

# MAXWELL'S EQUATIONS

SUMMARIZE WHAT WE KNOW:

(i) $\nabla \cdot \vec{E} = \rho / \epsilon_0$	GAUSS	} EXPERIMENTAL FACTS
(ii) $\nabla \cdot \vec{B} = 0$		
(iii) $\nabla \times \vec{E} = -\partial \vec{B} / \partial t$	FARADAY	
(iv) $\nabla \times \vec{B} = \mu_0 \vec{J}$	AMPÈRE	

INHERENT INCONSISTENCY

FIRST TAKE DIVERGENCE OF  $\nabla \times \vec{E} = -\partial \vec{B} / \partial t$

$$\nabla \cdot (\nabla \times \vec{E}) = \nabla \cdot \left( -\frac{\partial \vec{B}}{\partial t} \right) = -\frac{\partial}{\partial t} (\nabla \cdot \vec{B})$$

DIV. CURL = 0      CONSISTENT       $\nabla \cdot \vec{B} = 0$

NOW TAKE DIVERGENCE OF  $\vec{\nabla} \times \vec{B} = \mu_0 \vec{J}$

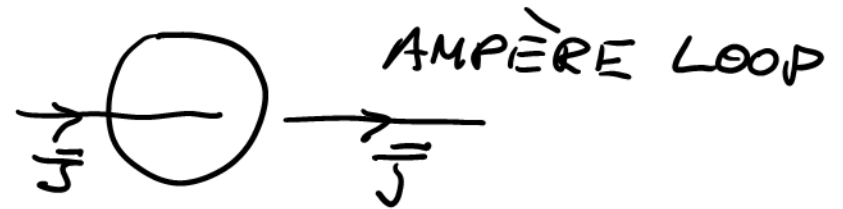
$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B}) = \mu_0 (\vec{\nabla} \cdot \vec{J})$$

$\underbrace{\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B})}_{\text{DIV CURL}} = 0$

$\rightarrow$  FOR STEADY CURRENT  
THIS IS ZERO

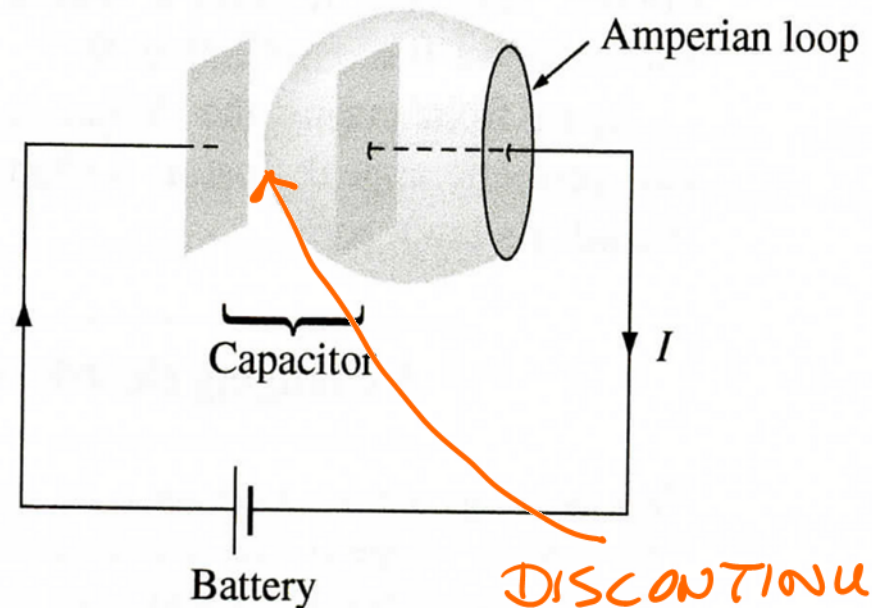
IF GO BEYOND MAGNETOSTATICS  
NOT STEADY CURRENT

$\rightarrow$  AMPERE CANNOT  
BE TRUE



CHARGE DOES NOT  
PILE UP

UP UNTIL NOW HAVE  
CONSIDERED CONTINUOUS  
CIRCUITS, CHARGE  
CANNOT PILE UP

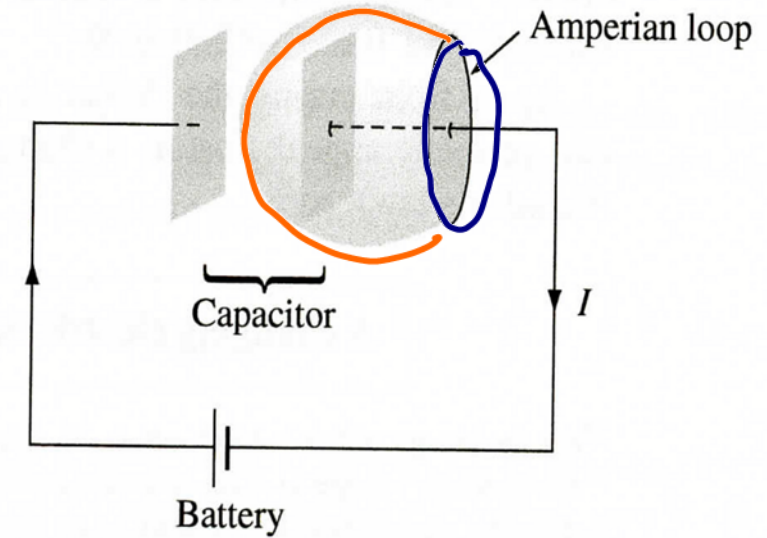


DISCONTINUOUS  
CIRCUIT

IN THIS CASE WE ARE CHARGING UP A CAPACITOR

IN GENERAL AMPÈRE SAYS

$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 I_{enc}$$



TOTAL CURRENT PASSING THROUGH LOOP

FOR SURFACE IN PLANE OF LOOP  $I_{en} = I$  OK ✓

FOR ORANGE SURFACE

↳ NO CURRENT PASSES THROUGH IT

SITUATIONS LIKE THIS

DO NOT ARISE IN MAGNETOSTATICS

$$I_{enc} = 0 \quad !$$

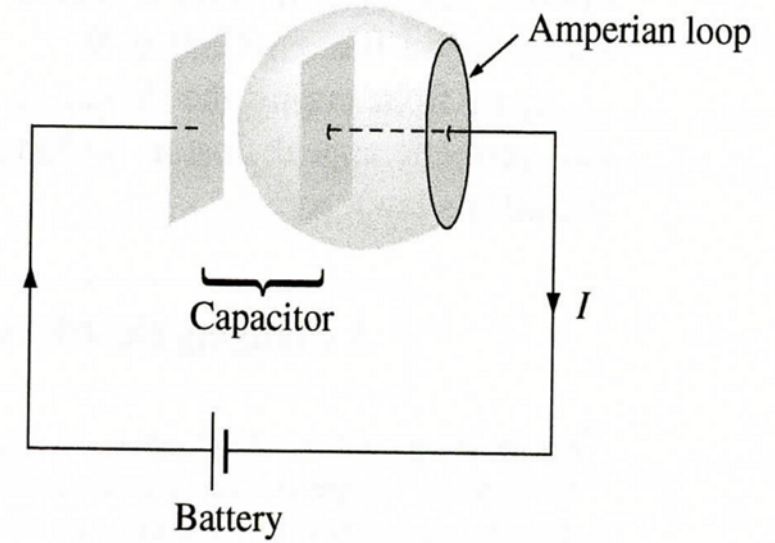
→ ONLY WHEN CHARGE PILING UP.

FOR NON-STEADY CURRENTS

→ CURRENT ENCLOSED BY LOOP

IS ILL-DEFINED

↳ DEPENDS ON SURFACE  
YOU CHOOSE



THIS IS A PURELY THEORETICAL PROBLEM

↳ AT THE TIME OF MAXWELL

THERE WAS NO EXPERIMENTAL REASON

TO DOUBT AMPÈRE'S LAW



GO BACK TO

$$\underbrace{\nabla \cdot (\nabla \times \vec{B})}_{\text{ZERO}} = \underbrace{\mu_0 (\nabla \cdot \vec{J})}_{\text{SHOULD BE ZERO BUT ISN'T.}}$$

SHOULD BE ZERO  
BUT ISN'T.

CONTINUITY EQUATION

$$\nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t}$$

GAUSS

$$\begin{aligned} -\frac{\partial \rho}{\partial t} &= -\frac{\partial}{\partial t} (\epsilon_0 \nabla \cdot \vec{E}) \\ &= -\nabla \cdot \left( \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right) \end{aligned}$$

IF

$$\vec{J} \rightarrow \vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\nabla \cdot \left( \vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right) = 0$$

$$\vec{J} \rightarrow \vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

REWRITE AMPÈRE AS!

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

THEN

$$\underbrace{\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B})}_0 = \underbrace{\vec{\nabla} \cdot (\mu_0 \vec{J})}_0 + \underbrace{\vec{\nabla} \cdot (\epsilon_0 \frac{\partial \vec{E}}{\partial t})}_0$$

THIS CHANGES NOTHING IN MAGNETOSTATICS

$$\vec{E} \text{ CONSTANT} \rightarrow \frac{\partial \vec{E}}{\partial t} = 0$$

EXTRA TERM HARD TO DETECT  $J \rightarrow$  LARGE  
 $\frac{\partial \vec{E}}{\partial t} \rightarrow$  SMALL

NOW:

$$\frac{\partial \vec{B}}{\partial t} \rightarrow \vec{E}$$

$$\frac{\partial \vec{E}}{\partial t} \rightarrow \vec{B}$$

CHANGING  $\vec{E}$  FIELD  
 RESULTS IN  $\vec{B}$  FIELD

MAXWELL CALLED EXTRA TERM DISPLACEMENT CURRENT

$$\vec{J}_d \equiv \epsilon_0 \partial \vec{E} / \partial t$$

GO BACK TO CHARGING CAPACITOR, PLATES CLOSE TOGETHER

ELECTRIC FIELD BETWEEN PLATES

GAUSS  $\rightarrow$   $\vec{E} = \frac{\sigma}{\epsilon_0}$   $\leftarrow$  SURFACE CHARGE DENSITY

$$\oint \vec{E} \cdot d\vec{a} = \frac{Q_{en}}{\epsilon_0}$$

$$EA = \frac{\sigma A}{\epsilon_0}$$

$$= \frac{1}{\epsilon_0} \frac{Q}{A}$$

$\leftarrow$  TOTAL CHARGE  
 $\leftarrow$  PLATE AREA

INTEGRAL FORM OF

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \partial \vec{E} / \partial t$$

BECOMES

$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 I_{enc} + \mu_0 \epsilon_0 \int \frac{\partial \vec{E}}{\partial t} \cdot d\vec{a}$$

$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 I_{enc} + \mu_0 \epsilon_0 \int \frac{\partial \vec{E}}{\partial t} \cdot d\vec{a}$$

FOR FLAT SURFACE

$$\vec{E} = 0 \text{ AND } I_{enc} = I$$

FOR ORANGE SURFACE

$$I_{enc} = 0$$

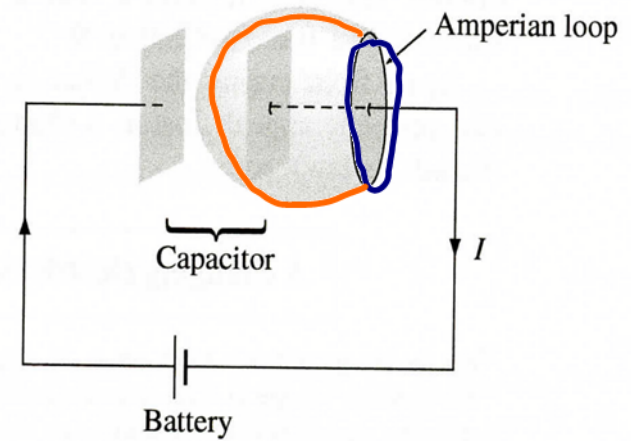
BUT 
$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 \epsilon_0 \int \frac{\partial \vec{E}}{\partial t} \cdot d\vec{a}$$

PUT 
$$\frac{\partial E}{\partial t} = \frac{I}{\epsilon_0 A}$$

$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 \epsilon_0 \int \frac{I}{\epsilon_0 A} \cdot d\vec{a}$$

$$= \frac{\mu_0 \epsilon_0}{\epsilon_0} \cdot \frac{I}{A} \int d\vec{a}$$

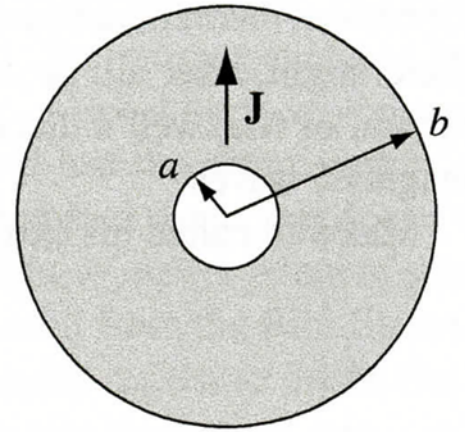
$$= \mu_0 I$$



COMES FROM  
CONDUCTION  
CURRENT

COMES FROM  
DISPLACEMENT CURRENT

EXAMPLE: TWO CONCENTRIC SPHERES  
 INNER, RADIUS  $a$  CARRIES A VARYING  
 CHARGE  $+Q(t)$ . OUTER, RADIUS  $b$ ,  
 CARRIES A CHARGE  $-Q(t)$ . THE SPACE  
 BETWEEN IS AN OHMIC MATERIAL, VOLUME  
 CONDUCTIVITY  $\sigma$



THERE IS A RADIAL (OHM'S LAW) CURRENT

$$\vec{J} = \sigma \vec{E} = \sigma \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \hat{r} \quad (1)$$

$$\begin{aligned} \vec{I} &= -\frac{dQ}{dt} = \int \vec{J} \cdot d\vec{a} \\ &= \frac{\sigma}{4\pi\epsilon_0} \cdot Q \int \frac{\hat{r}}{r^2} \end{aligned}$$

$$= \frac{\sigma}{4\pi\epsilon_0} \frac{Q}{R^2} \int r^2 \sin\theta \, d\theta \, d\phi$$

$$\vec{I} = \frac{\sigma Q}{4\pi\epsilon_0} \frac{1}{R^2} \cdot R^2 \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi = \frac{\sigma}{4\pi\epsilon_0} \cdot 4\pi$$

$$\vec{I} = \frac{\sigma Q}{\epsilon_0}$$

THIS GEOMETRY IS RADIALY SYMMETRIC  $\rightarrow \vec{B} = 0$

$\vec{B}$  WOULD BE RADIAL,  $\vec{\nabla} \cdot \vec{B} = 0 \rightarrow \oint \vec{B} \cdot d\vec{a} = \underbrace{B 4\pi r^2}_{=0}$

$Q, \vec{E}, \vec{J}$  ARE ALL FUNCTIONS OF TIME  $\therefore \vec{B} = 0$

AMPÈRE & BIOT-SAVART DO NOT APPLY

$\rightarrow$  USE DISPLACEMENT CURRENT

## DISPLACEMENT CURRENT

$$J_d = \epsilon_0 \frac{\partial E}{\partial t} = \frac{1}{4\pi\epsilon_0 r^2} \frac{dQ}{dt} \cdot \hat{r} \rightarrow = -\frac{\sigma Q}{\epsilon_0}$$

$$J_d = -\sigma \frac{Q}{4\pi\epsilon_0 r^2} \cdot \hat{r} \quad (2)$$

$$(1) - (2) = 0 = J_{\text{OHM}} - J_{\text{DISP}} \rightarrow \text{NO CURRENT} \\ \text{NO } \vec{B}$$

ALSO

$$\vec{\nabla} \times \vec{B} = \mu_0 J - \mu_0 \epsilon_0 \frac{\partial E}{\partial t} = 0$$



WE NOW HAVE

$$(i) \quad \vec{\nabla} \cdot \vec{E} = \rho / \epsilon_0$$

$$(ii) \quad \vec{\nabla} \cdot \vec{B} = 0$$

$$(iii) \quad \vec{\nabla} \times \vec{E} = -\partial \vec{B} / \partial t$$

$$(iv) \quad \vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$+ \quad \vec{F} = q (\vec{E} + \vec{v} \times \vec{B})$$

THIS IS ALL OF CLASSICAL ELECTROMAGNETISM

EVEN CONTINUITY EQUATION IS IN THERE

$$\underbrace{\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B})}_0 = \mu_0 \vec{\nabla} \cdot \vec{J} + \mu_0 \epsilon_0 \vec{\nabla} \cdot \partial \vec{E} / \partial t$$

$$\mu_0 \vec{\nabla} \cdot \vec{J} = -\mu_0 \epsilon_0 \cdot \frac{1}{\epsilon_0} \frac{\partial \rho}{\partial t} \Rightarrow \vec{\nabla} \cdot \vec{J} = -\frac{\partial \rho}{\partial t}$$

USUAL WAY OF WRITING MAXWELL SEEMS TO IMPLY

$\vec{E}$  PRODUCED BY  $\rho$  OR  $\partial\vec{B}/\partial t$

$\vec{B}$  PRODUCED BY  $\vec{J}$  OR  $\partial\vec{E}/\partial t$ .

BUT  $\partial\vec{B}/\partial t$ ,  $\partial\vec{E}/\partial t$  THEMSELVES PRODUCED BY  $\rho$ ,  $\vec{J}$   
MORE LOGICAL (?)

$$\vec{\nabla} \cdot \vec{E} = \rho/\epsilon_0, \quad \vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0$$

$$\vec{\nabla} \cdot \vec{B} = 0, \quad \vec{\nabla} \times \vec{B} - \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} = \mu_0 \vec{J}$$

FIELDS PRODUCED BY SOURCES

ALL OF  $E$  &  $H$  COMES FROM CURRENTS & CHARGES

MAXWELL  $\rightarrow$  SOURCES PRODUCE FIELDS

LORENTZ  $\rightarrow$  HOW FIELDS EFFECT CHARGES

# MAGNETIC CHARGES?

STRIKING SYMMETRY IN MAXWELL  $\begin{matrix} \rho \rightarrow 0 \\ \vec{J} \rightarrow 0 \end{matrix}$

$$\textcircled{1} \quad \vec{\nabla} \cdot \vec{E} = 0 \quad \vec{\nabla} \times \vec{E} = -\partial \vec{B} / \partial t$$

$$\textcircled{2} \quad \vec{\nabla} \cdot \vec{B} = 0 \quad \vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \partial \vec{E} / \partial t$$

IF  $\vec{E} \rightarrow \vec{B}$ ,  $\vec{B} \rightarrow -\mu_0 \epsilon_0 \vec{E}$   $\textcircled{1} \rightarrow \textcircled{2}$

SYMMETRY SPOILED BY  $\vec{\nabla} \cdot \vec{E} = \rho / \epsilon_0$ ,  $\vec{\nabla} \times \vec{B} = \mu_0 \vec{J}$

$$\vec{\nabla} \cdot \vec{E} = \rho_e / \epsilon_0 \quad \vec{\nabla} \times \vec{E} = -\mu_0 \vec{J}_m - \partial \vec{B} / \partial t$$

$$\vec{\nabla} \cdot \vec{B} = \mu_0 \rho_m \quad \vec{\nabla} \times \vec{B} = \mu_0 \vec{J}_e + \mu_0 \epsilon_0 \partial \vec{E} / \partial t$$

WHAT IF HAD

$\rho_e \leftrightarrow \rho_m \rightarrow$  BOTH CHARGES CONSERVED

$$\vec{\nabla} \cdot \vec{J}_m = -\frac{\partial \rho_m}{\partial t} \quad ; \quad \vec{\nabla} \cdot \vec{J}_e = -\frac{\partial \rho_e}{\partial t}$$

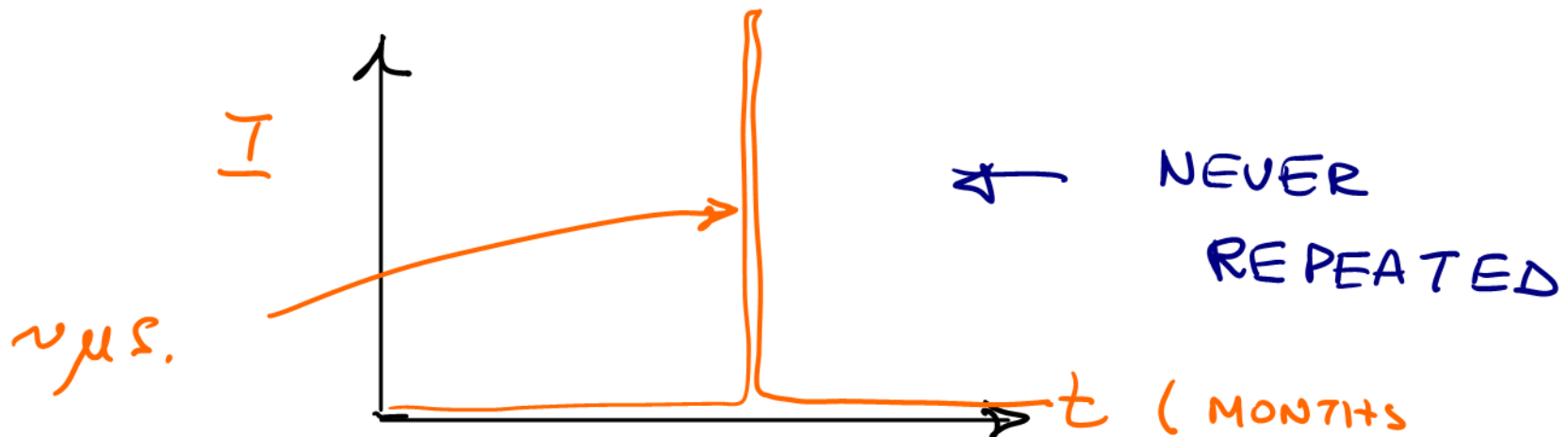
THERE APPEAR TO BE NO MAGNETIC MONOPOLES  
(CHARGES) IN OUR UNIVERSE

→ DIRAC SHOWED MAGNETIC MONOPOLES → ELECTRIC CHARGE  
QUANTIZATION

→ EXPECT MONOPOLES PRODUCED IN BIG BANG

↳ INFLATION SPREAD THEM OUT  $\sim 1$  IN OUR  
UNIVERSE

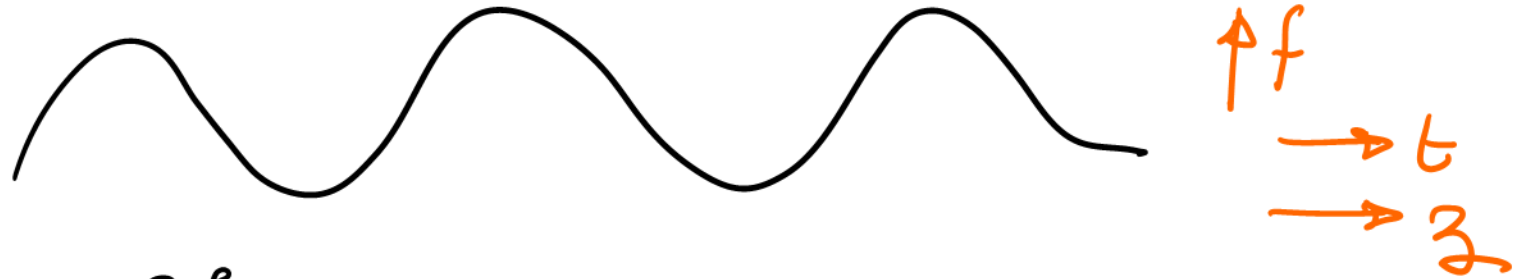
ONE EXPERIMENT CLAIMED TO HAVE OBSERVED



# ELECTROMAGNETIC WAVES

## 1-D WAVES

STRING



$$\frac{\partial^2 f}{\partial z^2} = \frac{1}{v^2} \frac{\partial^2 f}{\partial t^2}$$

$$f = f(z, t)$$

$$f(z, t) = A \cos [k(z - vt) + \delta]$$

Annotations: An arrow points from 'A' to 'AMPLITUDE'. An arrow points from the bracketed term to 'PHASE'. An arrow points from 'v' to 'VELOCITY'.

AMPLITUDE

$$f(z, t) = \text{Re} \left[ A e^{i(kz - \omega t + \delta)} \right]$$

Annotation: An arrow points from the exponent to 'FREQUENCY'.

FREQUENCY

$$v = \sqrt{\frac{T}{m}}$$

Annotations: An arrow points from 'v' to 'VELOCITY'. An arrow points from 'T' to 'TENSION'. An arrow points from 'm' to 'MASS PER UNIT LENGTH'.

# MAXWELL IN FREE SPACE NO CHARGE OR CURRENT

$$(i) \quad \bar{\nabla} \cdot \bar{E} = 0$$

$$(iii) \quad \bar{\nabla} \times \bar{E} = -\partial \bar{B} / \partial t$$

$$(ii) \quad \bar{\nabla} \cdot \bar{B} = 0$$

$$(iv) \quad \bar{\nabla} \times \bar{B} = \mu_0 \epsilon_0 \partial \bar{E} / \partial t$$

COUPLED 1<sup>ST</sup> ORDER DIFFERENTIAL EQUATIONS

DECOUPLE BY TAKING CURL OF (iii) (iv)

$$\begin{aligned} \bar{\nabla} \times (\bar{\nabla} \times \bar{E}) &= \bar{\nabla} (\bar{\nabla} \cdot \bar{E}) - \bar{\nabla}^2 \bar{E} = \bar{\nabla} \times \left( -\frac{\partial \bar{B}}{\partial t} \right) \\ &= -\frac{\partial}{\partial t} (\bar{\nabla} \times \bar{B}) = -\mu_0 \epsilon_0 \frac{\partial^2 \bar{E}}{\partial t^2} \end{aligned}$$

SINCE

$$\bar{\nabla} \cdot \bar{E} = 0$$

$$\nabla^2 \bar{E} = \mu_0 \epsilon_0 \frac{\partial^2 \bar{E}}{\partial t^2}$$

$$\begin{aligned}\bar{\nabla} \times (\bar{\nabla} \times \bar{B}) &= \bar{\nabla} (\bar{\nabla} \cdot \bar{B}) - \nabla^2 \bar{B} = \bar{\nabla} \times \left( \mu_0 \epsilon_0 \frac{\partial \bar{E}}{\partial t} \right) \\ &= \mu_0 \epsilon_0 \frac{\partial}{\partial t} (\bar{\nabla} \times \bar{E}) = -\mu_0 \epsilon_0 \partial^2 \bar{B} / \partial t^2\end{aligned}$$

$$\nabla^2 \bar{B} = \mu_0 \epsilon_0 \partial^2 \bar{B} / \partial t^2$$

AND HAD

$$\nabla^2 \bar{E} = \mu_0 \epsilon_0 \partial^2 \bar{E} / \partial t^2$$

3-D WAVE EQUATION  $\frac{\partial^2 f}{\partial t^2} = \frac{1}{v^2} \frac{\partial^2 f}{\partial t^2} \Rightarrow \nabla^2 f = \frac{1}{v^2} \frac{\partial^2 f}{\partial t^2}$

WAVE EQUATIONS FOR  $\bar{E}$  &  $\bar{B}$

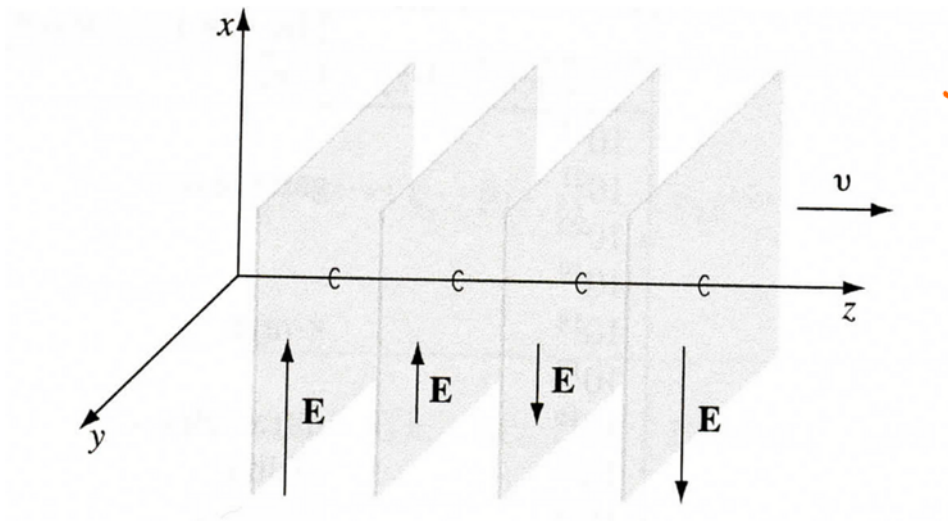
$$v = \frac{1}{\sqrt{\epsilon_0 \mu_0}} = 3 \times 10^8 \text{ m/sec}$$

WITHOUT  $\epsilon_0 \mu_0 \frac{\partial \bar{E}}{\partial t} \rightarrow$  NO WAVES

$\epsilon_0 \partial \bar{E} / \partial t \rightarrow$  DISPLACEMENT CURRENT



# MONOCHROMATIC PLANE WAVES



→  $v_z$

NO  $x, y$  DEPENDENCE

FIXED FREQUENCY  $\omega$

GENERALLY  $f(z, t) = \text{Re} \left[ A e^{i(kz - \omega t + \delta)} \right]$

IN OUR CASE

COMPLEX →  $\vec{E} = \vec{E}_0 e^{i(kz - \omega t)}$   
 $\vec{B} = \vec{B}_0 e^{i(kz - \omega t)}$

$$\vec{E} = \vec{E}_0 e^{i(kz - \omega t)}$$
$$\vec{B} = \vec{B}_0 e^{i(kz - \omega t)}$$

EVERY SOLUTION OF MAXWELL IS A WAVE

EVERY WAVE IS NOT A SOLUTION OF MAXWELL

MAXWELL  $\rightarrow$   $\vec{\nabla} \cdot \vec{E} = 0, \quad \vec{\nabla} \cdot \vec{B} = 0$

$$(\vec{E}_0)_z = (\vec{B}_0)_z = 0$$

ELECTROMAGNETIC WAVES ARE TRANSVERSE

NO COMPONENT OF FIELDS IN DIRECTION OF MOTION

ALSO  $\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t}$   $\bar{E}, \bar{B}$  AMPLITUDES ARE RELATED

$$-k(\tilde{E}_0)_y = -\omega(\tilde{B}_0)_x, \quad k(\tilde{E}_0)_x = \omega(\tilde{B}_0)_y$$

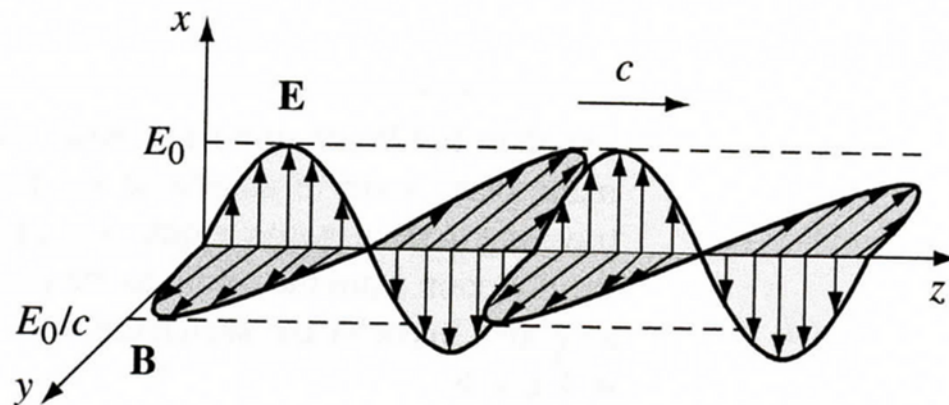
$$\tilde{B}_0 = \frac{k}{\omega} (\hat{z} \times \tilde{E}_0)$$

$\bar{E}$  &  $\bar{B}$  IN PHASE & PERPENDICULAR

AMPLITUDES RELATED  $B_0 = \frac{k}{\omega} E_0 = \frac{1}{c} E_0$

$$\mathbf{E}(z, t) = E_0 \cos(kz - \omega t + \delta) \hat{x}, \quad \mathbf{B}(z, t) = \frac{1}{c} E_0 \cos(kz - \omega t + \delta) \hat{y}.$$

(9.48)



<b>The Electromagnetic Spectrum</b>		
Frequency (Hz)	Type	Wavelength (m)
$10^{22}$		$10^{-13}$
$10^{21}$	gamma rays	$10^{-12}$
$10^{20}$		$10^{-11}$
$10^{19}$		$10^{-10}$
$10^{18}$	x-rays	$10^{-9}$
$10^{17}$		$10^{-8}$
$10^{16}$	ultraviolet	$10^{-7}$
$10^{15}$	visible	$10^{-6}$
$10^{14}$	infrared	$10^{-5}$
$10^{13}$		$10^{-4}$
$10^{12}$		$10^{-3}$
$10^{11}$		$10^{-2}$
$10^{10}$	microwave	$10^{-1}$
$10^9$		1
$10^8$	TV, FM	10
$10^7$		$10^2$
$10^6$	AM	$10^3$
$10^5$		$10^4$
$10^4$	RF	$10^5$
$10^3$		$10^6$

<b>The Visible Range</b>		
Frequency (Hz)	Color	Wavelength (m)
$1.0 \times 10^{15}$	near ultraviolet	$3.0 \times 10^{-7}$
$7.5 \times 10^{14}$	shortest visible blue	$4.0 \times 10^{-7}$
$6.5 \times 10^{14}$	blue	$4.6 \times 10^{-7}$
$5.6 \times 10^{14}$	green	$5.4 \times 10^{-7}$
$5.1 \times 10^{14}$	yellow	$5.9 \times 10^{-7}$
$4.9 \times 10^{14}$	orange	$6.1 \times 10^{-7}$
$3.9 \times 10^{14}$	longest visible red	$7.6 \times 10^{-7}$
$3.0 \times 10^{14}$	near infrared	$1.0 \times 10^{-6}$

NOW YOU HAVE TOOLS TO STUDY MANY FIELDS  
OF PHYSICS

→ OPTICS

→ PLASMAS

→ ELECTRICITY GENERATION

→ EM WAVES IN TECHNOLOGY

→ EM WAVES IN ASTROPHYSICS

⋮