

14 December 2018

Thermal Physics

↳ What is it about?

- almost everything

the world around us

- room

- weather

- climate

- refrigerator

- furnace

- stars

- big bang.

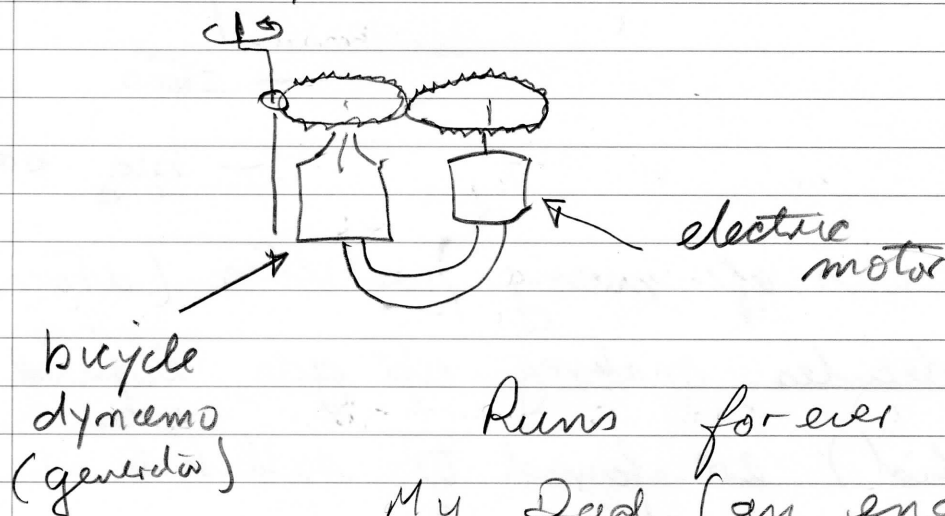
Systems of many particles (atoms, molecules, making up gas, liquid, solid) are found to exhibit "universal behaviour"

↳ you can predict a lot of the behaviour of system, without knowing the detailed internal structure.

Empirically Thermal Physics
originated in the 19th century
with the INDUSTRIAL REVOLUTION
PRACTICAL QUESTIONS, LIKE:

→ Can we build a perpetual
motion machine 1st kind

I thought you could when I was
a kid,



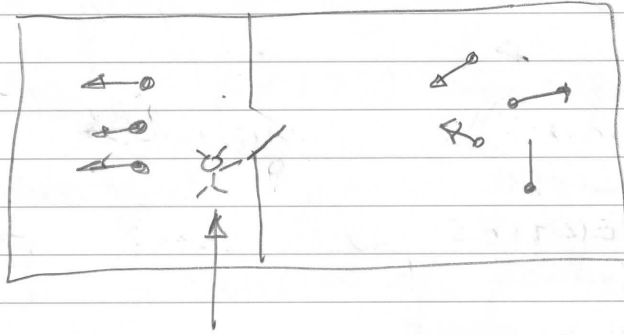
Runs forever
My Dad (an engineer)
said it would never
work

free energy

1st Law of thermodynamics
Conservation of Energy

→ Perpetual motion 2ND kind

useful energy from rectifying random motion of air molecules

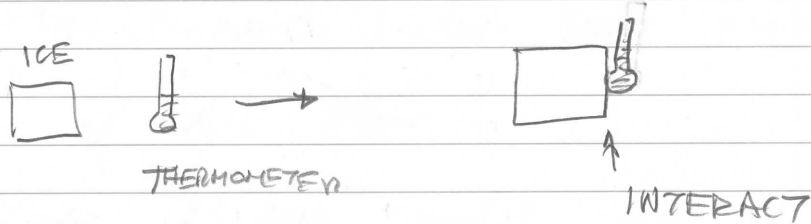


Maxwell demon

2ND Law of Thermodynamics

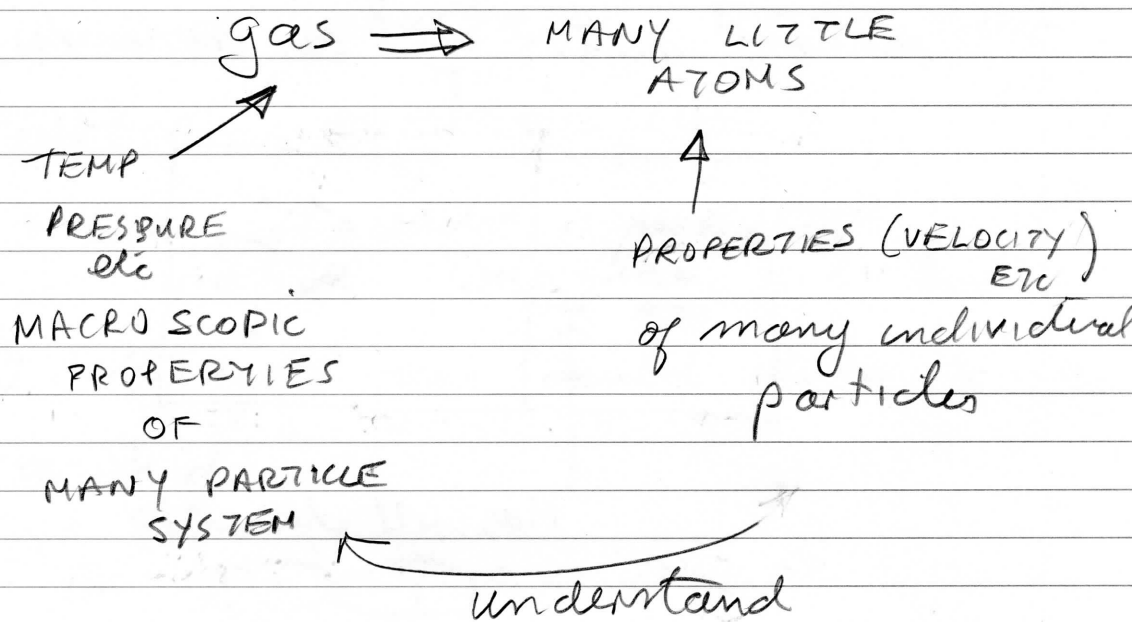
↳ Increase of Entropy

1) THERMODYNAMIC — study of general properties of systems of MANY PARTICLES which can INTERACT WITH EACH OTHER



2) STATISTICAL MECHANICS

Make a MICROSCOPIC model



Statistical Mechanics offers a way
to derive thermodynamics from
FIRST PRINCIPLES.

How many particles are 'MANY'?



$$m_{\text{H-atoms}} \approx 10^{-27} \text{ kg}$$

so one gram of hydrogen $\sim 10^{24}$ atoms

Avogadro — 6×10^{23} atoms per mole

$$N \sim 10^{24}$$

↳ number of particles in a system

How do you describe a system of $\sim 10^{24}$ particles?

Use classical mechanics?

$$m \ddot{\vec{r}}_i = \vec{F}(\vec{r}_1, \dots, \vec{r}_N, \dot{\vec{r}}_1, \dots, \dot{\vec{r}}_N)$$

$$i = 1 \dots N = 10^{24}$$

↳ 3×10^{24} equations.

Presently supercomputers can model motion of 10^{12} stars.

But even if we could do that computation \rightarrow it would be 6×10^{24} numbers for every moment in time

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How could you possibly make sense of that?

For a system of so many particles the individual behaviour of particles is of no real interest. Empirically you "feel" the "average action" of many particles in the air.

M MACROSCOPIC QUANTITIES

- pressure P average force / unit area
- N particles
- Volume V
- Temperature T

↳ quantity measured by a thermometer

As opposed to MICROSCOPIC QUANTITIES

$$\bar{r}_i(t) \quad \dot{\bar{r}}_i(t)$$

$$i = 1 \dots 10^{10^4}$$

6×10^2
position
and velocities

SUMMARY:

- Description of a system via MACROSCOPIC QUANTITIES is the subject of THERMODYNAMICS "COURSE GRAINED"
- Mechanics (classical or quantum) uses MICROSCOPIC QUANTITIES to describe a system — DETAILED, FINE GRAINED
- STATISTICAL MECHANICS, relates the two of them

Concrete example of a SYSTEM

↳ IDEAL GAS

will define what we mean by this

↳ obeys IDEAL GAS LAWS

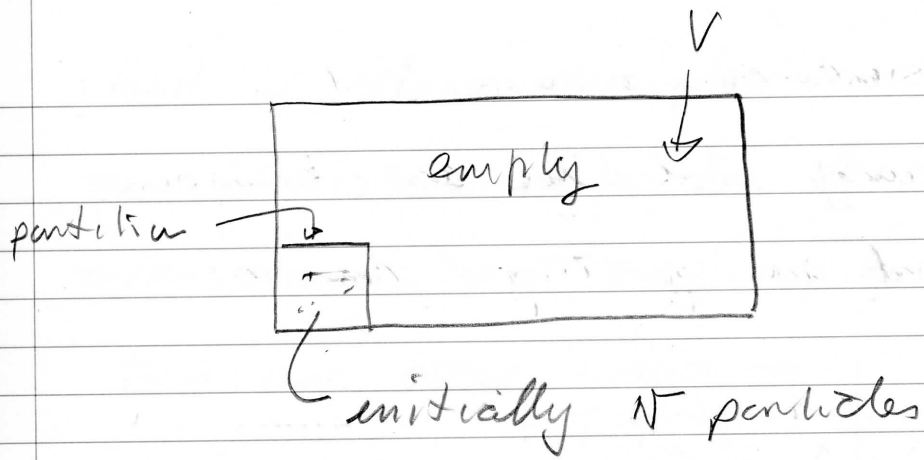
Important concept for any thermodynamic system is THERMODYNAMIC EQUILIBRIUM.

Experimentally a gas (N particles in volume V) is left on its own for a sufficiently long time it will

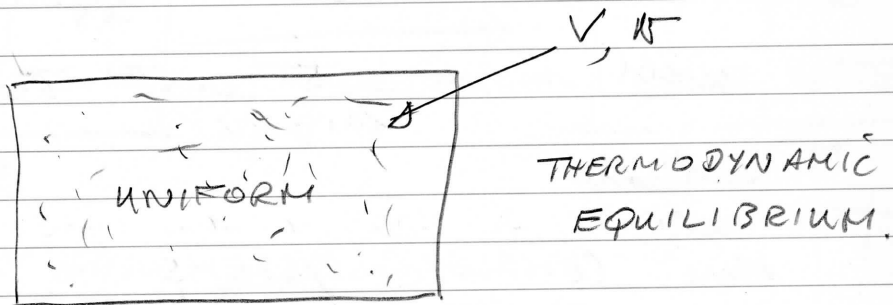
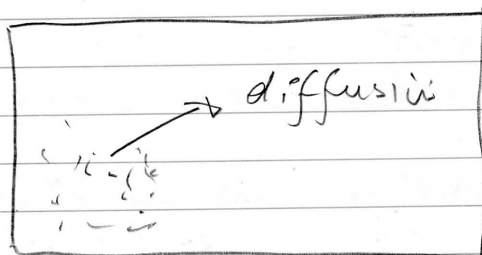
THERMODYNAMIC
EQUILIBRIUM.

- 1) be uniform, have same density, throughout the volume V
- 2) PRESSURE + TEMPERATURE will not change with time, and will be uniform throughout the volume
- 3) No microscopic fluxes (currents) winds

↳ individual particles will still be bouncing around but in no preferred direction



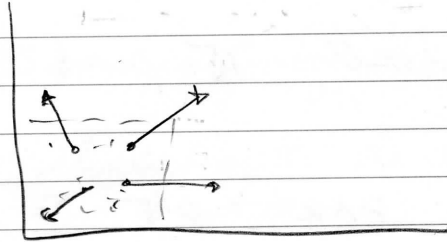
then remove PARTITION



- GAS uniformly fills box
- pressure and temperature and density are all UNIFORM IN SPACE
CONSTANT IN TIME

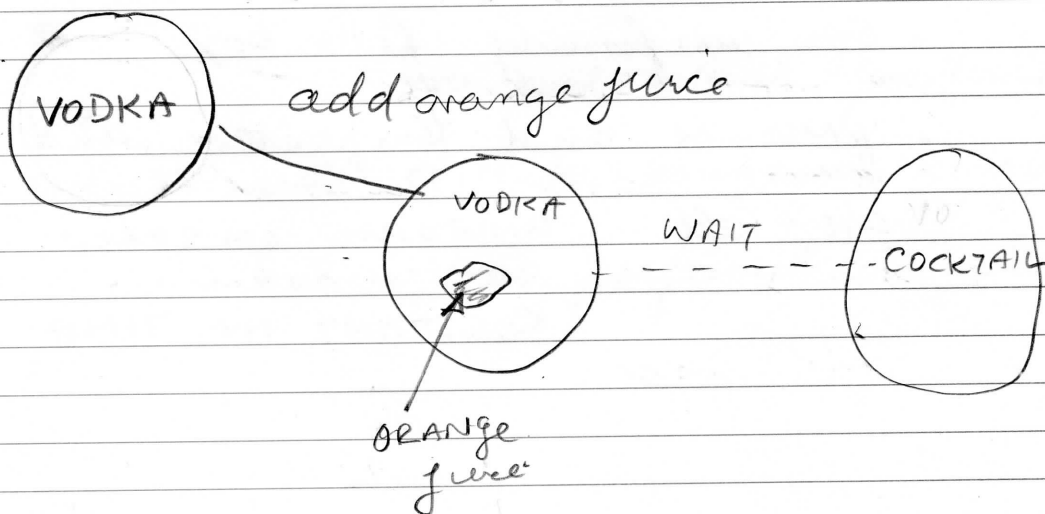
DIFFUSION

initially molecules are bouncing around in partitioned corner



Initially there were some molecules with velocities directed outwards when there is no partition, they will move out into the empty volume.

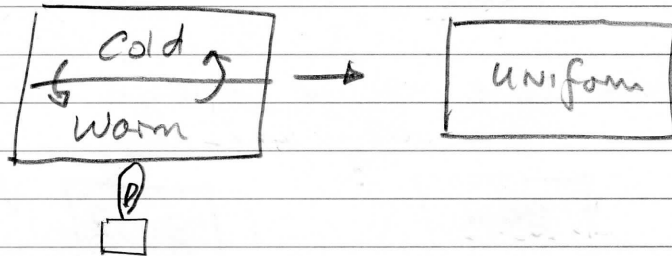
More COMPLEX SYSTEMS



other mechanisms leading to

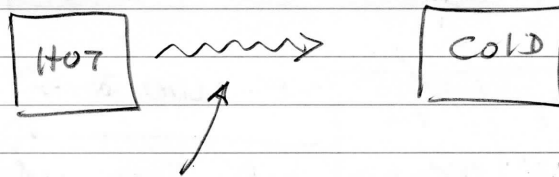
THERMODYNAMIC EQUILIBRIUM

CONVECTION

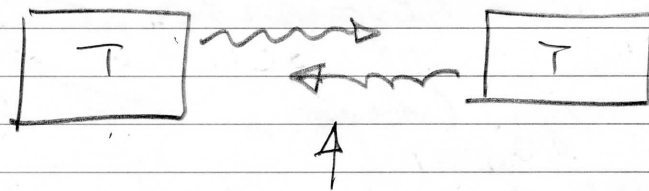


Heating water, warmer rises, colder falls

RADIATION

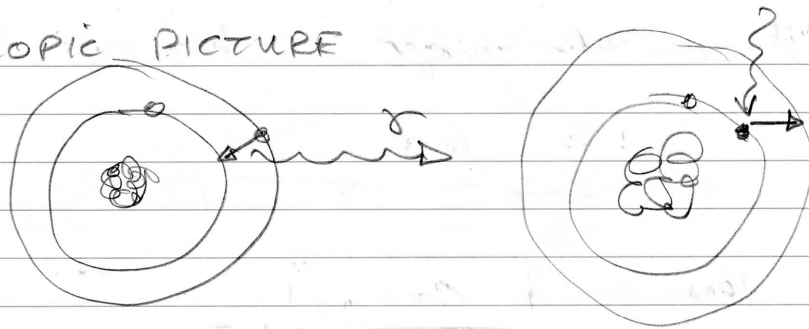


electromagnetic radiation carries
carries energy from HOT to COLD
eventually temperatures equalize



each is producing some
energy from radiation

MICROSCOPIC PICTURE



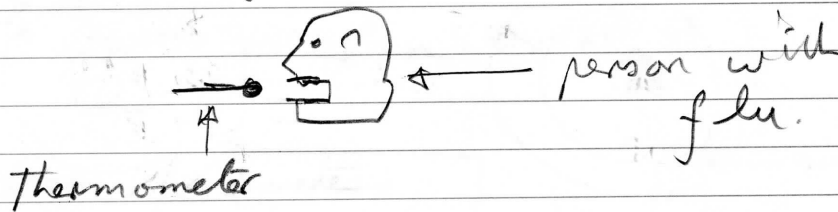
In thermodynamics and statistical mechanics we are usually not concerned with the detailed microscopic mechanism leading to the establishment of thermodynamic Equilibrium.

↳ should be aware of the general nature of the processes which lead to equilibrium

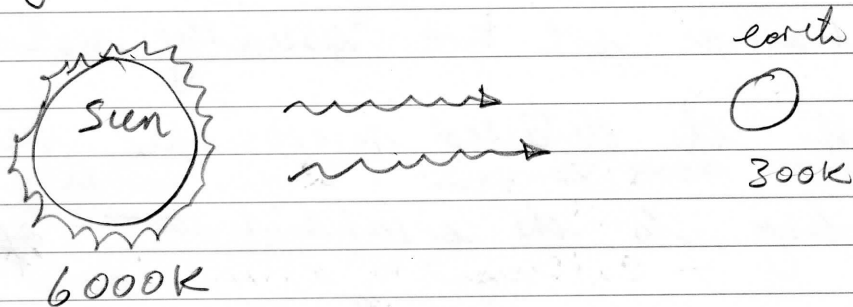
↳ heat exchange, diffusion, convection, radiation - - -

Also question of how long it takes to equalize pressure, density, temperature is not trivial — depends on mechanisms involved

It can be fast



very slow



not in equilibrium, luckily

PHYSICAL KINETICS studies rates
of diffusion, heat exchange, radiation etc

→ section 1-7 Schröder.

Concrete System - IDEAL CLASSICAL GAS

obeys ideal gas laws

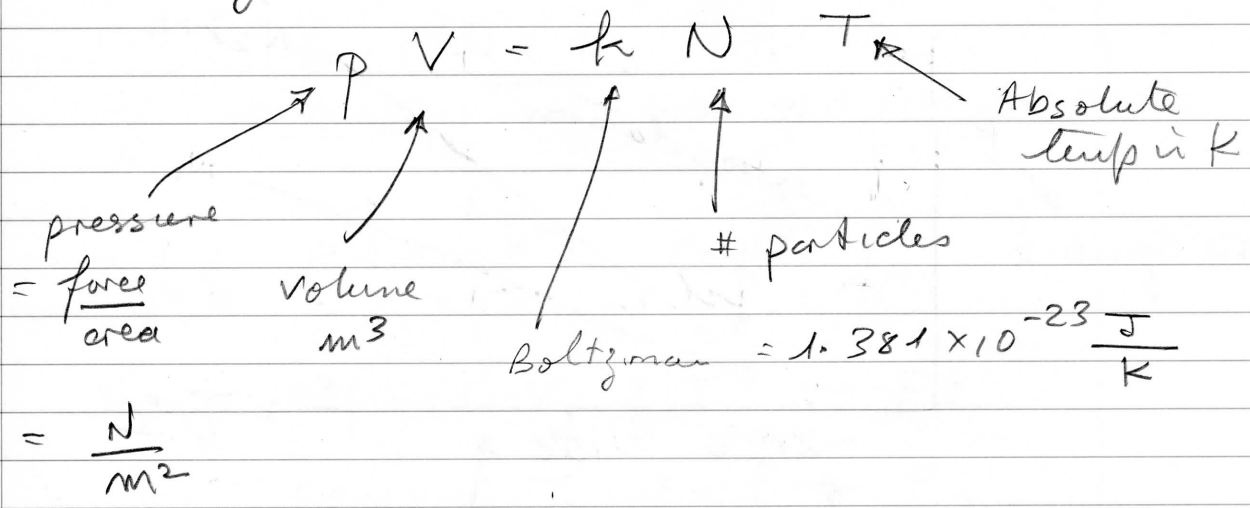
point like particles

$N \sim 10^{24}$

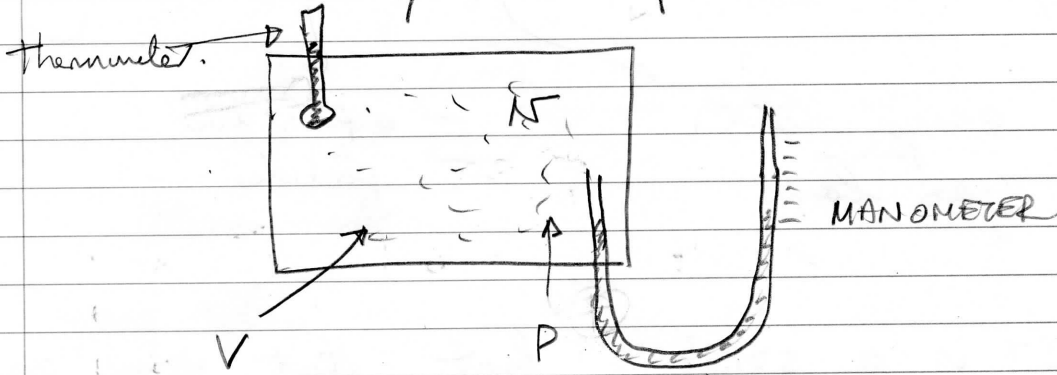
IT is a striking fact that any gas which is

- sufficiently hot
 - sufficiently dilute
- } obeys ideal gas law.

to great accuracy.

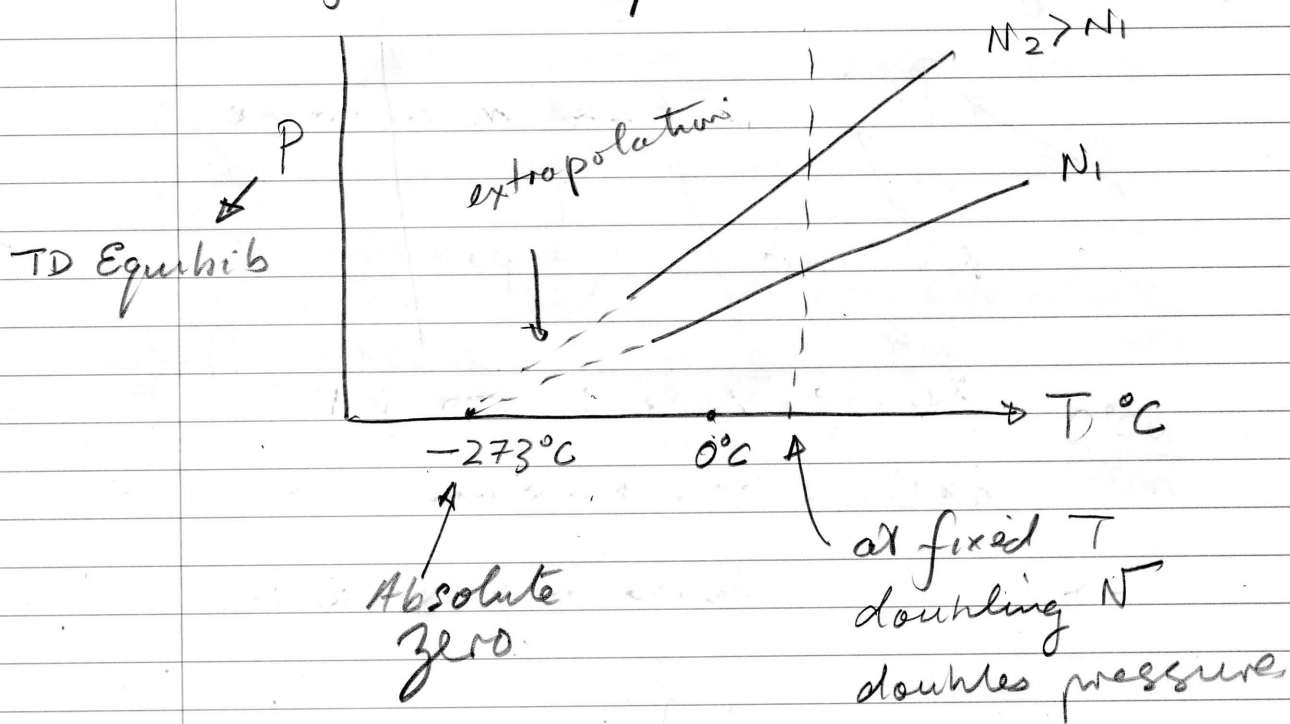


This is an experimental fact (18th → 19th Century?)



$$\text{pressure} \sim 1 \text{ atm} \sim 10^5 \frac{\text{N}}{\text{m}^2}$$

put, at fixed V , different amounts of molecules and measure pressure at different temperatures



$$pV = nRT$$

of moles of gas

$$1 \text{ Mole} = N_A = 6.02 \times 10^{23} \text{ molecules or atoms}$$

eg $N_A = \# \text{ of atoms in } 12\text{g of } ^{12}\text{C carbon}$

Comparing with $pV = nRT$

$$R = N_A \cdot k = 8.31 \frac{\text{J}}{\text{mol K}}$$

$$nR = N_A k$$

$$R = \frac{N_A \cdot k}{n} \rightarrow n=1 \rightarrow \text{per mole}$$

$$pV = nRT \text{ simple!}$$

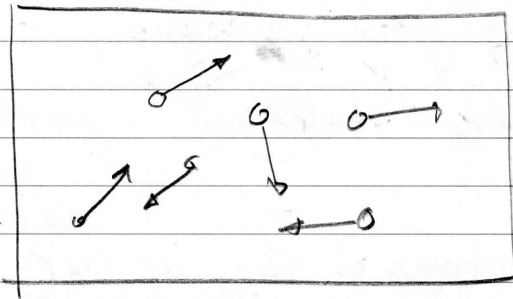
at fixed V ; if $T \uparrow \rightarrow P \uparrow$

at fixed T , if $V \uparrow \rightarrow P \downarrow$

\rightarrow common sense \rightarrow or is it experience

\downarrow
see SP441

Can deduce from SIMPLE MOLECULAR
MODEL OF GAS



$\sim 10^{24}$ molecules randomly
bouncing around.

→ uniformly distributed

→ velocities isotropically distributed

in any given volume containing
a large number of molecules, the
number having a given direction
of \vec{v} is the same.

↳ if this were not true
would have a macroscopic
flow in direction \vec{v} → this
would not be thermodynamic
Equilibrium

Important assumption is that there is some
AVERAGE SPEED

Call average velocity \bar{v}_x , because \bar{v} s are
isotropically distributed.

average $\bar{v}_x = \bar{v}_y = \bar{v}_z = 0$

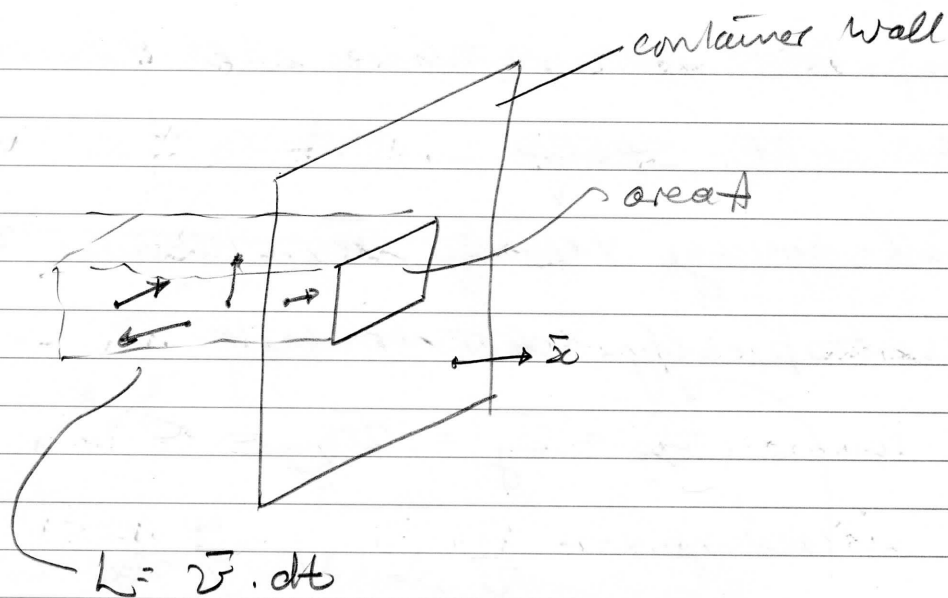
$$\bar{v}_x = \frac{1}{N} \sum_{i=1}^N v_{xi}$$

because as many move left as right
we can expect $\bar{v}_x = 0$, also \bar{v}_y, \bar{v}_z

\bar{v}_x^2 need not vanish, and we assume
that it does not.

$$\bar{v}_x^2 = \frac{1}{N} \sum_{i=1}^N v_{xi}^2$$

$$\sqrt{\bar{v}_x^2} = \bar{v} \quad \text{is our definition of } \bar{v}$$



In box LA volume there are

$$\frac{N}{V} \cdot \underbrace{L \cdot A}_{\text{volume}}$$

↑
density

expect that $\frac{1}{2}$ of molecules have velocity pointing in \bar{x} direction

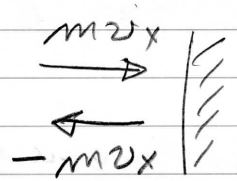
$L = \bar{v} dt$, in time dt all of the molecules inside the box will hit the wall

So, in time dt there will be

$$\frac{1}{2} \frac{N}{V} \cdot L \cdot A = \frac{1}{2} \frac{N}{V} v dt A \quad (1)$$

collisions with the wall

when a particle with x -momentum mv_x bounces off the wall, the change in its momentum is $-2mv_x$



$$|\Delta p| = |mv_x - (-mv_x)|$$

$$= 2mv_x$$

but it is in the $-ve$ x direction

$$\therefore \Delta p = -2mv_x \quad (3)$$

↳ this is how molecules exert pressure on the wall

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Ignore the fact that some molecules will have $v_y, v_z \neq 0$ and may not hit A

From ① and ③ we have that

the total change in momentum of
the wall of area A

= # collisions in
time dt \times change in momentum
of A for each
collision

$$\left(\frac{1}{2} \frac{N}{V} \bar{v} dt A \right) \times (2m\bar{v}) \quad \text{③'}$$

Change in momentum of A is, by
Newton's Second Law = force on
the wall

$$\frac{dp_x}{dt} = F_x \quad \text{④}$$

$$\text{③' } \rightarrow \text{force} \cdot dt = dp_x = \frac{1}{2} \frac{N}{V} 2m\bar{v}^2 A \cdot dt$$

$$\text{④ } \text{force} = \frac{dp_x}{dt} = \frac{N}{V} m\bar{v}^2 A$$

\hookrightarrow force on A , perpendicular
to it.

Pressure is "force per unit area"

$$\frac{\text{force}}{\text{Area}} = \frac{N}{V} m \bar{v}^2 A / A$$

$$\text{pressure } P = \frac{N}{V} m \bar{v}^2 \quad (5)$$

Remember ideal gas law.

$$pV = N kT$$

$$pV = N m \bar{v}_x^2 \quad \rightarrow \text{actually } v_x^2$$

So, we conclude that

$$kT = m \bar{v}_x^2$$

But we chose the x-direction arbitrarily so by isotropy $\bar{v}_x^2 = \bar{v}_y^2 = \bar{v}_z^2$

$$\text{So } \bar{v}_x^2 = \frac{1}{3} (\bar{v}^2) = \frac{1}{3} (\bar{v}_x^2 + \bar{v}_y^2 + \bar{v}_z^2)$$

More generally.
$$\bar{v}^2 = \frac{1}{N} \sum_{i=1}^N \bar{v}_i^2$$

so we have that

$$kT = m \bar{v}_x^2 = \frac{1}{3} m \bar{v}^2$$

$$kT = \frac{1}{3} m \overline{v^2}$$

$$\text{kinetic energy} = \frac{1}{2} m v^2$$

$$kT = \frac{2}{3} \frac{m \overline{v^2}}{2}$$

$$\frac{m \overline{v^2}}{2} = \frac{3}{2} kT$$

In thermodynamic equilibrium
the average kinetic energy per
particle is $\frac{3}{2} kT$ in an IDEAL
GAS

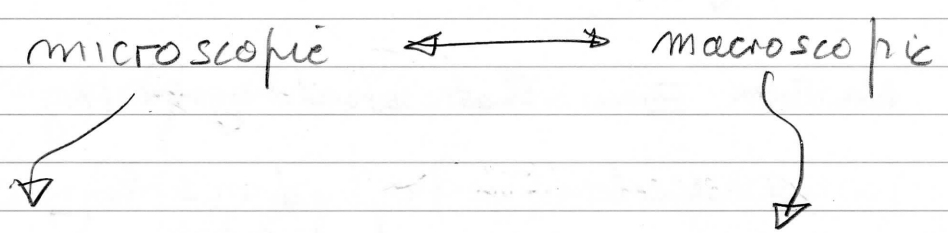
$$T \propto \frac{1}{2} m \overline{v^2}$$

So now we have a more

"scientific" definition of TEMPERATURE

TEMP \rightarrow is a measure of the
average kinetic energy
of the particles in
the gas

this is a nice example of:



Kinetic energy per particle ↔ Temperature of gas

↳ under the assumption of a uniform, random in direction Velocity distribution

Numerical stuff is always important and interesting

what is $\sqrt{\langle v^2 \rangle}$ Root Mean Square

$$v_{RMS} = \sqrt{\frac{3kT}{m}}$$

\swarrow k \searrow T
 \downarrow \downarrow

$$= \left(\frac{3 \times 1.381 \times 10^{-23} \frac{\text{joules}}{\text{Kelvin}} \times 300 \text{ Kelvin}}{2 \times 14 \times 1.6605 \times 10^{-27} \text{ atomic mass} \rightarrow \text{kilo}} \right)^{1/2}$$

atomic mass of N_2

$$= 517 \text{ m/s}$$

Let's just check dimensions

Newton = Force that would give 1kg an
acceleration of 1 m s^{-2}

$$F = ma$$

$$[\text{newton}] = [\text{kg}] \frac{[\text{m}]}{[\text{s}^2]}$$

Joule = work done by one Newton
over 1m

$$[\text{Joule}] = [\text{Newton}] \times [\text{m}]$$

$$\therefore [\text{Joule}] = \frac{[\text{kg}] \times [\text{m}] \times [\text{m}]}{[\text{s}^2]}$$

$$\text{So } \frac{[\text{Joule}]}{[\text{kg}]} = \frac{[\text{m}^2]}{[\text{s}^2]}$$

$$\left(\frac{[\text{Joule}]}{[\text{kg}]} \right)^{\frac{1}{2}} = \frac{\text{m}}{\text{s}} \quad \checkmark$$

SUMMARY

- Thermodynamic Equilibrium means uniform, isotropic, no MACROSCOPIC flows

- All dilute and hot gases obey

$$PV = NkT$$

→ This is a universal law, it does not depend on the detailed composition of the gas [as long as it is IDEAL]

- $\frac{m \langle \vec{v}^2 \rangle}{2} =$ average translational kinetic energy

$$= \frac{3}{2} kT$$

so $kT =$ energy

UNITS & DIMENSIONS

$$PV = nRT - k$$

$\frac{N}{m^2}$ (pointing to P)
 m^3 (pointing to V)
 # of moles (pointing to n)
 Joules (pointing to R)
 gas constant (pointing to k)

$$R = \frac{N \cdot m}{m \cdot T - k}$$

$$PV = [N][m] = [\text{mole}][k][R]$$

$$\text{so } [R] = \frac{[N \cdot m] \leftarrow J}{[\text{mole}][k]}$$

Above $[N]$ is Newtons

can use N for number of particles

$$n[\text{moles}] = \frac{N}{N_A} \quad \text{AVOGADRO}$$

So can write

$$PV = \frac{N}{N_A} \cdot R \cdot T$$

$$PV = \frac{N}{N_A} \cdot R \cdot T$$

Boltzmann

Also have $PV = N k T$

$$k = \frac{R}{N_A} = \frac{[\text{J}]}{[\text{mole}] \cdot [\text{K}]} \cdot [\text{mole}]$$

$$k = \frac{[\text{J}]}{[\text{K}]} \cdot \frac{[\text{Joules}]}{[\text{Kelvin}]}$$

$$R = 8.31 \frac{\text{J}}{\text{mole} \cdot \text{K}}$$

$$k = \frac{8.31}{N_A} \frac{\text{J}}{\text{mole} \cdot \text{K}} = 1.38 \times 10^{-23}$$

Extra Points from Schröder

average translational kinetic energy

$$= \bar{K}_{\text{trans}} = \frac{3}{2} kT$$

Nice interpretation of k . It has just the right units to convert between temperature and energy

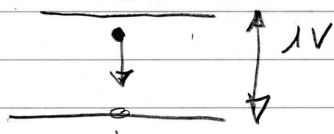
$$[k] = \left[\frac{\text{J}}{\text{K}} \right]$$

at Room temperature of 300 K

$$kT = 1.38 \times 10^{-23} \frac{\text{J}}{\text{K}} \cdot 300 \text{ K}$$
$$= 4.14 \times 10^{-21} \text{ J}$$

this is TINY, but nice molecules are tiny - that's just what we expect.

As a particle physicist, I like to use ELECTRON VOLTS



electron energies with kinetic energy of 1 volt.

$$1 \text{ eV} = 1.6 \times 10^{-19} \text{ J}$$

in these units $k = 8.62 \times 10^{-5} \text{ eV/K}$

so, at room temperature

$$kT = \frac{8.62 \times 10^{-5} \text{ eV}}{\text{K}} \cdot 300 \text{ K} = 0.026 \text{ eV}$$

Even in electron volts molecular energies are pretty small numbers

$$v_{\text{rms}} \sim \bar{v}$$

Derivation breaks down if molecules exert forces on each other, or if collisions with walls are inelastic

↳ gas becomes so dense that space occupied by molecules is a substantial fraction of the volume of the container

↳ then approximation of molecules flying in straight lines breaks down

$K_{\text{trans}} = \frac{3}{2} kT$ remains true even for liquids & solids