

BOUNDARY CONDITIONS

TYPICALLY — KNOW ρ
WANT \vec{E}

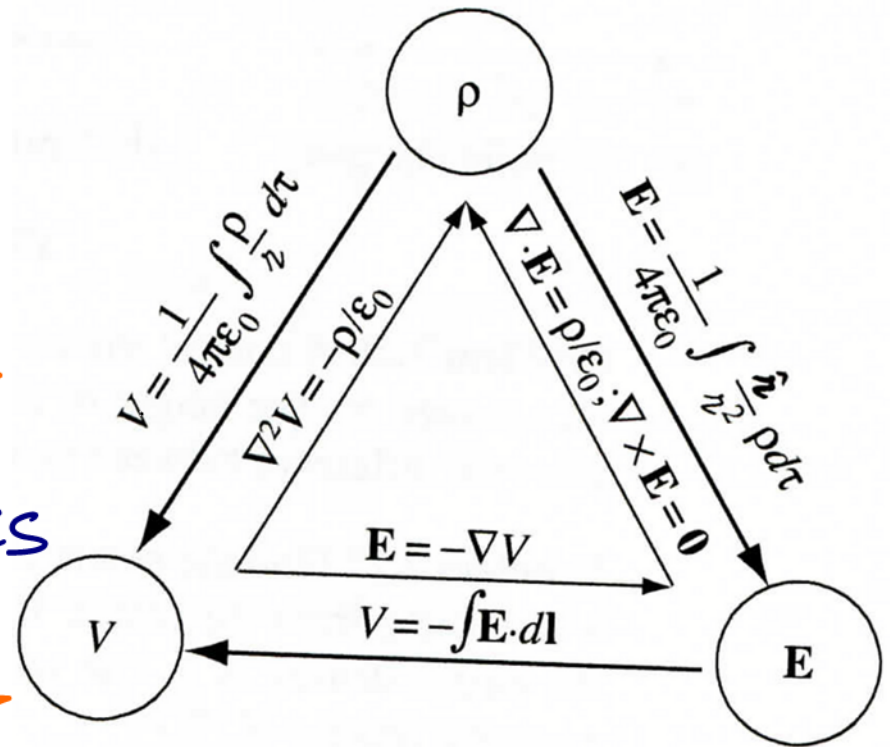
IF NO SYMMETRY

↳ BEST FIND POTENTIAL

3 FUNDAMENTAL QUANTITIES
IN ELECTROSTATICS

ρ , \vec{E} , V

SIX FORMULAS RELATING THEM



THESE FORMULAE COME FROM 2 EXPERIMENTAL
FACTS

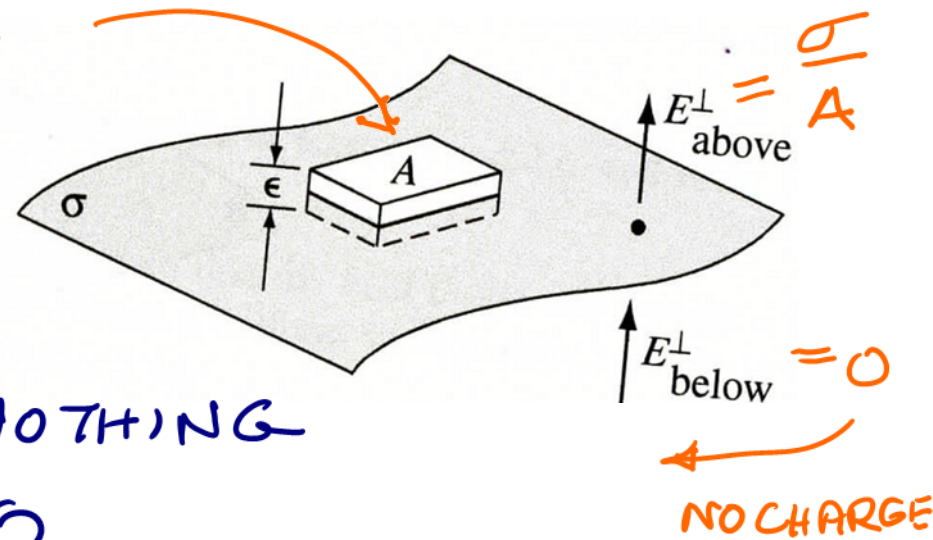
- 1) SUPERPOSITION
- 2) COULOMBS LAW

ELECTRIC FIELD ALWAYS UNDERGOES DISCONTINUITY
WHEN YOU CROSS A SURFACE CHARGE

CAN FIND AMOUNT OF CHANGE IN E AT BOUNDARY

VERY THIN GAUSSIAN PILL BOX
GAUSS'S LAW

$$\oint \vec{E} \cdot d\vec{a} = \frac{Q_{\text{ENC}}}{\epsilon_0} = \frac{\sigma A}{\epsilon_0}$$



SIDES OF BOX CONTRIBUTE NOTHING
AS THICKNESS OF BOX $\rightarrow 0$

$$E_{\text{ABOVE}}^{\perp} - E_{\text{BELOW}}^{\perp} = \frac{\rho}{\epsilon_0}$$

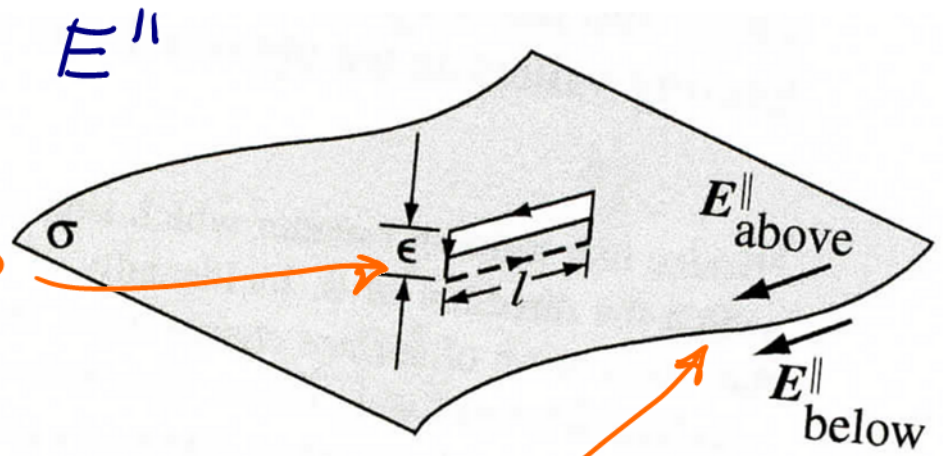
NORMAL COMPONENT OF \vec{E} IS DISCONTINUOUS BY
AN AMOUNT σ/ϵ_0 AT ANY BOUNDARY

\rightarrow NO SURFACE CHARGE E^{\perp} CONTINUOUS

TANGENTIAL COMPONENT E_{\parallel}
IS ALWAYS CONTINUOUS

THIN RECTANGULAR LOOP

$$\oint \vec{E} \cdot d\vec{l} = 0$$



ENDS GIVE NOTHING, SIDES $E_{\parallel \text{ABOVE}} = E_{\parallel \text{BELOW}}$

BOUNDARY CONDITION COMBINED

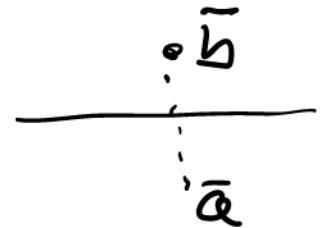
$$\vec{E}_{\text{ABOVE}} - \vec{E}_{\text{BELOW}} = \frac{\rho}{\epsilon_0} \hat{n}$$

POTENTIAL IS CONTINUOUS ACROSS BOUNDARY

$$V_{\text{ABOVE}} - V_{\text{BELOW}} = - \int_a^b \vec{E} \cdot d\vec{l}$$

AS $dl \rightarrow 0$

$V_{\text{ABOVE}} \rightarrow V_{\text{BELOW}}$



$$V_{\text{ABOVE}} = V_{\text{BELOW}}$$

GRADIENT OF V DISCONTINUOUS $\rightarrow \vec{E}$

$$\vec{E} = -\vec{\nabla} V$$

$$\vec{\nabla} V_{\text{ABOVE}} - \vec{\nabla} V_{\text{BELOW}} = -\frac{1}{\epsilon_0} \sigma \hat{n}$$

WRITE $\frac{\partial V}{\partial n} = \vec{\nabla} V \cdot \hat{n}$

NORMAL DERIVATIVE
 \hookrightarrow RATE OF CHANGE
NORMAL TO
SURFACE

$$\frac{\partial V_{\text{ABOVE}}}{\partial n} - \frac{\partial V_{\text{BELOW}}}{\partial n} = -\frac{1}{\epsilon_0} \sigma$$

DERIVATIVES ARE LIMITING VALUES

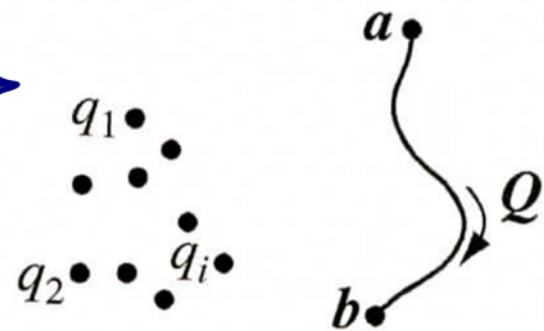
\hookrightarrow INFINITESIMALLY CLOSE TO THE
BOUNDARY SURFACE

WORK & ENERGY IN ELECTROSTATICS

WORK TO MOVE A CHARGE

HOW MUCH WORK DOES IT TAKE TO MOVE FROM $a \rightarrow b$

FORCE ON CHARGE $\vec{F} = Q\vec{E}$



WORKING AGAINST FORCE $-Q\vec{E}$

$$\text{WORK} = \int_{\vec{a}}^{\vec{b}} \vec{F} \cdot d\vec{e} = -Q \int_{\vec{a}}^{\vec{b}} \vec{E} \cdot d\vec{e} = Q [V(\vec{b}) - V(\vec{a})]$$

POTENTIAL

WORK DONE \rightarrow INDEP of PATH \rightarrow CONSERVATIVE FORCE

$$V(\vec{b}) - V(\vec{a}) = W/Q$$

POTENTIAL DIFF BETWEEN \vec{a} , \vec{b} IS WORK PER UNIT CHARGE $\vec{a} \rightarrow \vec{b}$. TO BRING CHARGE FROM

FAR AWAY $\rightarrow \infty$ $W = Q [V(\vec{r}) - V(\infty)]$

$$W = Q [V(r) - V(\infty)]$$

IF HAVE SET REFERENCE POINT AT ∞

$$V(\infty) = 0$$

$$W = QV(r)$$

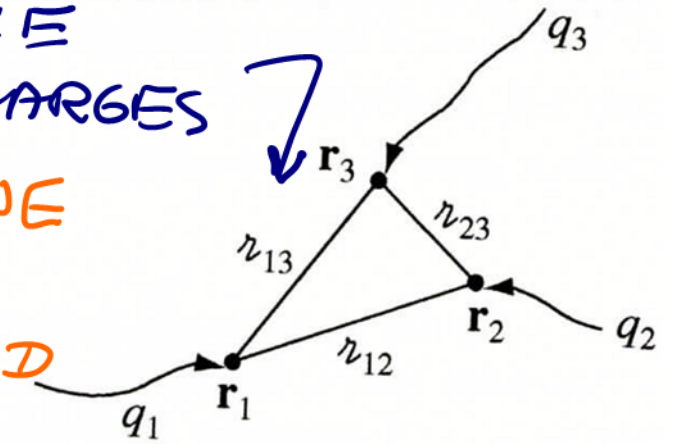
POTENTIAL IS POTENTIAL ENERGY PER UNIT CHARGE

→ WORK IT TAKES TO ASSEMBLE SYSTEM

REMEMBER → FIELD IS FORCE / UNIT CHARGE

ENERGY OF A POINT CHARGE DISTRIBUTION

HOW MUCH WORK DOES IT TAKE
TO ASSEMBLE A GROUP OF CHARGES
BRING THEM FROM ∞ ONE BY ONE



- q_1 FIRST CHARGE - NO WORK NO FIELD

- q_2 $W = QV(\vec{r}) \rightarrow W_2 = \frac{1}{4\pi\epsilon_0} q_2 \left(\frac{q_1}{r_{12}} \right)$

$$W_3 = \frac{1}{4\pi\epsilon_0} q_3 \left(\frac{q_1}{r_{13}} + \frac{q_2}{r_{23}} \right)$$

$$W_4 = \frac{1}{4\pi\epsilon_0} q_4 \left(\frac{q_1}{r_{14}} + \frac{q_2}{r_{24}} + \frac{q_3}{r_{34}} \right)$$

TOTAL WORK = $\frac{1}{4\pi\epsilon_0} \left(\frac{q_1 q_2}{r_{12}} + \frac{q_1 q_3}{r_{13}} + \frac{q_1 q_4}{r_{14}} + \frac{q_2 q_3}{r_{23}} + \frac{q_2 q_4}{r_{24}} + \frac{q_3 q_4}{r_{34}} \right)$

$$W = \frac{1}{4\pi\epsilon_0} \sum_{i=1}^n \sum_{j>i}^n \frac{q_i q_j}{r_{ij}}$$

DOESN'T COUNT EACH PAIR TWICE

DOES COUNT EACH PAIR TWICE

$$W = \frac{1}{8\pi\epsilon_0} \sum_{i=1}^n \sum_{j \neq i}^n \frac{q_i q_j}{r_{ij}}$$

$$W = \frac{1}{2} \sum_{i=1}^n q_i \left(\sum_{j \neq i}^n \frac{1}{4\pi\epsilon_0} \frac{q_j}{r_{ij}} \right)$$

POTENTIAL AT

q_i DUE TO ALL OTHER CHARGES

$$W = \frac{1}{2} \sum_{i=1}^n q_i V(\vec{r}_i)$$

YOU GET THIS WORK BACK IF YOU DISMANTLE ASSEMBLY

ENERGY STORED IN CONFIGURATION OF CHARGES

ENERGY OF CONTINUOUS CHARGE DISTRIBUTION

$$W = \frac{1}{2} \sum_{i=1}^n q_i V(\vec{r}_i) \rightarrow \frac{1}{2} \int \rho V d\tau$$

← VOLUME CHARGE DENSITY

USE GAUSS LAW $\rho \rightarrow \vec{E}$

$$\rho = \epsilon_0 \vec{\nabla} \cdot \vec{E} \rightarrow W = \frac{\epsilon_0}{2} \int (\vec{\nabla} \cdot \vec{E}) V d\tau$$

$$\int \text{BY PARTS} \rightarrow \int_V f(\vec{\nabla} \cdot \vec{A}) d\tau = - \int_V \vec{A} \cdot (\vec{\nabla} f) d\tau + \oint f \vec{A} \cdot d\vec{a}$$

$$W = \frac{\epsilon_0}{2} \left[- \int \vec{E} \cdot (\vec{\nabla} V) d\tau + \oint V \vec{E} \cdot d\vec{a} \right]$$

$$\text{BUT } \vec{\nabla} V = -\vec{E}$$

$$W = \frac{\epsilon_0}{2} \int_V E^2 d\tau + \oint_S V \vec{E} \cdot d\vec{a}$$

$$W = \frac{\epsilon_0}{2} \int_V E^2 d\tau + \oint_S V \vec{E} \cdot d\vec{a}$$

THIS CAME FROM

$$\frac{1}{2} \int \rho V d\tau$$

∫ OVER LARGER SPACE W DOES NOT CHANGE

$\rho = 0$
OUTSIDE
VOLUME

AS INCREASE VOLUME $\int_V \vec{E}$ INCREASES

\int_S DECREASES → W CONSTANT

W = ∫ CORRECT FOR ANY VOLUME

$$E \sim \frac{1}{r^2}, \quad V \sim \frac{1}{r}, \quad \text{SURFACE} \sim r^2$$

∫ OVER ALL SPACE → SURFACE → 0

$$W = \frac{\epsilon_0}{2} \int_{\text{ALL SPACE}} E^2 d\tau$$

FIND ENERGY OF UNIFORMLY CHARGED SPHERICAL SHELL

① USE $W = \frac{1}{2} \int \sigma V da = \frac{1}{2} \int \sigma \cdot \frac{1}{4\pi\epsilon_0} \frac{q}{R} da$

$$= \frac{1}{8\pi\epsilon_0} \frac{q}{R} \int \sigma da = \frac{1}{8\pi\epsilon_0} \frac{q^2}{R}$$

OR ②

$$W = \frac{\epsilon_0}{2} \int_{\text{ALL SPACE}} E^2 d\tau$$

INSIDE SPHERE $\vec{E} = \vec{0}$

OUTSIDE $\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r} \rightarrow E^2 = \frac{q^2}{(4\pi\epsilon_0)^2} \cdot \frac{1}{r^4}$

$$W = \frac{\epsilon_0}{2(4\pi\epsilon_0)^2} \int_{\text{OUTSIDE}} \left(\frac{q^2}{r^4} \right) (r^2 \sin\theta dr d\theta d\phi)$$

SPHERICAL COORDS

$$= \frac{1}{32\pi^2\epsilon_0} \cdot \frac{q^2}{r^2} \cdot 4\pi \int_R^{\infty} \frac{1}{r^2} dr = \frac{1}{8\pi\epsilon_0} \frac{q^2}{R}$$

SOME COMMENTS ON ES ENERGY

$$W = \frac{\epsilon_0}{2} \int E^2 \tau \rightarrow \text{IMPLIES } W \text{ IS ALWAYS +VE}$$

$$\text{BUT } W = \frac{1}{2} \sum_{i=1}^n q_i V(\vec{r}_i) \quad \text{COULD BE -VE}$$

$q_i \rightarrow \pm$

↳ DOES NOT TAKE INTO ACCOUNT
WORK TO CREATE CHARGE

↳ JUST WORK TO BRING THEM
TOGETHER

$$W = \frac{\epsilon_0}{2} \int E^2 d\tau \quad \text{IMPLIES ENERGY OF POINT
CHARGE IS } \infty$$

$$W = \frac{\epsilon_0}{2} \int \left(\frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \right)^2 r^2 \sin\theta \, dr \, d\theta \, d\phi = \frac{\epsilon_0 q^2}{2(4\pi\epsilon_0)^2} \int \frac{1}{r^2} \underbrace{\sin\theta \, d\theta \, d\phi}_{4\pi} \, dr$$

$$= \frac{\epsilon_0 q^2 4\pi}{2(4\pi\epsilon_0)^2} \int_0^\infty \frac{1}{r^2} \, dr \rightarrow \infty$$

∞ SELF ENERGY OF ELECTRON PUZZLED MANY
GREAT 19TH CENTURY PHYSICISTS

↳ RENORMALIZATION IN QCD

FOR US → CHARGES GIVEN TO US

↳ WE JUST MOVE THEM AROUND

↳ ∞ SELF ENERGY IRRELEVANT

ENERGY IS STORED IN FIELD

$$\left(\begin{array}{c} \text{ENERGY PER} \\ \text{UNIT VOLUME} \end{array} \right) = \frac{\epsilon_0}{2} E^2$$

SUPERPOSITION & ENERGY

ELECTROSTATIC ENERGY QUADRATIC IN \vec{E}

→ DOES NOT OBEY SUPERPOSITION

ENERGY SYSTEM $\neq \sum$ OF PARTS

$$(E_1 + E_2)^2 = E_1^2 + E_2^2 + 2E_1E_2$$

$$W_{\text{TOT}} = \frac{\epsilon_0}{2} \int E^2 d\tau = \frac{\epsilon_0}{2} \int (\vec{E}_1 + \vec{E}_2)^2 d\tau$$

$$= \frac{\epsilon_0}{2} \int (E_1^2 + E_2^2 + 2\vec{E}_1 \cdot \vec{E}_2) d\tau$$

$$= W_1 + W_2 + \epsilon_0 \int \vec{E}_1 \cdot \vec{E}_2 d\tau$$

DOUBLE CHARGE → QUADRUPLE ENERGY