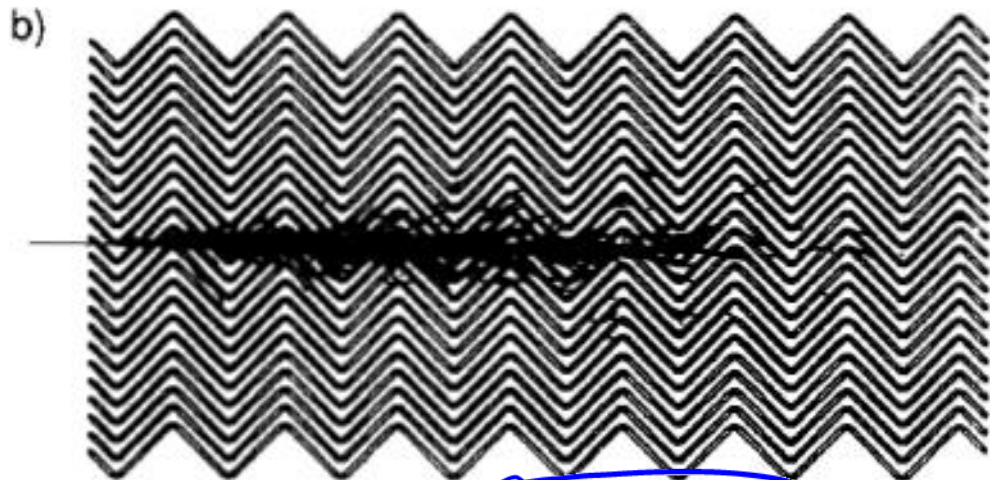
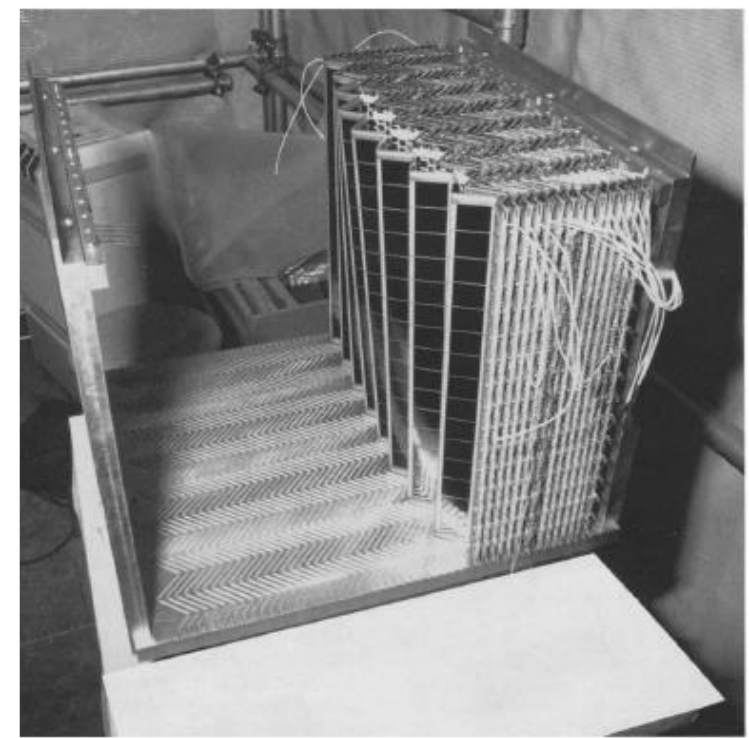
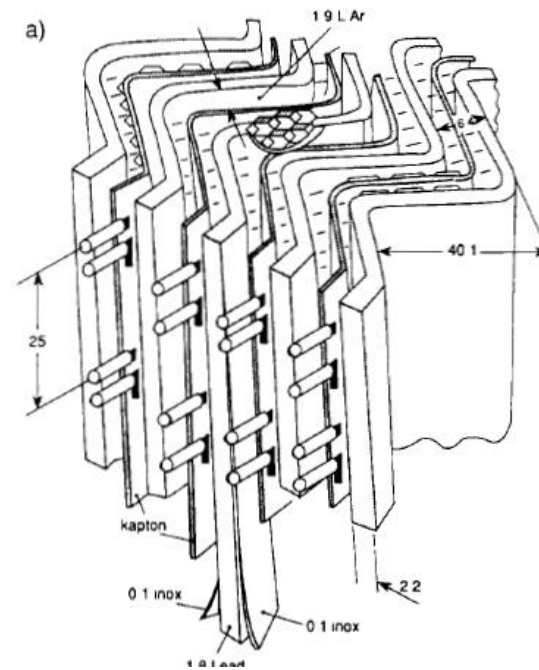


Fig. 2. (a) Artist's view of the accordion calorimeter geometry. (b) Development of a 40 GeV electron shower (Monte Carlo simulation). Only charged tracks above 10 MeV are shown.

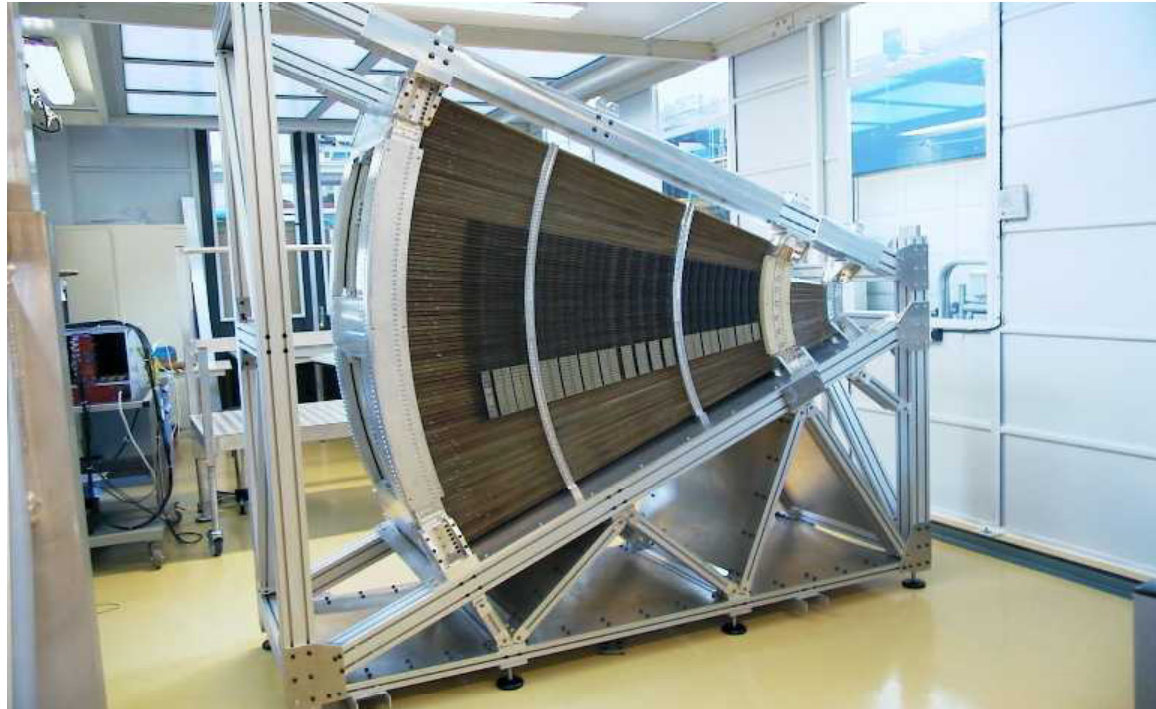
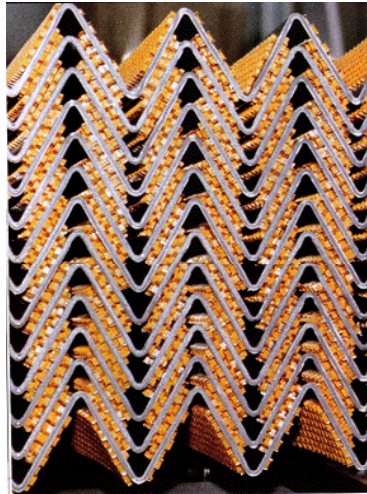




# ATLAS

## Prototype of EM endcap

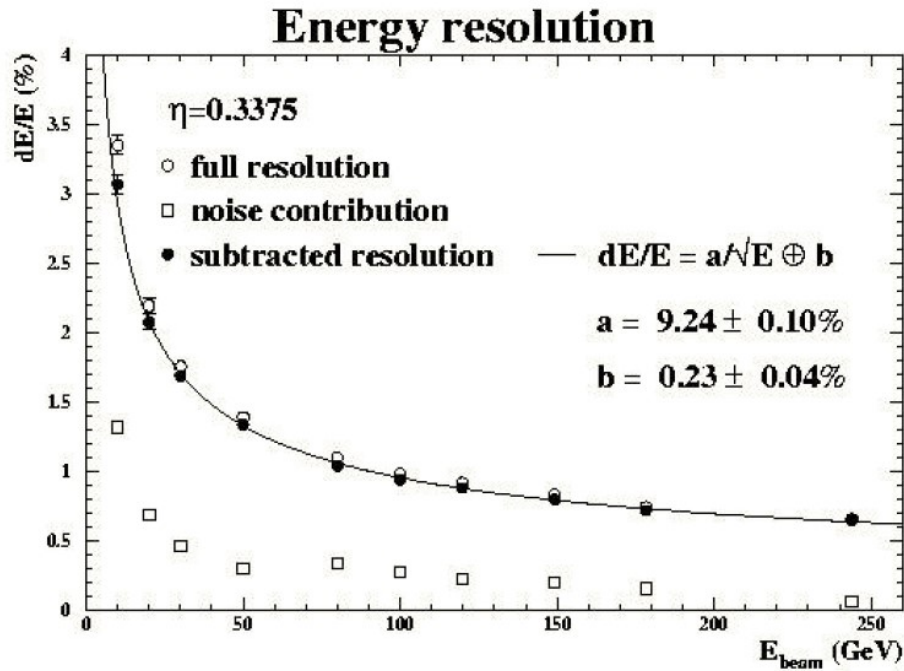
Detail of Kaptons



↑  
ACCORDIAN  
STRUCTURE

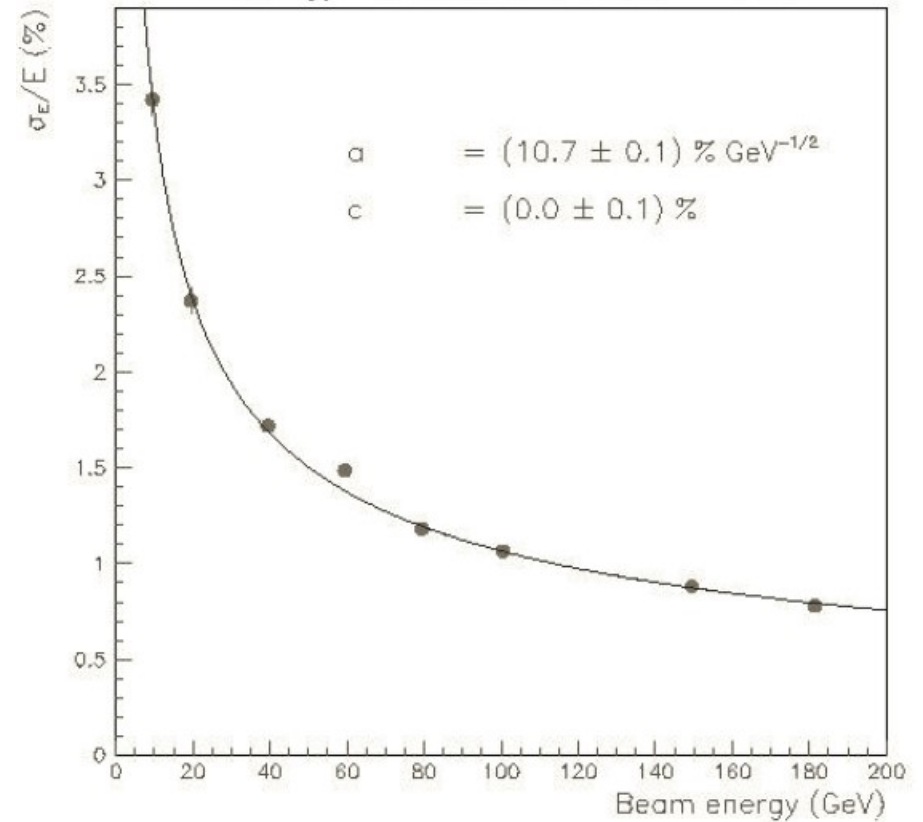
# ATLAS EM CAL ENERGY RESOLUTION

BARREL



END CAP

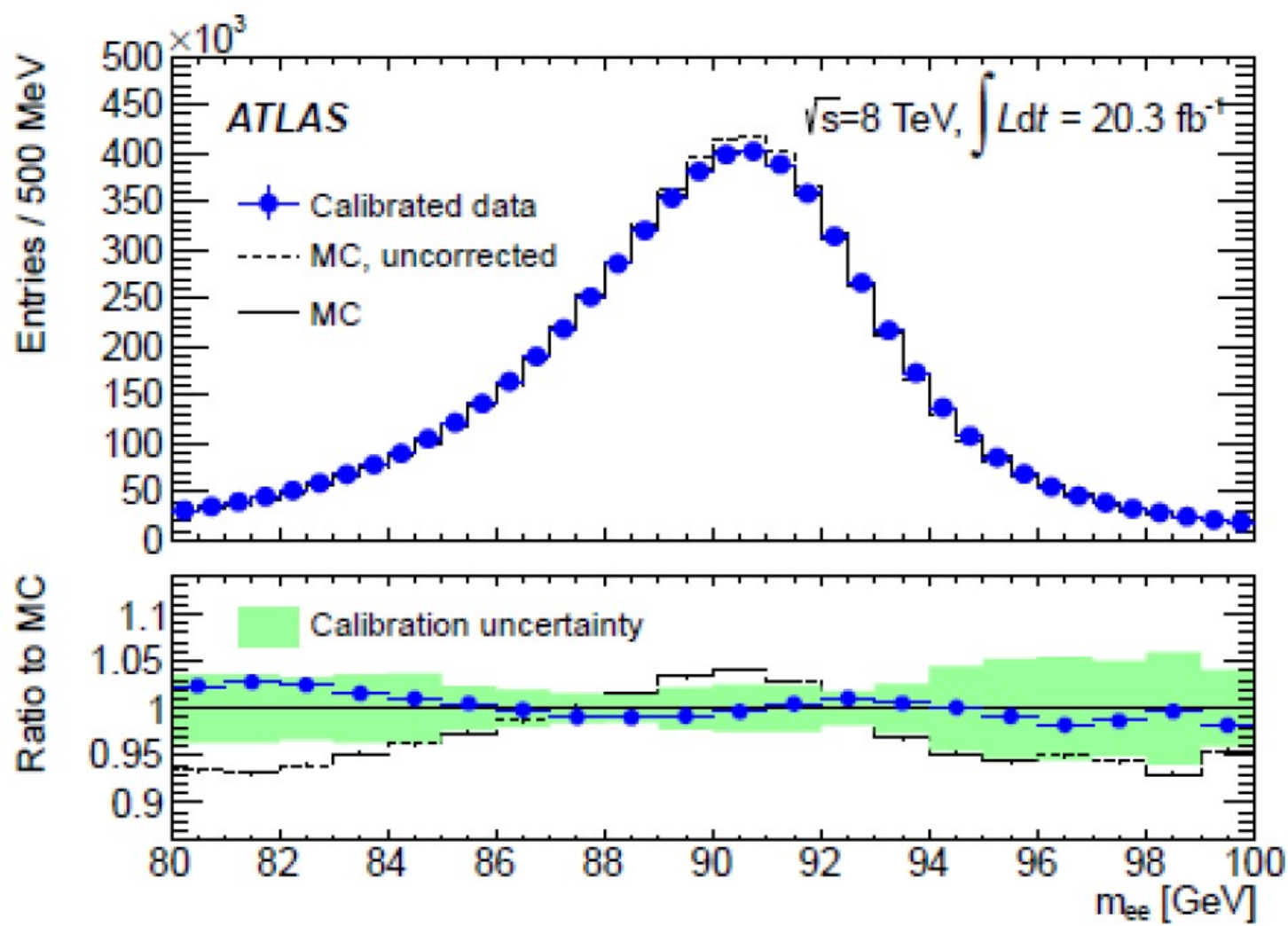
Energy resolution at  $\eta = 1.9$



ATLAS

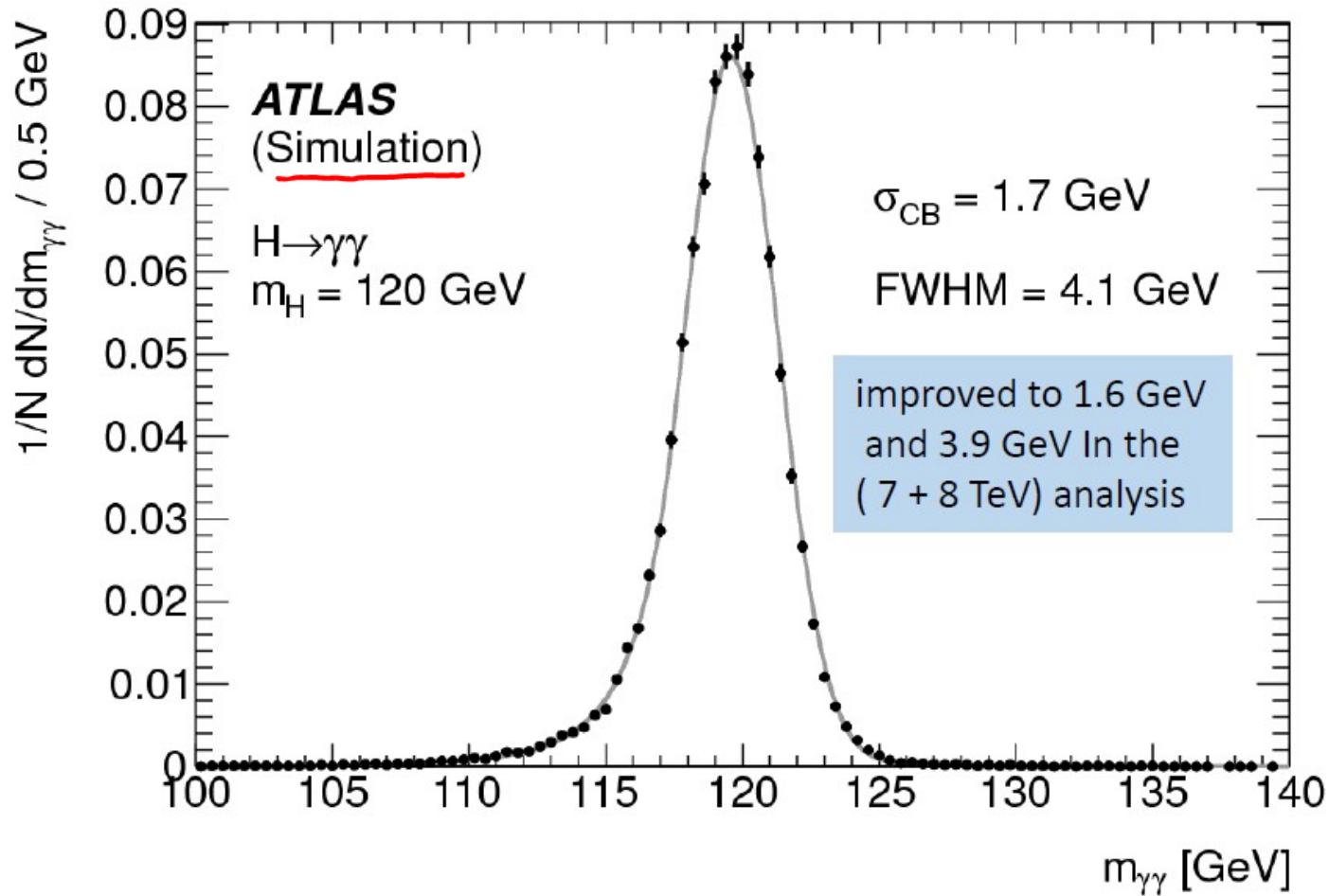
EM

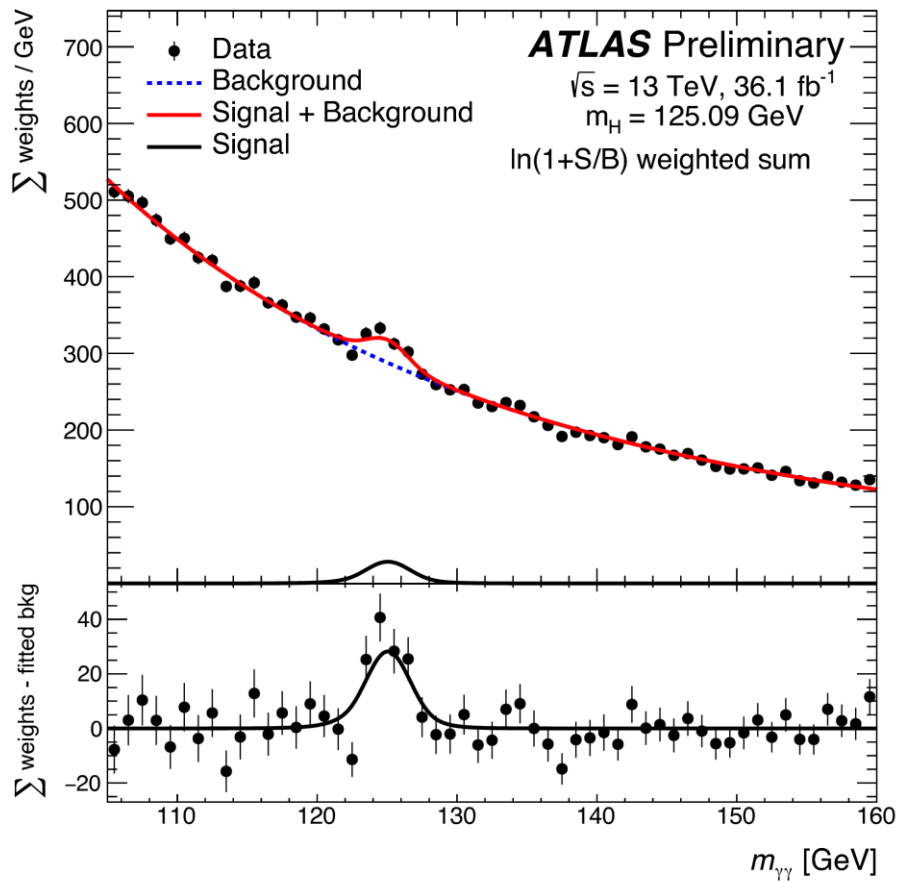
$Z^0 \rightarrow e^+e^-$



ATLAS EM CAL

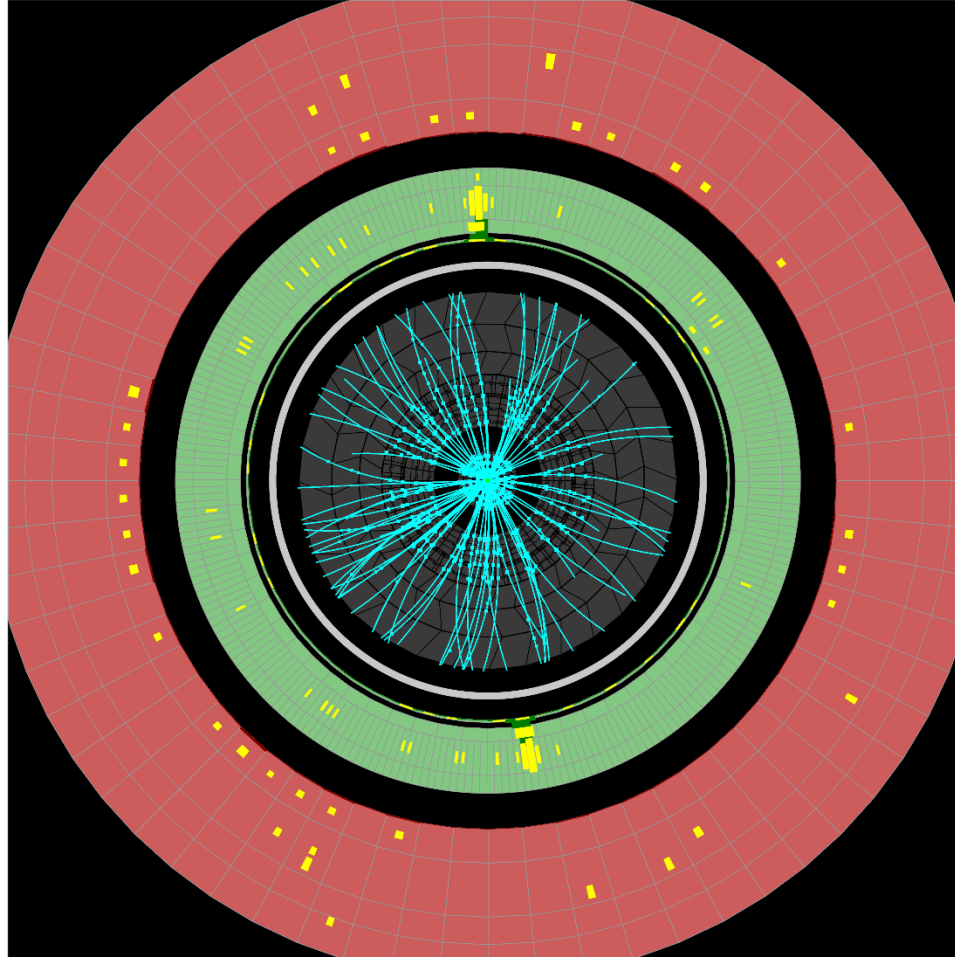
$H \rightarrow \gamma\gamma$





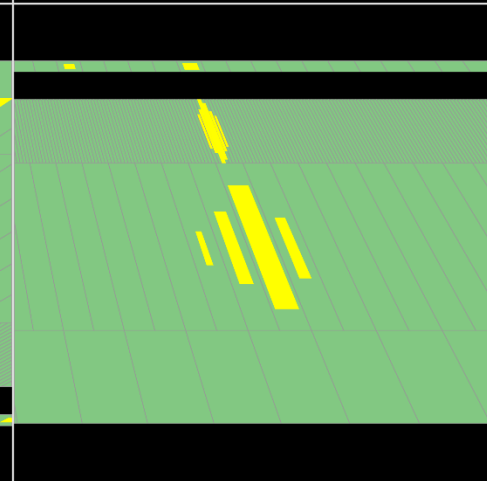
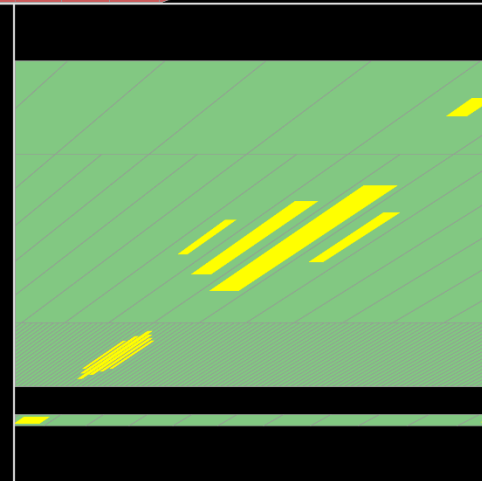
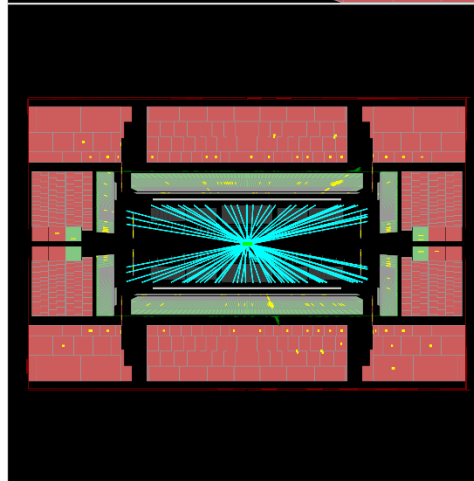
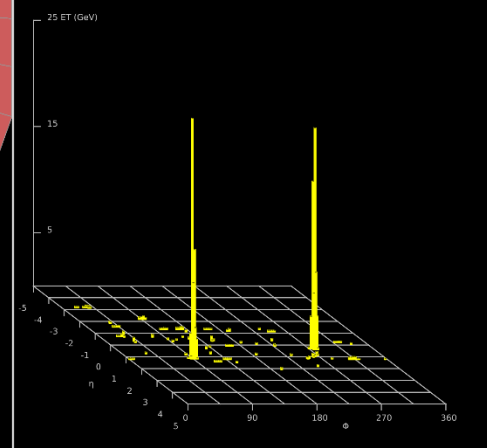
$H \rightarrow \gamma\gamma$

From RD-Schaffer's CERN seminar Sept 5th  
 And CONF-2017-047

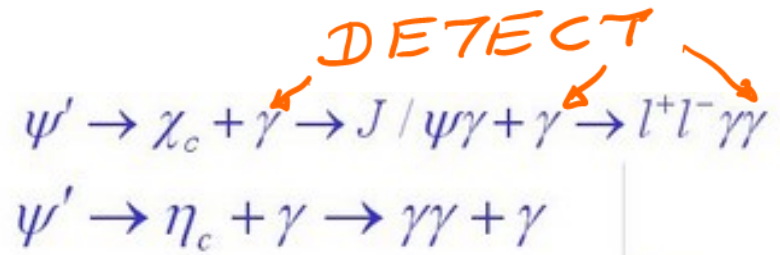
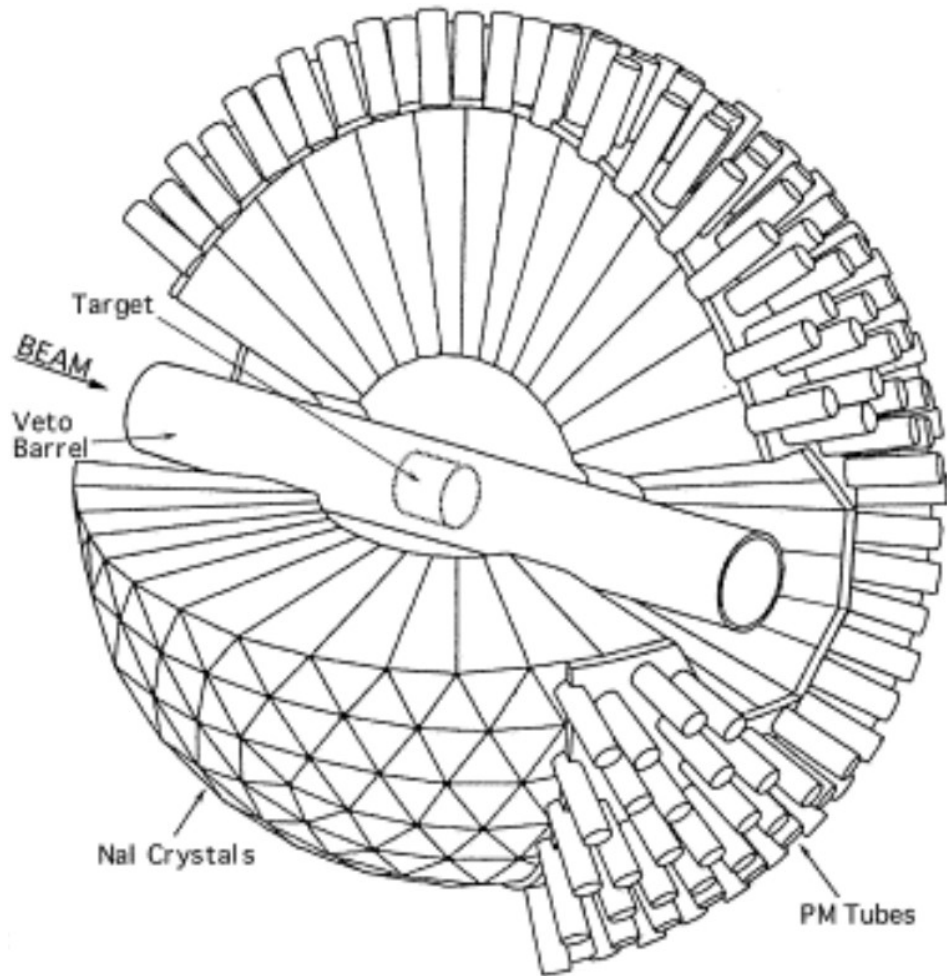


Run Number: 203779, Event Number: 56662314

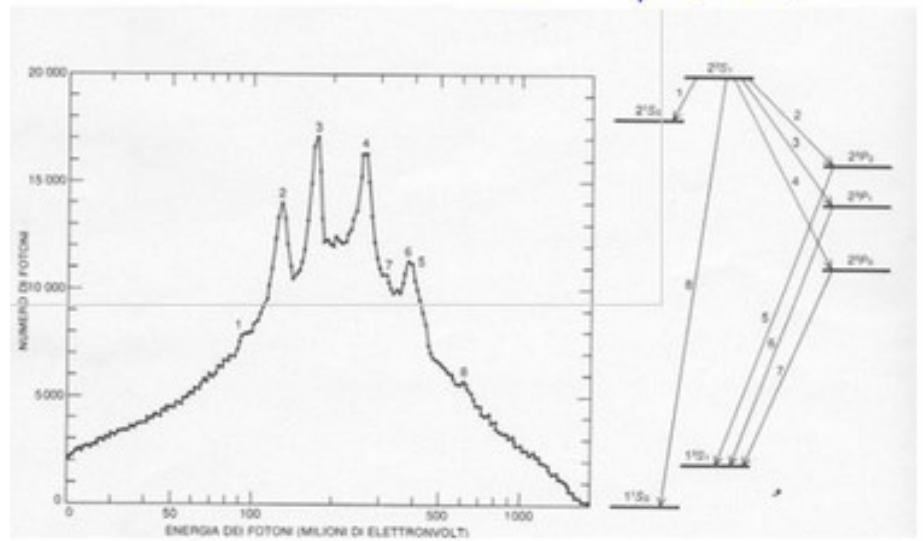
Date: 2012-05-23 22:19:29 CEST



# CRYSTAL BALL AT SPEAR $e^+e^-$

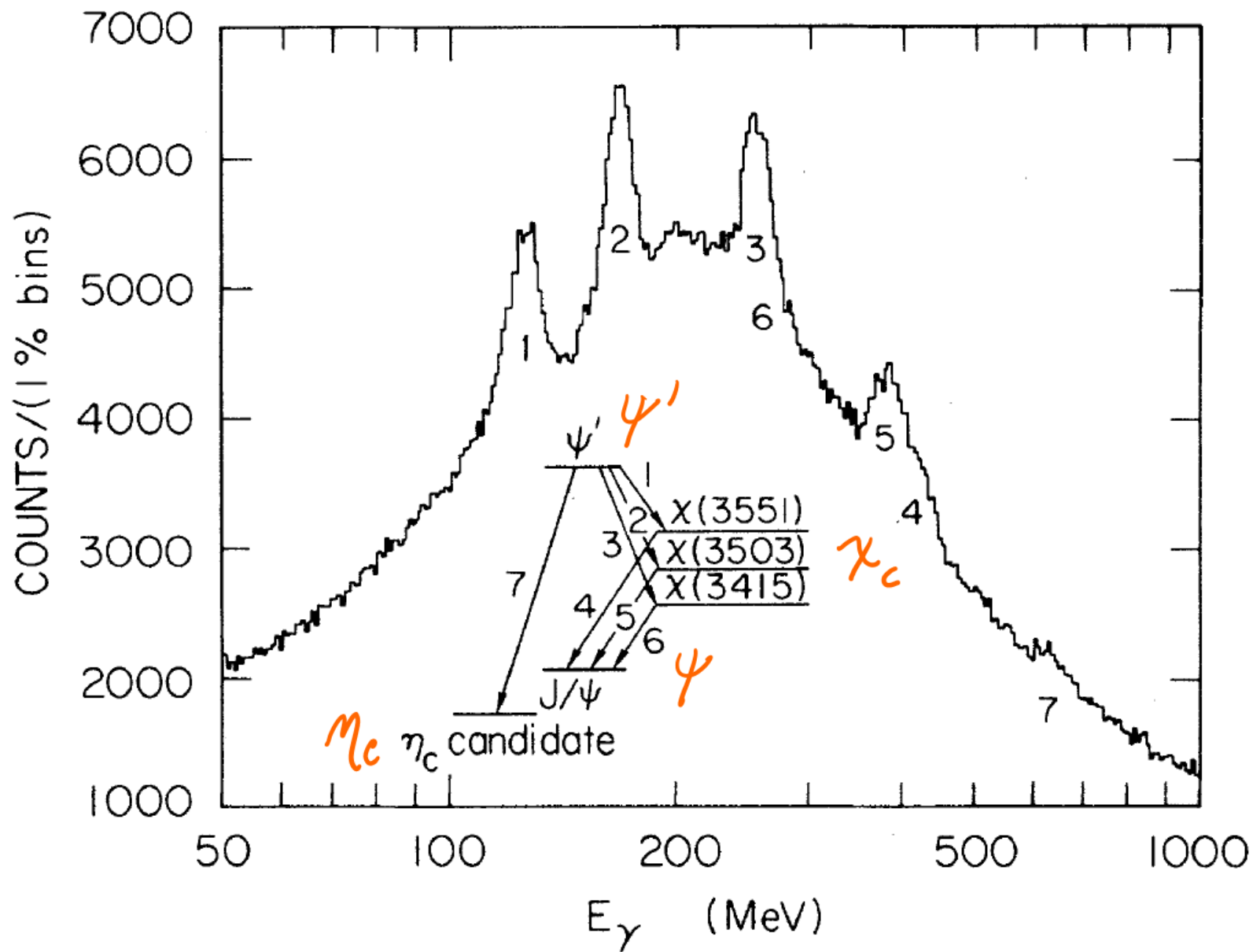


energy resolution:  $\frac{\sigma(E)}{E} = \frac{2.8\%}{\sqrt[4]{E(\text{GeV})}}$

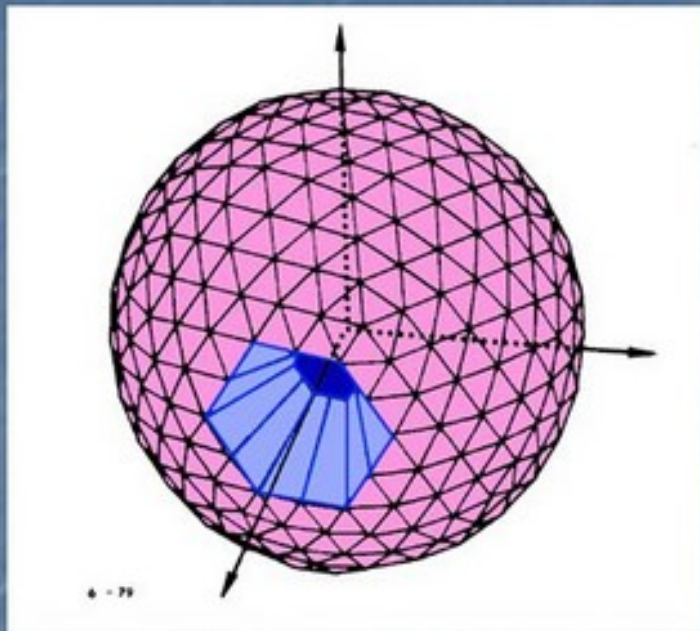


DISCOVERY of  $\chi_c$   $\eta_c$

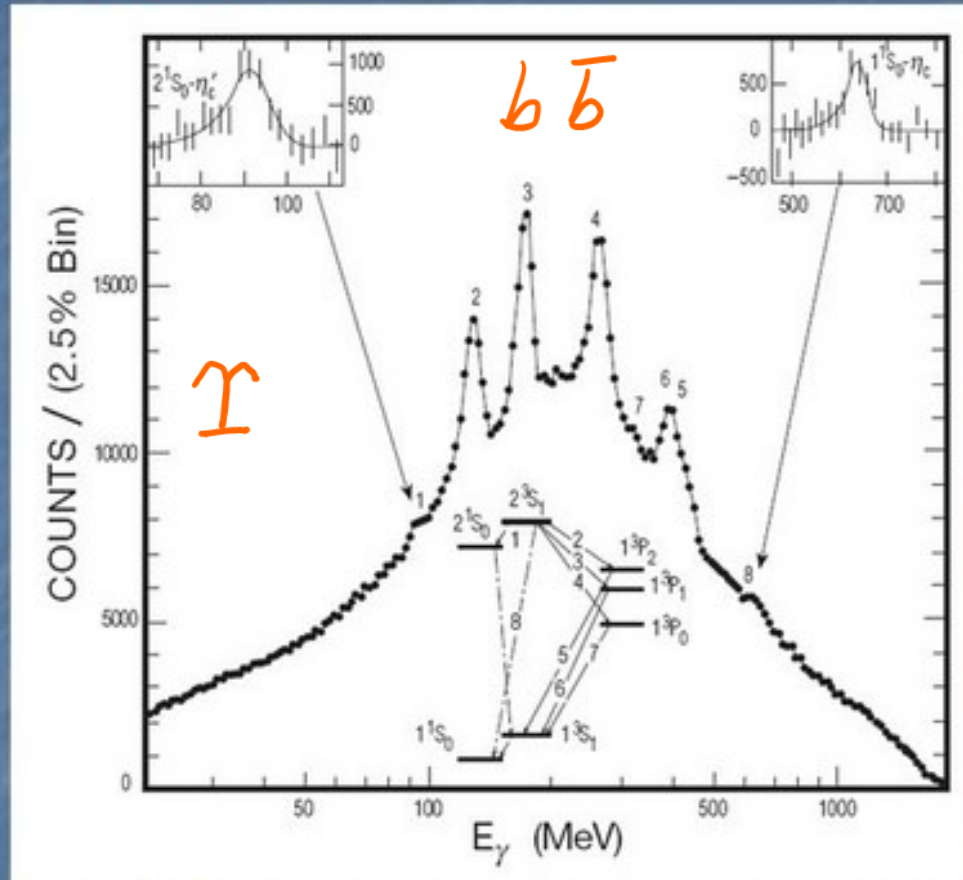




# Crystal Ball NaI(Tl) Calorimeter



Number of crystals 672  
 Inner radius 25.4 cm  
 Outer radius 66.0 cm  
 Thickness  $16 X_0$   
 Solid angle coverage 93%  
 Photodetector PMT  
 Noise 0.05 MeV  
 Dynamic range  $10^4$



$$\frac{\sigma_E}{E} = \frac{2.8\%}{\sqrt[4]{E(\text{GeV})}}$$

Inclusive photon spectrum at  $\psi(2S)$  resonance



**Evidence for a Narrow Massive State  
in the Radiative Decays of the Upsilon\***

July 1984  
(T/E)

Crystal Ball Collaboration

above (see Fig. 2b) now yields a significance of 4.2 standard deviations for the signal. The signal-parameters become

$$E_\gamma = (1072 \pm 8 \pm 21) \text{ MeV}$$

$$M = (8319 \pm 10 \pm 24) \text{ MeV}$$

$$\text{Counts} = 87.1 \pm 20.5$$

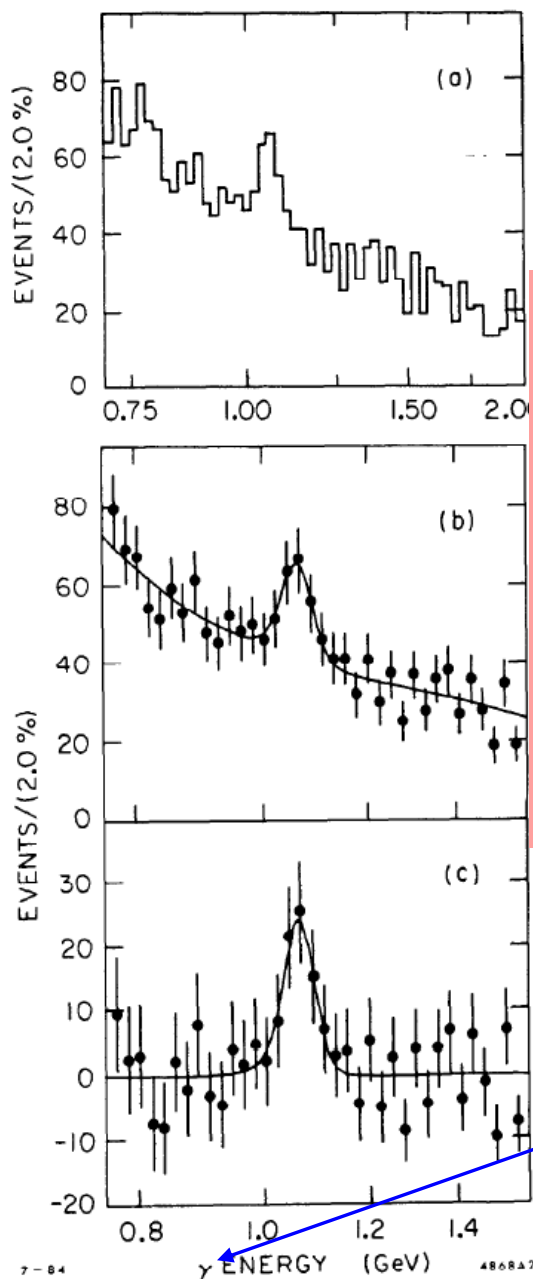
$$\chi^2 = 24.8 \text{ for 32 degrees of freedom,}$$

(1)

where the first error in  $E_\gamma$  or  $M$  is statistical and the second is systematic.<sup>(6)</sup>

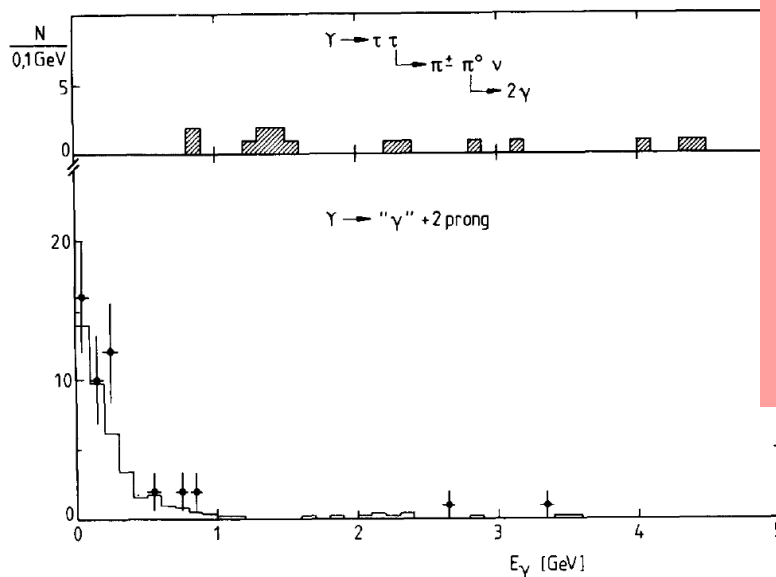
$$\Upsilon(1S) \rightarrow \gamma H^0$$

Is the Higgs mass 8.3 GeV ???  
At 4.2 sigma???



## SEARCH FOR NARROW STATES COUPLING TO $\tau$ PAIRS IN RADIATIVE $\Upsilon$ DECAYS

The ARGUS Collaboration



In summary, we have observed no indication for narrow objects produced in radiative  $\Upsilon$  decays and decaying into a  $\tau$  pair. The present sensitivity is an order of magnitude too small to check the predictions from the standard model, if only one scalar Higgs particle is assumed. However, the result puts improved constraints on models with a more complicated Higgs structure.

Is the Higgs mass 8.3 GeV ???  
At 4.2 sigma???

No, it isn't!

Fig. 3. Photon spectrum from the decay  $\Upsilon \rightarrow \gamma +$  two prongs used in search for the decay  $\Upsilon \rightarrow \gamma X$ ,  $X \rightarrow \tau^+ \tau^-$ . The hatched histogram shows the observed contribution from the decay  $\Upsilon \rightarrow \tau^+ \tau^-$  with one  $\tau$  decaying into  $\rho \nu$ , the open histogram shows expected background contributions from other sources as described in the text.