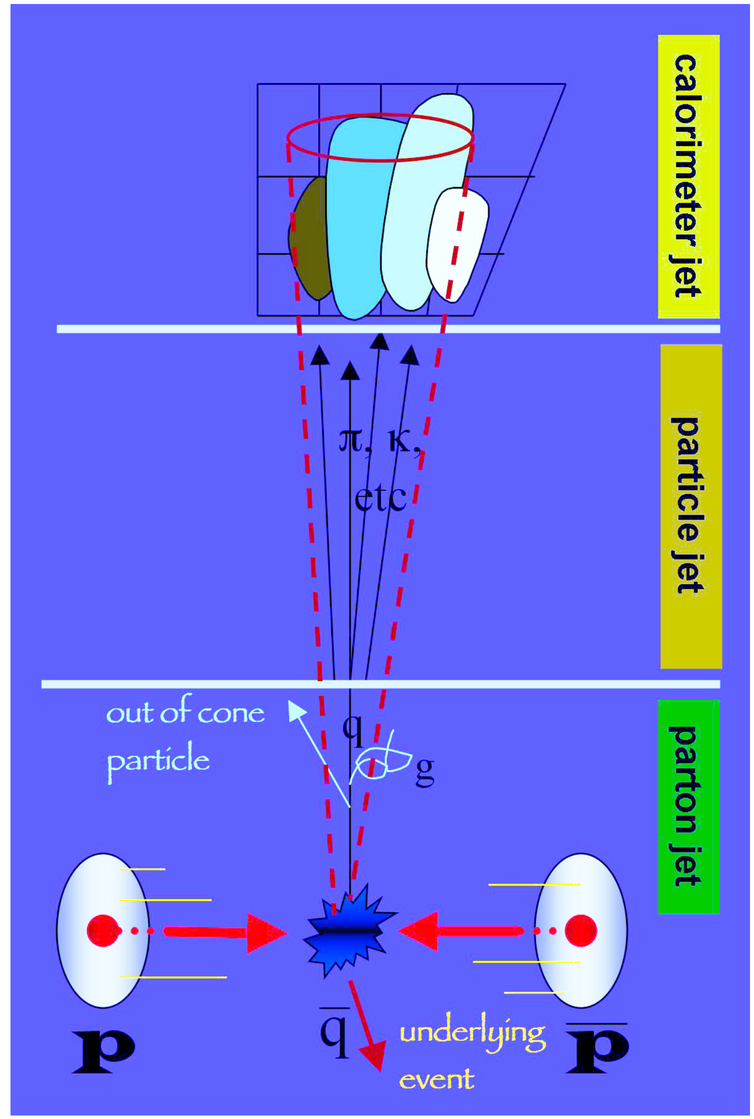
# Hard Scattering in Hadron-Hadron Collisions: Physics and Anatomy

## Section 4: Production & Identification of Jets

1. Definitions of Basic Physics Processes
2. Anatomy of a Jet
3. Jet-Finding Algorithms
4. Resolutions and Efficiencies
5. Heavy Quark Tagging
6. Example: Quark Substructure

# Definitional Issues

- **Confinement in QCD ensures that high $P_T$ quarks & gluons undergo**
  - **Fragmentation -- ie, dissociation into a "jet" of coloured partons**
  - **Hadronization -- ie, the partons form colourless, observable hadrons**
- **Study of jets motivated by**
  - **Understanding QCD**
  - **Studying of heavy quarks**
    - b/c quarks that fragment & hadronize before decay
    - Top quarks that decay before fragmentation/hadronication
  - **Searching for new interactions that couple to quarks/gluons**
  - **Jets as a background source to e, μ, τ & ν**

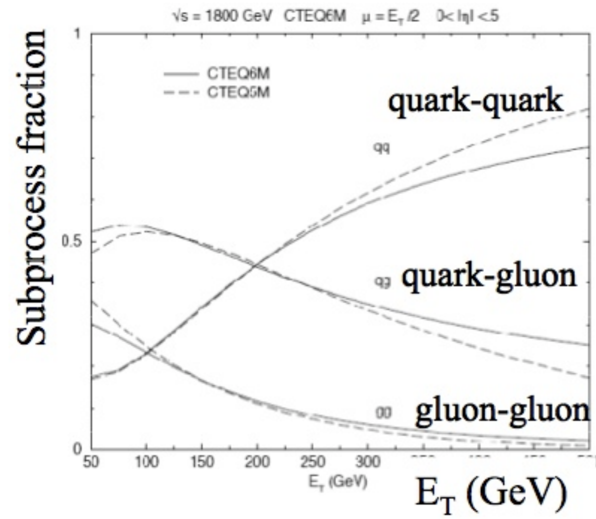


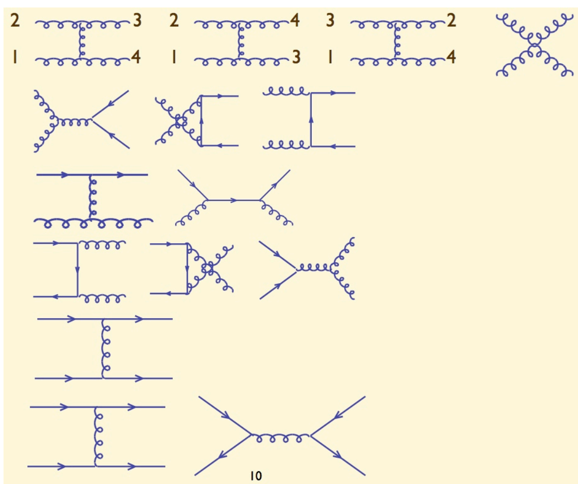
# Fundamentals of Jet Physics

- **Basic production mechanism in pQCD starts with**

$$\sigma = \sum_{\substack{partons\ i \\ colour\ j}} C_{ij} \int_0^1 d\tau \int_\tau^1 \frac{dx_1}{\tau} \left[ f_1(x_1) f_2(\tau / x_1) \right] \hat{\sigma}(\tau s)$$

  - **Leading-order (LO) diagrams already complex**



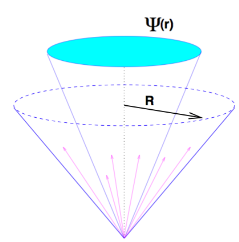



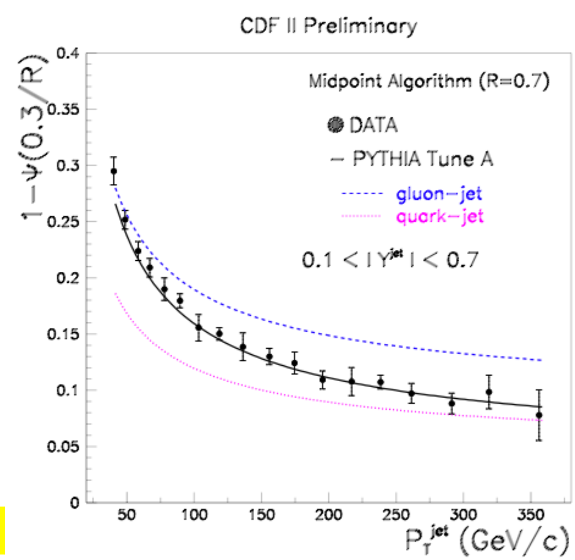
| Process | $\bar{\Sigma}\lvert\mathcal{M}\rvert^2/g^4$ | Numerical value for 90° |
|---|---|---|
| $qq' \to qq'$ | $\frac{4}{9}\frac{\hat{s}^2+\hat{u}^2}{\hat{t}^2}$ | 2.22 |
| $q\bar{q}' \to q\bar{q}'$ | $\frac{4}{9}\frac{\hat{s}^2+\hat{u}^2}{\hat{t}^2}$ | 2.22 |
| $qq \to qq$ | $\frac{4}{9}\left(\frac{\hat{s}^2+\hat{u}^2}{\hat{t}^2}+\frac{\hat{s}^2+\hat{t}^2}{\hat{u}^2}\right)-\frac{8}{27}\frac{\hat{s}^2}{\hat{u}\hat{t}}$ | 3.26 |
| $q\bar{q} \to q'\bar{q}'$ | $\frac{4}{9}\frac{\hat{t}^2+\hat{u}^2}{\hat{s}^2}$ | 2.22 |
| $q\bar{q} \to q\bar{q}$ | $\frac{4}{9}\left(\frac{\hat{s}^2+\hat{u}^2}{\hat{t}^2}+\frac{\hat{t}^2+\hat{u}^2}{\hat{s}^2}\right)-\frac{8}{27}\frac{\hat{u}^2}{\hat{s}\hat{t}}$ | 2.59 |
| $q\bar{q} \to gg$ | $\frac{32}{27}\frac{\hat{t}^2+\hat{u}^2}{\hat{t}\hat{u}}-\frac{8}{3}\frac{\hat{t}^2+\hat{u}^2}{\hat{s}^2}$ | 1.04 |
| $gg \to q\bar{q}$ | $\frac{1}{6}\frac{\hat{t}^2+\hat{u}^2}{\hat{t}\hat{u}}-\frac{3}{8}\frac{\hat{t}^2+\hat{u}^2}{\hat{s}^2}$ | 0.15 |
| $gq \to gq$ | $\frac{4}{9}\frac{\hat{s}^2+\hat{u}^2}{\hat{s}\hat{u}}+\frac{\hat{u}^2+\hat{s}^2}{\hat{t}^2}$ | 6.11 |
| $gg \to gg$ | $\frac{9}{2}\left(3-\frac{\hat{t}\hat{u}}{\hat{s}^2}-\frac{\hat{s}\hat{u}}{\hat{t}^2}-\frac{\hat{s}\hat{t}}{\hat{u}^2}\right)$ | 30.4 |

# What Have We Learned?

- **Definition of jets critical**
  - **Much evolution in algorithms**
  - **Driven in large measure by theoretical considerations**
- **Calibration of jets requires data-driven techniques**
  - **Developed several techniques to calibrate *in situ***
  - **Still "work in progress"**
- **Approach to jet-finding and calibration driven by physics**
  - **Best example is comparison between**
    - QCD tests
    - Reconstruction of heavy objects (top and Higgs)
- **Need data to understand jets as backgrounds**
  - **Examples include**
    - Lepton ID
    - MET measurement
    - Heavy quark tagging
  - **Use to "calibrate" MC/simulation**
- **Bottom line: SM Picture of QCD works well**

$$\Psi(r) \equiv \frac{1}{N_{jet}} \sum_{jets} \frac{P_T(0,r)}{P_T(0,R)}$$

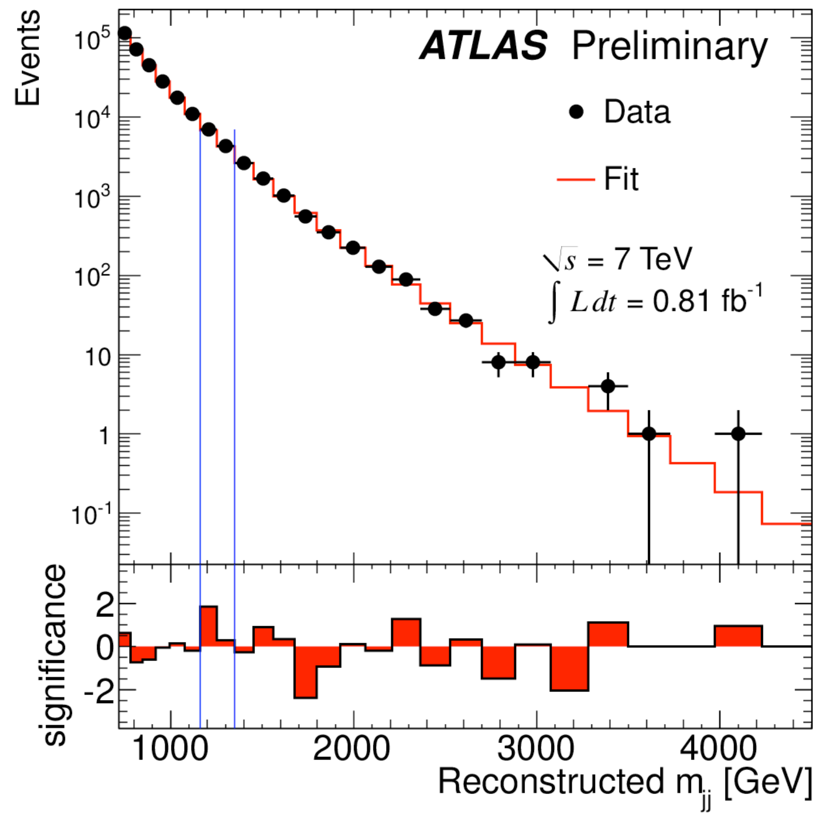


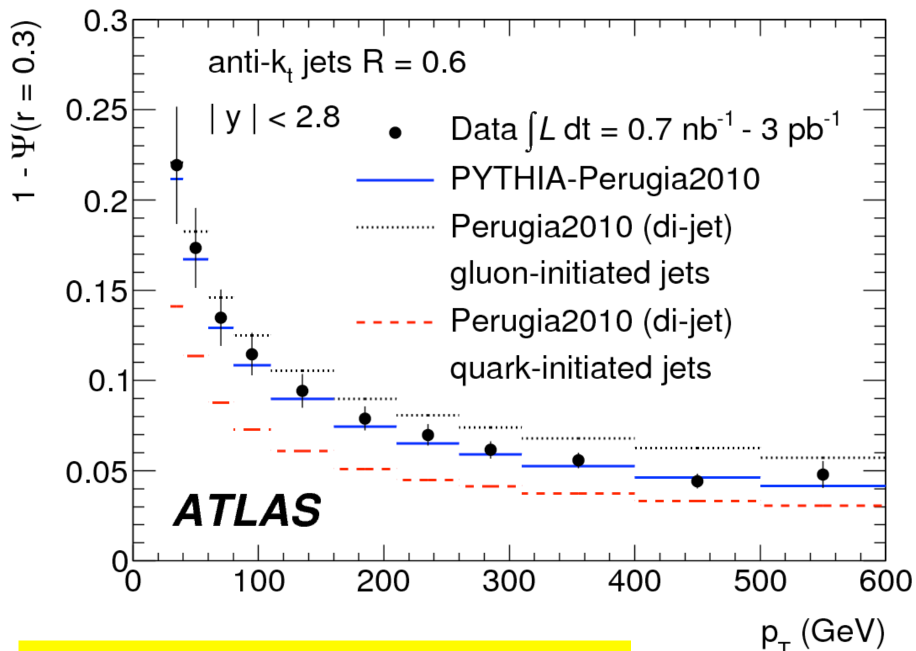

D. Acosta et al. (CDF), Phys. Rev. D 71, 112002 (2005)

# LHC Lessons

- **LHC studies have reproduced many of these effects**
  - **However, much higher jet momenta**
    - Jets with $p_T \sim 2$ TeV
  - **Focus has been on searches**
    - Looking for resonances in dijet mass
    - Sensitive to excited quarks




ATLAS-CONF-2011-096, 18 July 2011




ATLAS Collaboration, Phys. Rev. D 83, 052003 (2011)

# Jet Anatomy

- A jet arises from 2 different physical phenomena
  - Happen at different energy scales
    - Fragmentation of initial parton
      - QCD radiation of a coloured object
      - Creates a "cluster" of coloured partons
      - In principle, not independent of rest of event
      - Energy scale >> 1 GeV
    - Hadronization of "cluster"
      - Formation of colourless objects -- mesons & baryons
      - Responsible for the real observables
      - Energy scale ~ 1 GeV
- Have to worry about
  - What defines a jet (algorithm)?
  - What its properties are (recombination scheme)?

G. Blazey et al., FERMILAB-CONF-00-092-E and hep-ex/0005012, May 2000.

- First, tackle easiest part: What is a jet's observable properties?
  - Assume you have a collection of final state mass-less "particles" detected in calorimeter towers *i*

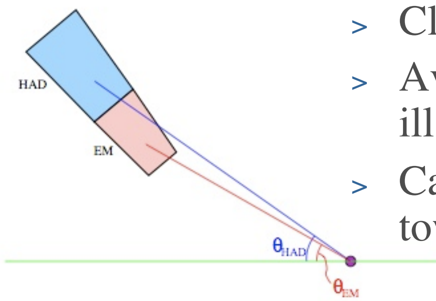
$$\vec{p}_J \equiv \left(E^J, p_x^J, p_y^J, p_z^J\right) \equiv \sum_i \left(E^i, p_x^i, p_y^i, p_z^i\right)$$

$$p_T^J \equiv \sqrt{\left(p_x^J\right)^2 + \left(p_y^J\right)^2}$$

$$M^J \equiv \sqrt{\left(E^J\right)^2 - \left(p^J\right)^2}$$

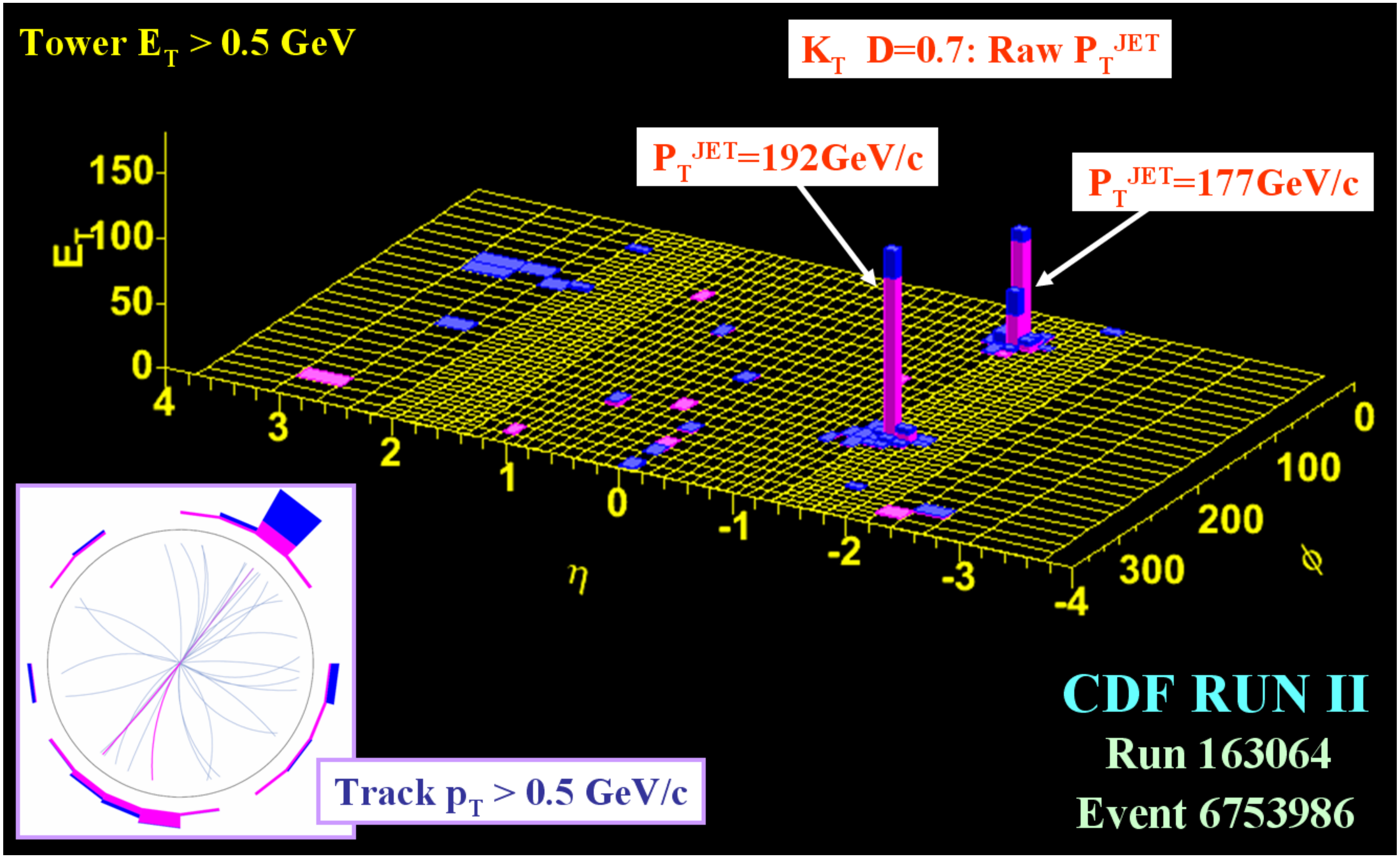
$$y^J \equiv \frac{1}{2}\ln\frac{E^J + p_z^J}{E^J - p_z^J}$$

$$\varphi^J \equiv \tan^{-1} p_y^J / p_x^J$$

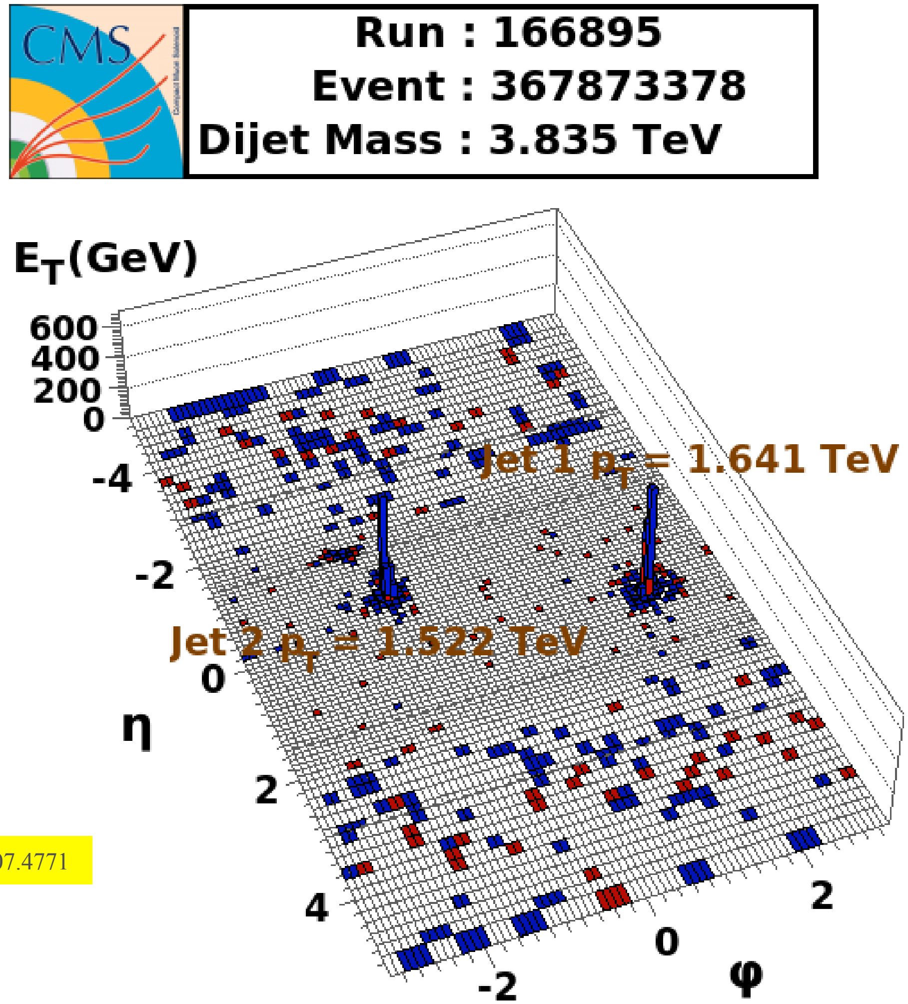



  - Advantages:
    - Clear Lorentz behaviour
    - Avoids use of $E_T$ which has ill-defined definition
    - Can generalize to "cells", towers, charged particles, etc.

# A Real Jet Event



# An LHC CMS Jet Event




Highest dijet Mass event Observed by CMS

CMS Collaboration, hep-ex/1107.4771

# Parton Shower Evolution

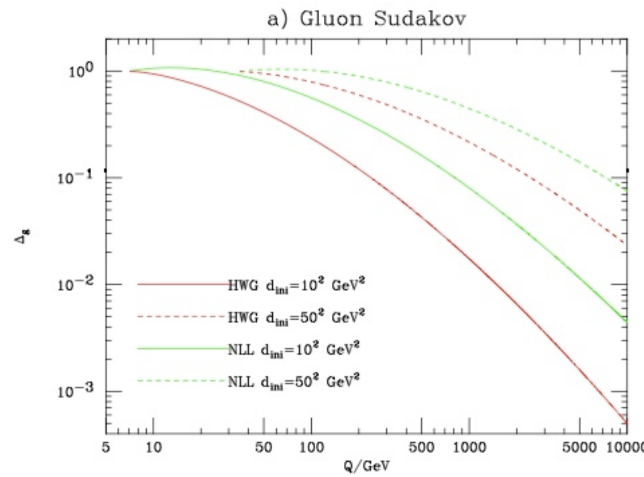
- **Start with a parton (q/g) with virtuality $\mu^2$**
  - **Probability of emission with daughter carrying $z$ fraction of parent momentum**

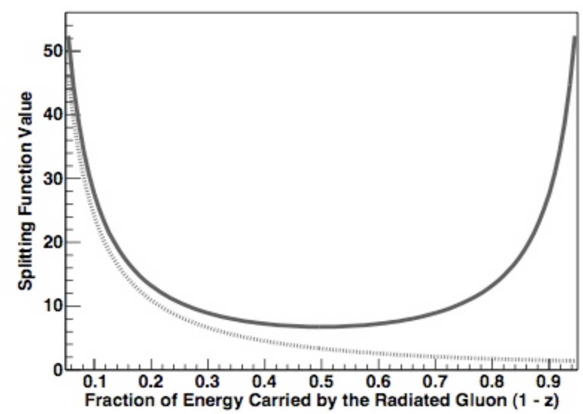
$$d^2P_a\left(z,\mu^2\right)=\frac{d\mu^2}{\mu^2}\frac{\alpha_s}{2\pi}P_{a\to bc}(z)dz$$

  - **Order these using Sudakov factor, relating $\mu^2 \sim Q^2$**

$$P_a^{no}\left(Q_{\max}^2,Q^2\right)=\exp\left(-\int_{Q^2}^{Q_{\max}^2}dQ'\int_{z_{\min}}^{z_{\max}}dz'\,P_a\left(z',Q'^2\right)\right)$$

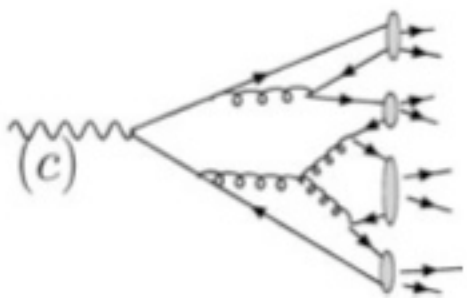
  - **Deal with infrared & collinear divergences**
    - Define minimum $\mu$ -- $\mu_0$
  - **Ensure colour coherence of multiple emissions**
    - Typically do this by angular ordering, selective vetoing, etc.
    - Must be respected when hadronization is performed

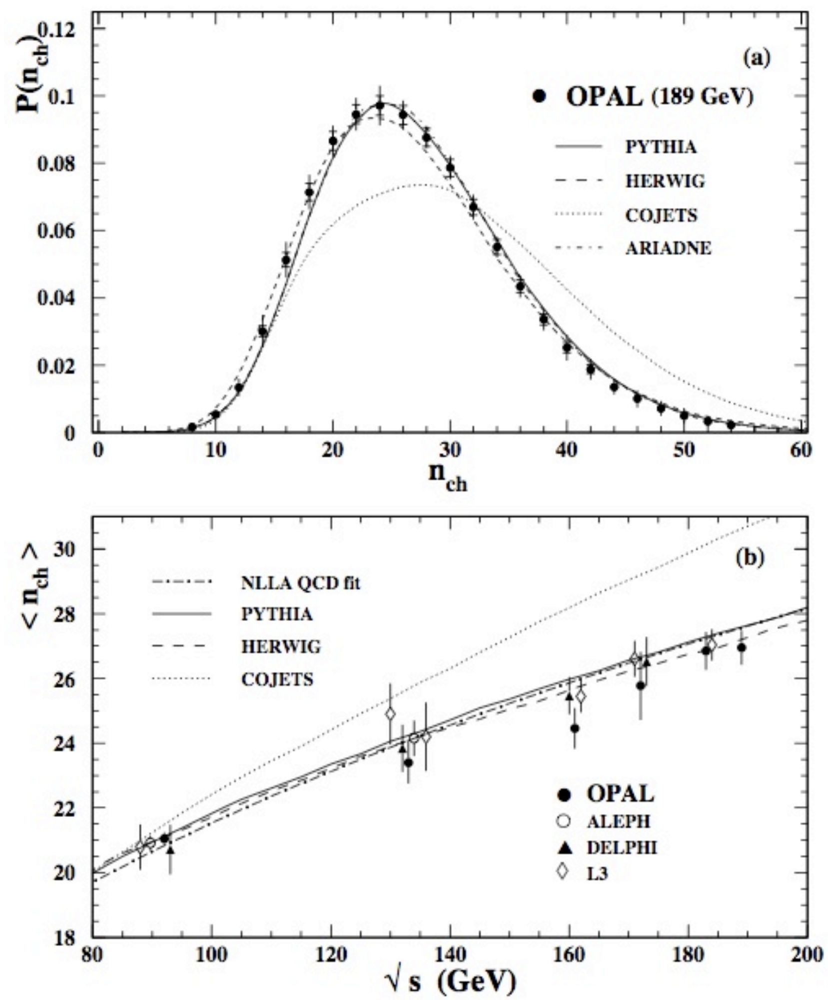
# Hadronization of Showers

- **Hadronization is then performed**
  - **Invoke "parton-hadron duality"**
  - **Several models**
    - String fragmentation (eg., PYTHIA)
    - Cluster fragmentation (eg. HERWIG)




  - **Have various parameters that need to be tuned to data**
    - Best constraints from LEP
      - **Tevatron results confirm these, but don't really add much power**
      - **Challenging to measure without significant systematics**
    - Remains a source of systematic uncertainty

OPAL, Eur. Phys. J C16, 185 (2000)



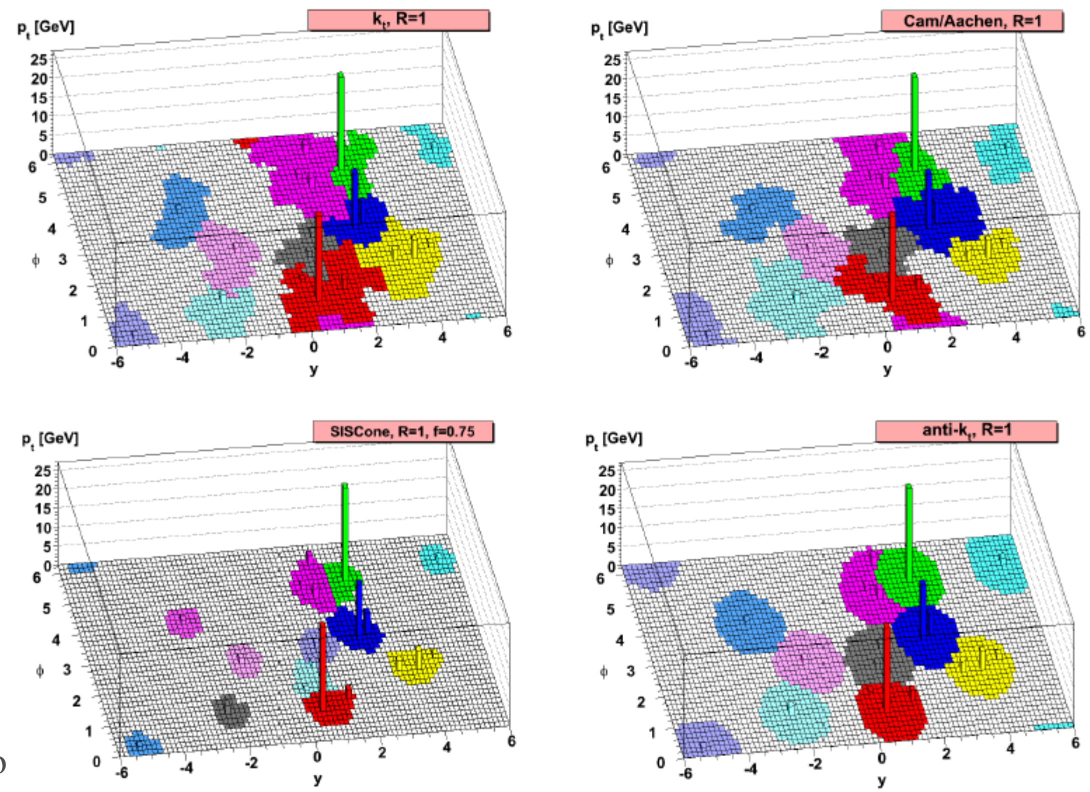
# Jet Algorithms

- Jet clustering algorithms have been focus of much effort
  - Goals of any algorithm can be divided into
    - Theoretically motivated:
      - Fully specified
      - Detector independent
      - Theoretically well-behaved
      - Order independent
    - Experimentally motivated:
      - Fully specified
      - Detector independent
      - Optimal resolution and efficiency
      - Ease of calibration
      - Computationally efficient
- Various efforts to develop consistent frameworks
  - Snowmass Accord (1990)
  - Les Houches Accord (1999)
- Raz Alon (see talk below) has done a nice job of summarizing current Jet Algorithm codes
  - Key observations:
    - In principle, prefer some algorithms over others
      - Seedless cone-based algorithms
      - $K_T$ algorithms
    - Computational efficiency is a concern in some cases
      - But largely an issue of optimization
    - Selection of "best" algorithm requires evaluation of ultimate systematic uncertainties
      - Need data, as certain choices will depend on performance of calorimeter
      - Example is noise and pileup
  - Good news is that we are not limited by lack of ideas

R. Alon, http://indico.cern.ch/conferenceDisplay.py?confId=52628

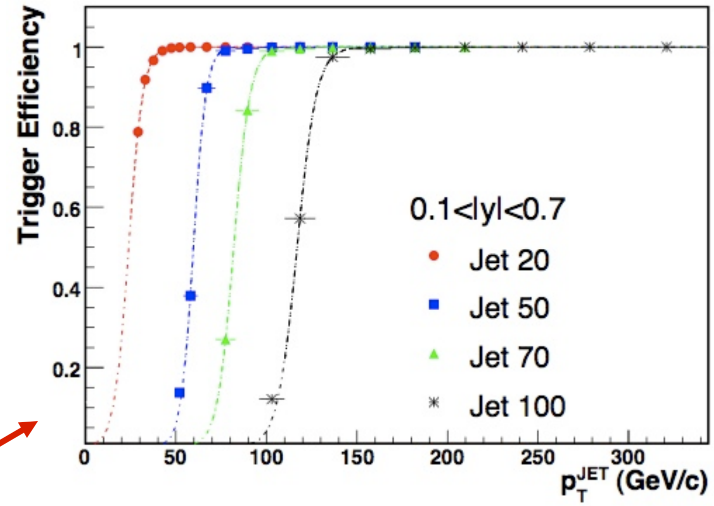
# Clustering Effects

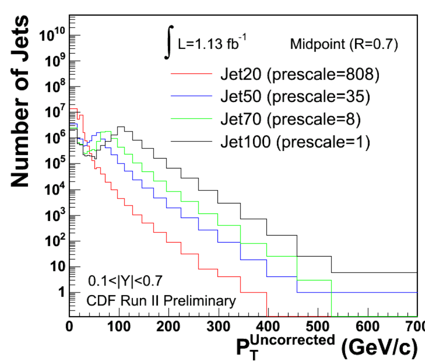
- **Illustrate by one example (from ATLAS studies)**
  - **Compare results of several different algorithms**
    - $K_T$ with R=1
    - Angular-ordering (Cam/Aachen)
    - SISCone
    - Anti-$K_T$
  - **Things to be concerned about**
    - Cluster sizes determined by data will present challenges to calibrate
    - Cluster merging/splitting will continue to be a challenge
    - Optimization of resolution/ systematic uncertainties will require effort
  - **Things not to worry about**
    - Angular resolution (though need to check for any biases)!

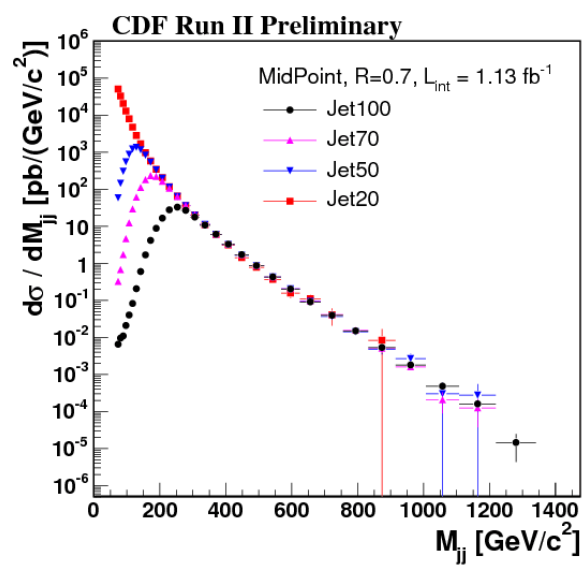


# Jet Finding Efficiencies

- Efficiency of finding jets limited primarily by two effects:
  - Detector energy response & resolution
  - Physical size of jets
    - For cone algorithms, these two compete with each other
- Further complicated by the fact that jets are produced with sharply falling spectrum
  - Means that efficiencies become an issue already at the trigger level
  - Manage these at Tevatron with variety of triggers
    - Prescale lower-energy jet triggers
    - Lower energy jets used primarily for
      - Background studies
      - Calibration



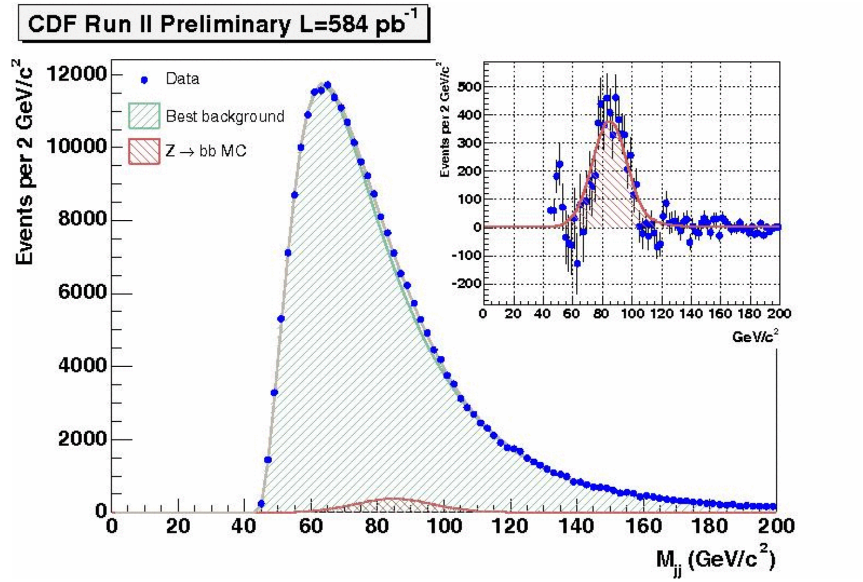
# Jet Energy Resolutions

- MC + simulation give estimates of energy resolution
  - Resolution is determined primarily by convolution of
    - Intrinsic calorimeter response
    - Jet fragmentation & hadronization effects
    - Jet algorithm + pileup + ….
  - In reality, need to measure the resolution in data
- Four *in situ* measurements of resolution developed at Tevatron
  - γ+jet balancing
  - W to qq in top quark decays
  - Dijet balancing (more of a constraint than anything else)
  - Z to bb decays
    - Require two jets, each with secondary vertex b-tag
      - Possible due to L2 vertex trigger
- Taking the FWHM ~ 25 GeV/c², obtain

$$\sigma_Z \sim 12\%\ M_Z$$

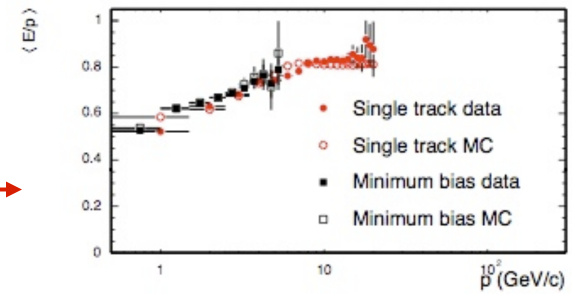
$$\Rightarrow \frac{\sigma}{P_T^J} \sim 17\%$$

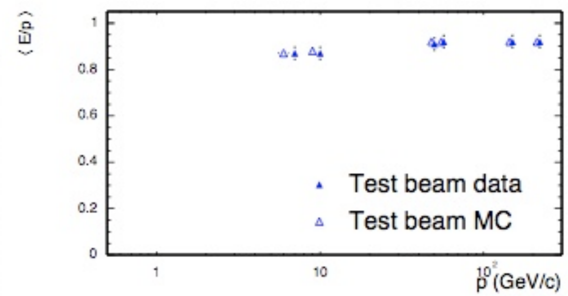
  - Or about 50% more than intrinsic energy resolution of calorimeter

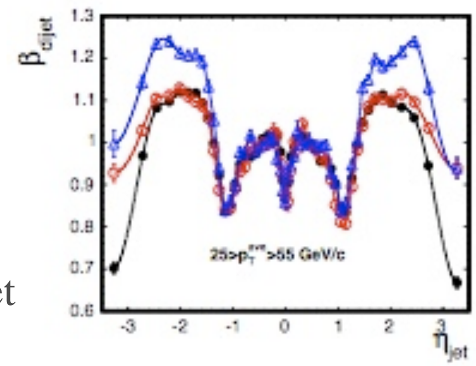


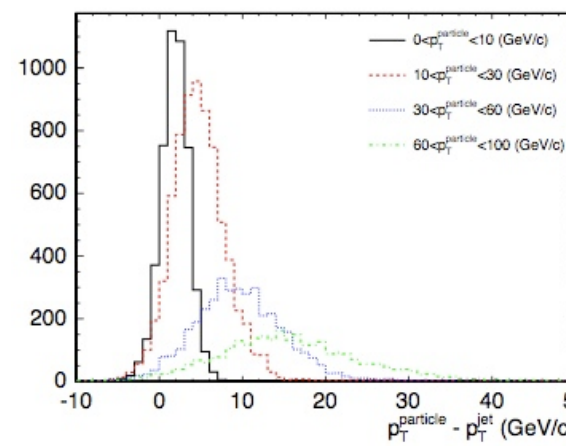
# Jet Energy Calibration

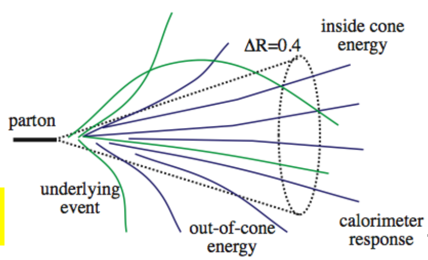
- **To calibrate jet energy scale:**
  - **1. Determine intrinsic response to particles**
    - Combination of in situ measurements & test beam data
  - **2. Dijet balancing to get uniform η response**
    - Primarily dijet data
    - "Tune" MC and simulation
  - **3. Determine absolute response to "particle jet"**
    - Define particle jet as all real particles in cone of jet
    - Account for calorimeter nonlinearity, cracks, etc.
  - **4. Take into account "out-of-cone" effects, multiple interactions**
    - Use combination of MC and data



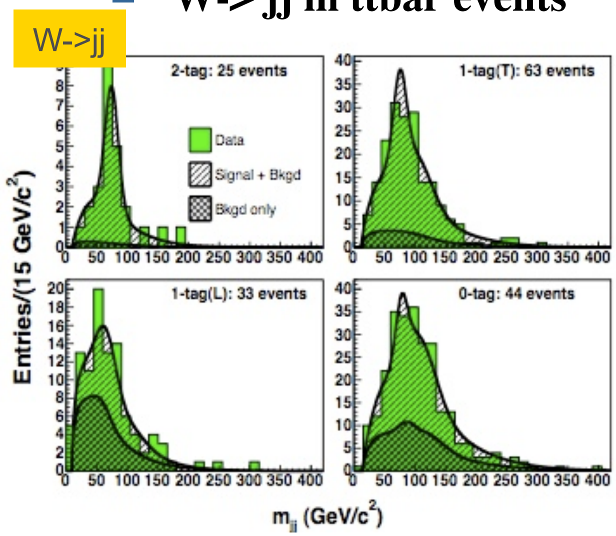



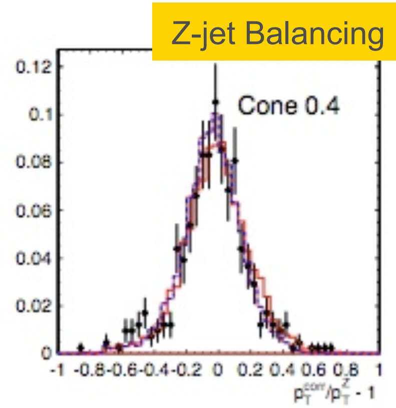



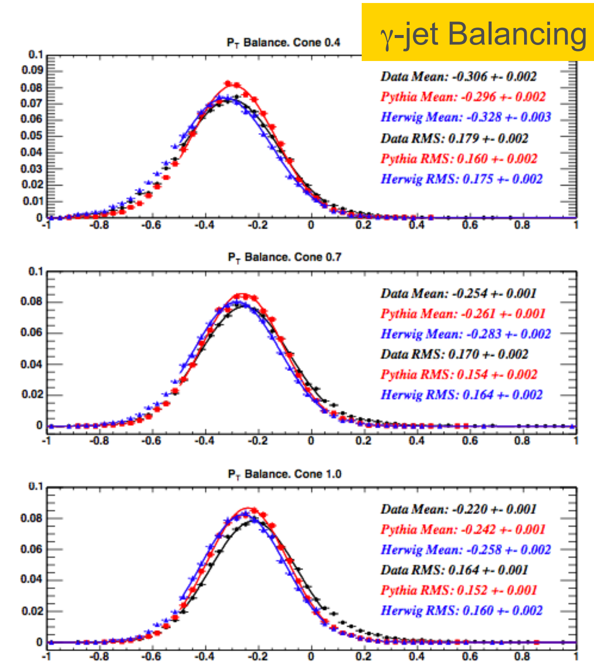

A. Bhatti et al., Nucl. Instrum. Meth. A566, 375 (2006)

# Final Steps in Energy Calibration

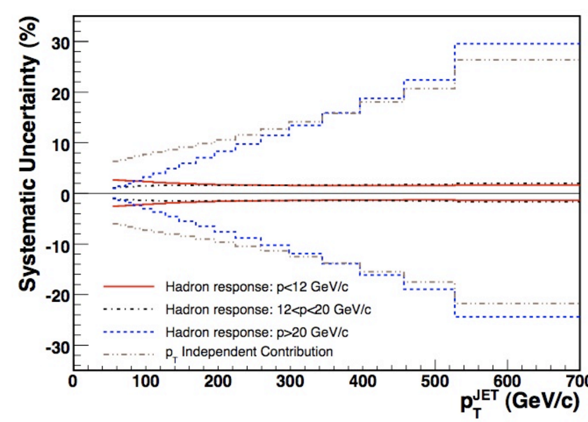
- Cross check using, for example,
  - Z+jet & γ+jet balancing
  - Dijet balancing
  - W-> jj in ttbar events



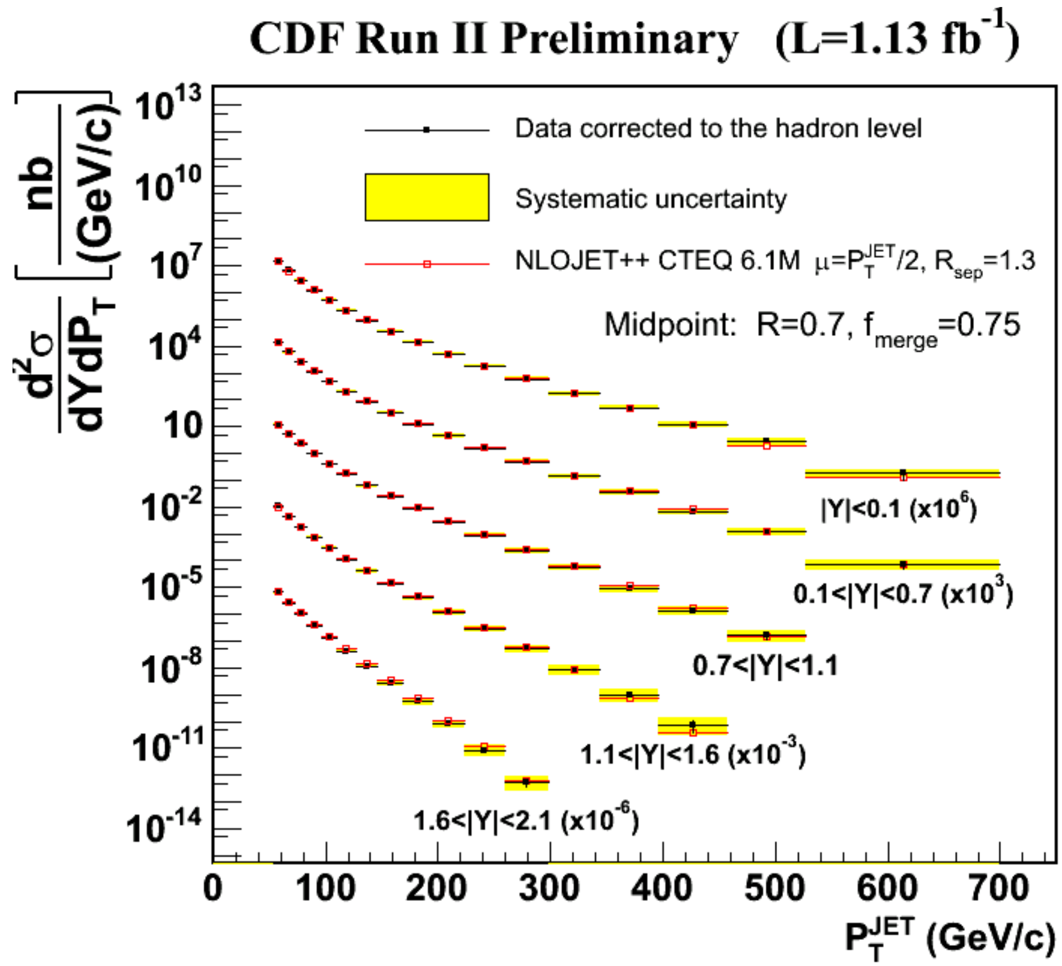




- Estimate systematic uncertainties
  - Estimate each source independently
  - Struggle with the fact that we cannot measure high $P_T$ jet response



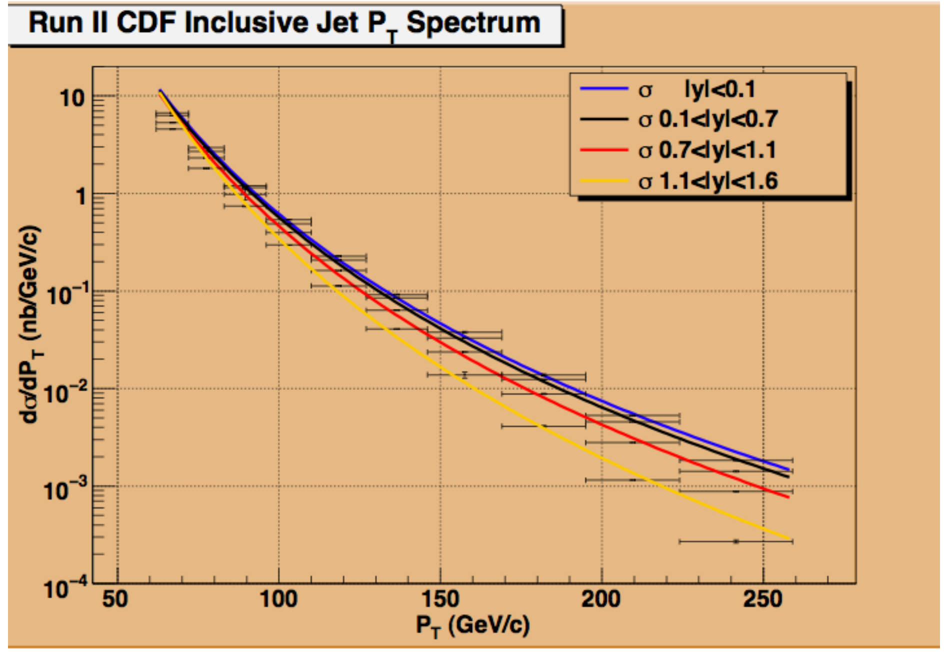
# Production Cross Sections

- Recent CDF analysis of ~1.13 $fb^{-1}$ of jet data
  - Used mid-point algorithm with R=0.7, $f_{merge}$=0.75
  - Data is scaled in plot to avoid overlapping
- Provide a strong test of QCD
  - Theoretically "clean" to model
  - Compare with NLO calculations
    - Fill in details!
  - Generally a trend of small excess of events at higher $P_T$
  - Not statistically significant given systematic uncertainties



# Total Jet Production Rates

- **An "Exercise to Reader" – what is total cross section?**
  - **To answer this question**
    - Fit the spectrum in each y bin to power law using ROOT
    - Use fit to extrapolate over various $P_T$ ranges
      - **Was lazy, only did the first four bins**
    - Generally, differential cross section falls with $(P_T)^{-6}$
      - **And gets a little steeper as $P_T$ increases**
      - **Means that higher $P_T$ jets tend to be more central**
- **Note large cross section at low $P_T$**
  - **This is the source of backgrounds to other objects**
  - **Also note that these are quite uncertain given the extrapolation!**
    - Eg., just changing range of fit
      - $\Delta\sigma(P_T>10)\sim30\%$



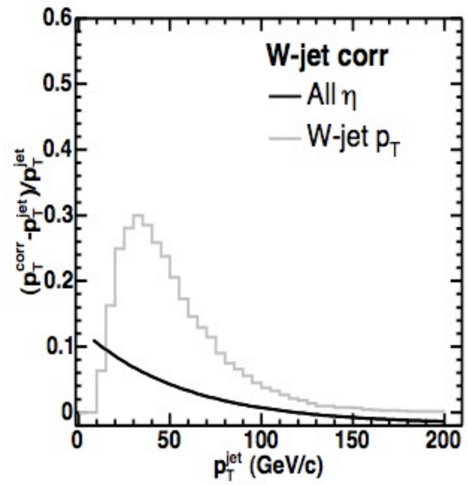

**Cross Section (in nb)**

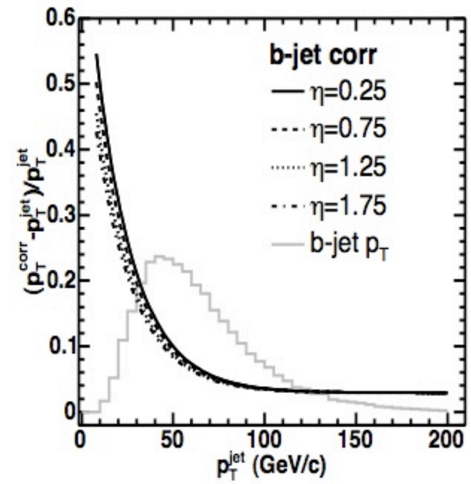
| | PT > 62 GeV | PT > 30 GeV | PT > 10 GeV |
|---|---|---|---|
| \|y\| < 0.1 | 122 | 5,600 | 1,800,000 |
| 0.1 < \|y\| < 0.7 | 111 | 5,600 | 2,000,000 |
| 0.7 < \|y\| < 1.1 | 96 | 6,100 | 3,000,000 |
| 1.1 < \|y\| < 1.6 | 93 | 8,900 | 8,900,000 |
| | 422 | 26,200 | 15,700,000 |

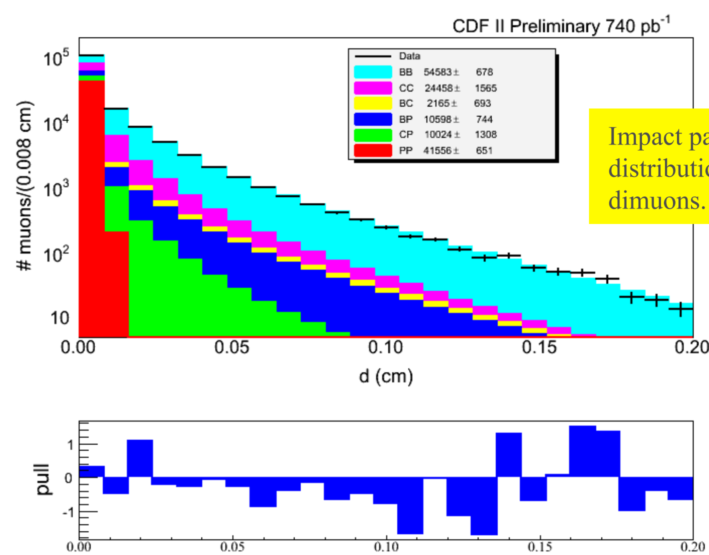
Note: Another ~5-10% in rapidity interval 1.6 < |y| < 2.1

# Heavy Quark Jets

- Heavy quarks (b/c) also manifest themselves as jets
  - Different fragmentation process
  - Different hadronization
    - Result in kinematics that differ from light quark & gluon jets
  - "rich" in $\nu$'s and charged leptons
    - Used for identification
    - But also affect efficiency and & energy resolution
  - Relatively long lifetimes allow for tagging using secondary vertices
    - Become "standard" technique
- Bottom quarks have been particularly important
  - Essential for top quark studies
  - Result in unique capabilities at hadron colliders
    - Good example is $B_s$ studies



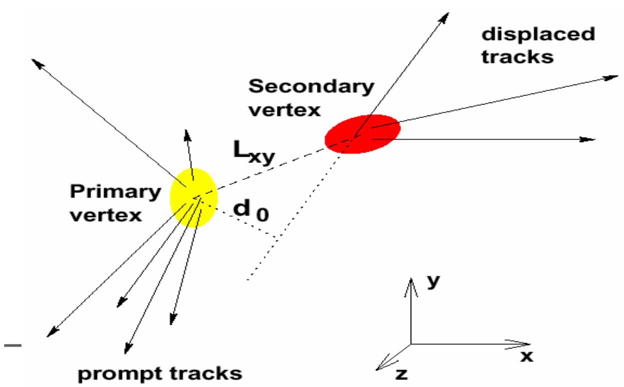




Impact parameter distribution, CDF dimuons.

# Heavy Flavour Tagging

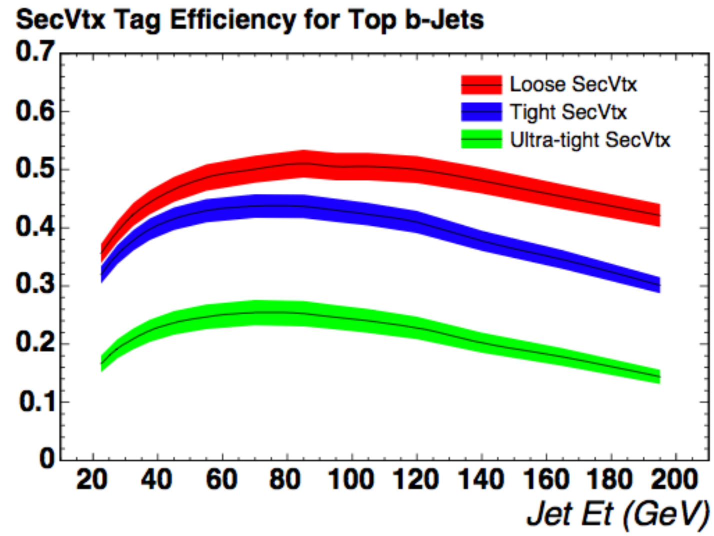
- Heavy flavour tagging has been essential tool at Tevatron
  - Top quark search
  - Search for Higgs
  - Studies of bottom/charm production
- Two methods developed
  - Semileptonic tagging
    - 20% of b's decay inclusively to μ or e
      - Another 20% have leptons from charm decay
    - Challenge is purity of tagging scheme
      - CDF couldn't get fake rates below about 3-4%
  - Secondary vertex tagging most powerful

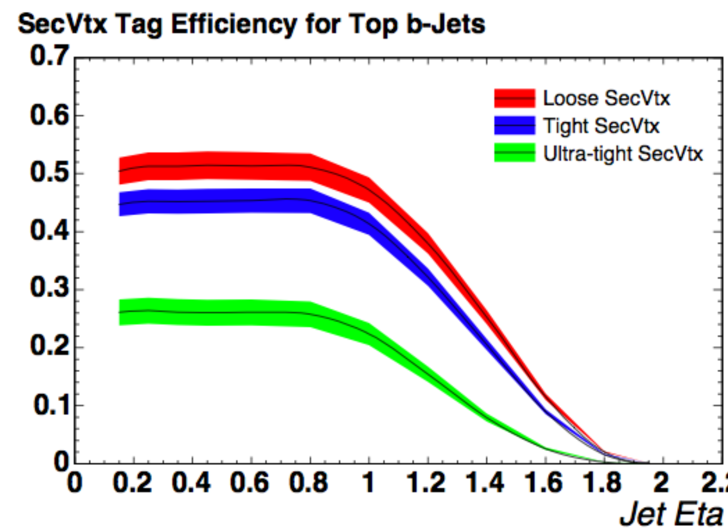
- Basic strategy is to use well-measured tracks
  - Select those with large impact parameter
    - Typically reconstruct average primary beam position in (x,y)
  - Require 2+ tracks with impact parameter > 2s and high quality
    - Attempt to create a secondary vertex
    - If successful, see if secondary vertex is sufficiently far from primary
      - Tag when secondary vtx found
      - Also "fake tag" when tag found, but in wrong direction

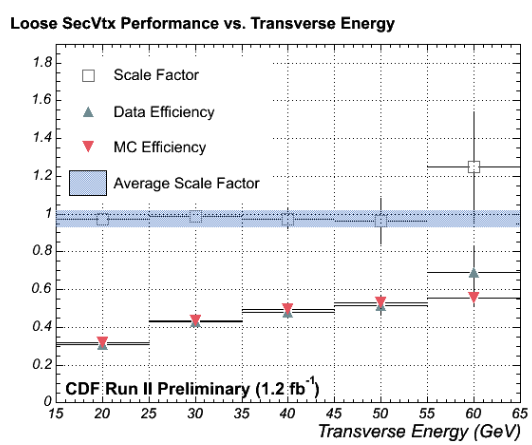


# Tagging Efficiencies

- Tagging efficiency difficult to model via simulation
  - **Requires excellent knowledge of tracking resolution & efficiency**
  - **Strategy:**
    - Measure efficiency and “mistag” rates in data
      - **Inclusive electrons and muons**
        - **Estimate b quark fraction**
      - **Tag fully reconstructed Bs**
    - Compare with simulation & compute a scale factor
      - **SF = $\varepsilon_{Data}/\varepsilon_{MC}$ ~ 0.95 ± 0.05 for “tight” SECVTX**

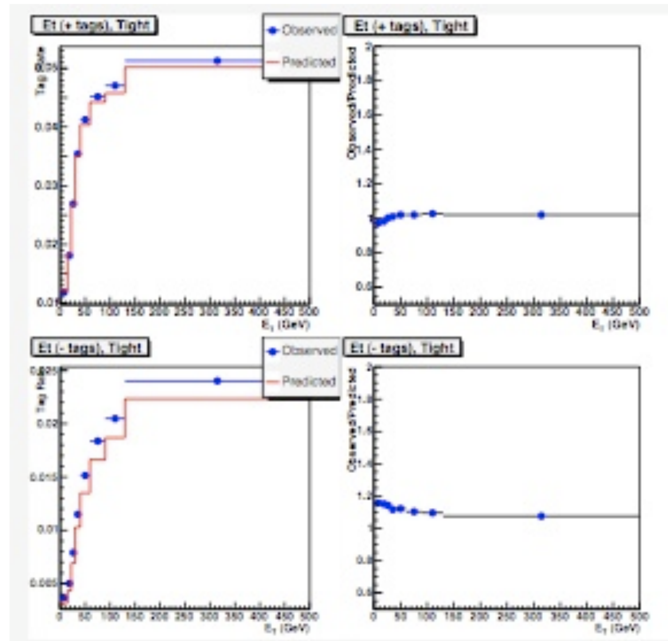


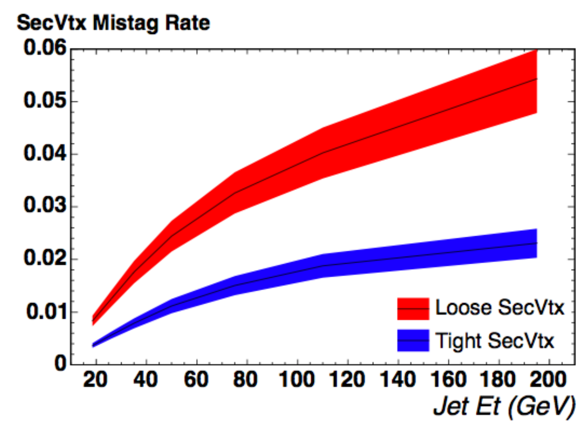



# Tagging Fake Rates

- **B tagging fake rates measured from data**
  - **Take samples of dijet data, and then create a "fake matrix"**
    - Function of 6 variables
    - Measure both +ve and -ve tag rates for "taggable jets"
      - **Use -ve tag rates as mistag rate**
    - Apply mistag rate to the jets in data sample before tagging



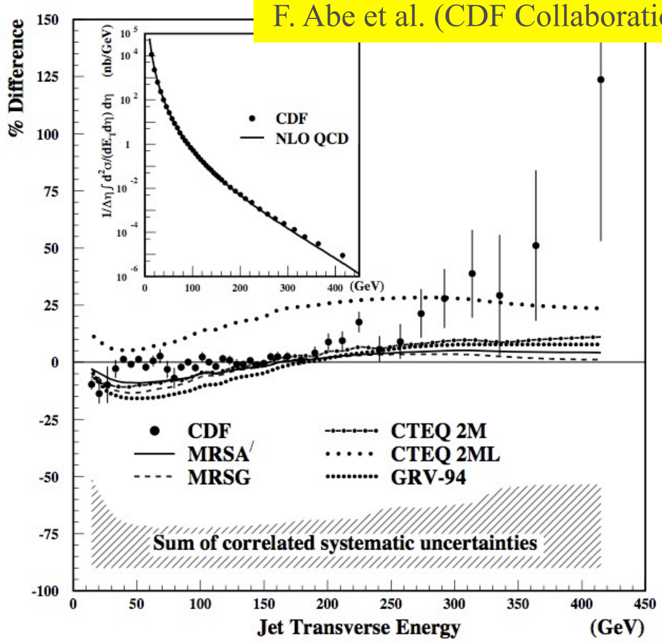

| Variable | Bin Edges |
|---|---|
| Jet $E_T$ (GeV) | 0, 15, 22, 30, 40, 60, 90, 130, 1000 |
| Num Tracks/jet | 0, 1, 2, 3, 4, 5, 6, 7, 8 10, 13, 100 |
| $\lvert\eta_{\rm jet}\rvert$ | 0.0, 0.4, 0.8, 1.1, 2.4 |
| $n_{\rm Primary Vertex}$ | 1 – 6 |
| $\sum E_t$ (GeV) | 0, 80, 140, 220, 1000 |
| $z_{\rm prim}$ (cm) | -25, -10, 10, 25 |

Table 1: Variables and binning used in the mistag matrices.

# Example:  Quark Substructure

- **Search for quark substructure a long-standing tradition at high energies**
  - **Eichten, Lane & Peskin**
    - PRL 50, 811 (1983)
  - **Introduced "contact term" $\Lambda_C$**
  - **CDF obliged in 1996**
    - $\Lambda_C \sim 1.6$ TeV

F. Abe et al. (CDF Collaboration), PRL 77, 438 (1996)



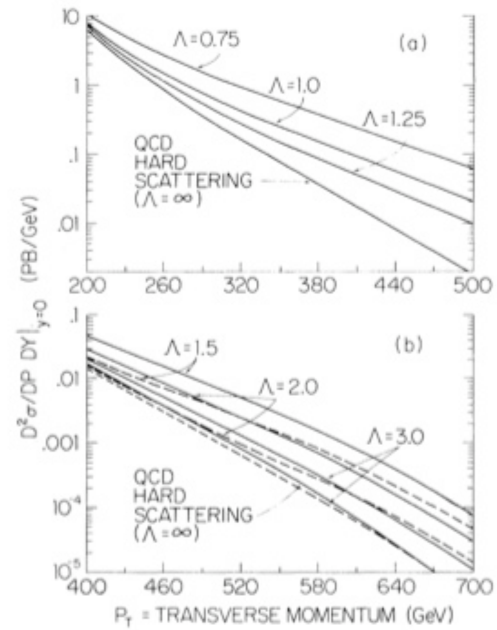




FIG. 3. The jet production cross section (in picobarns/gigaelectronvolt) at rapidity $y = 0$ vs transverse momentum at $\sqrt{s} = 2$ TeV in (a) $\bar{p}p$ collisions and (b) $pp$ collisions for various $\Lambda$ (in teraelectronvolts). The solid and dashed lines in (b) refer, respectively, to the plus and minus signs in Eq. (5). As a result of a cancellation near $y = 0$, the interference is negligible in (a).

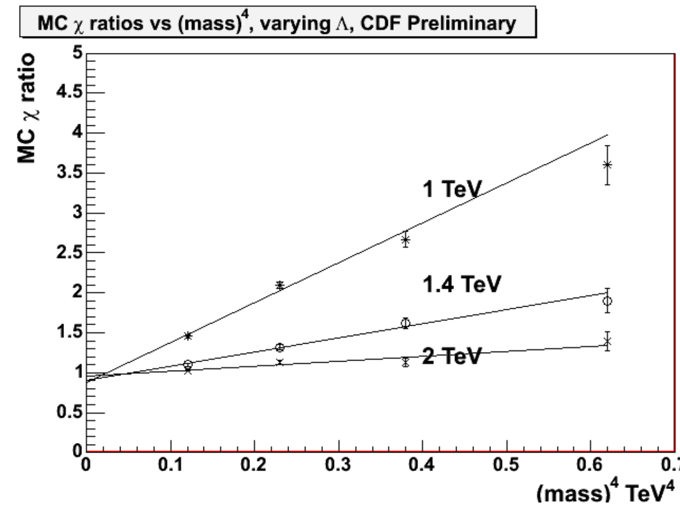
- **Later shown to be described by different PDF behaviour at large x**

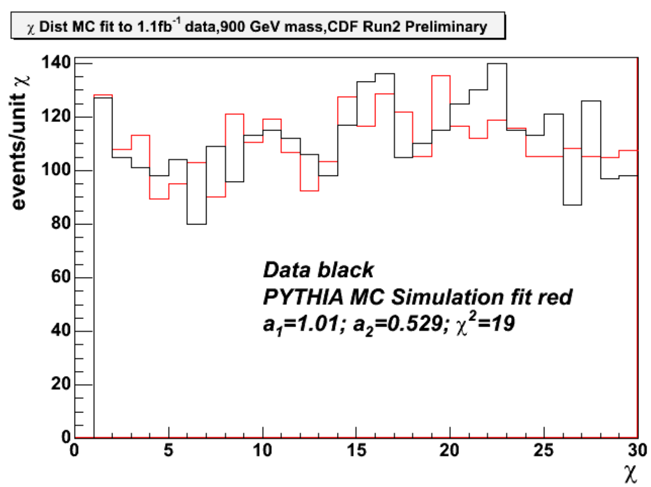
# More Sensitive Study

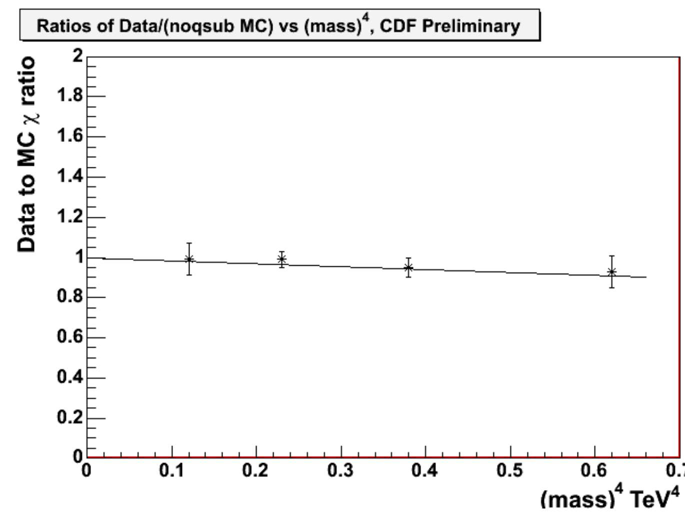
- **Employ angular distribution in dijet scattering**

$$\chi \equiv \exp|\eta_1 - \eta_2|$$

  - **Look at this as a function of dijet invariant mass**
    - 100 GeV mass bins
  - **More sensitive to $\Lambda_C$**
    - Less sensitive to PDFs
    - $\Lambda_C > 2.4$ TeV at 95% CL



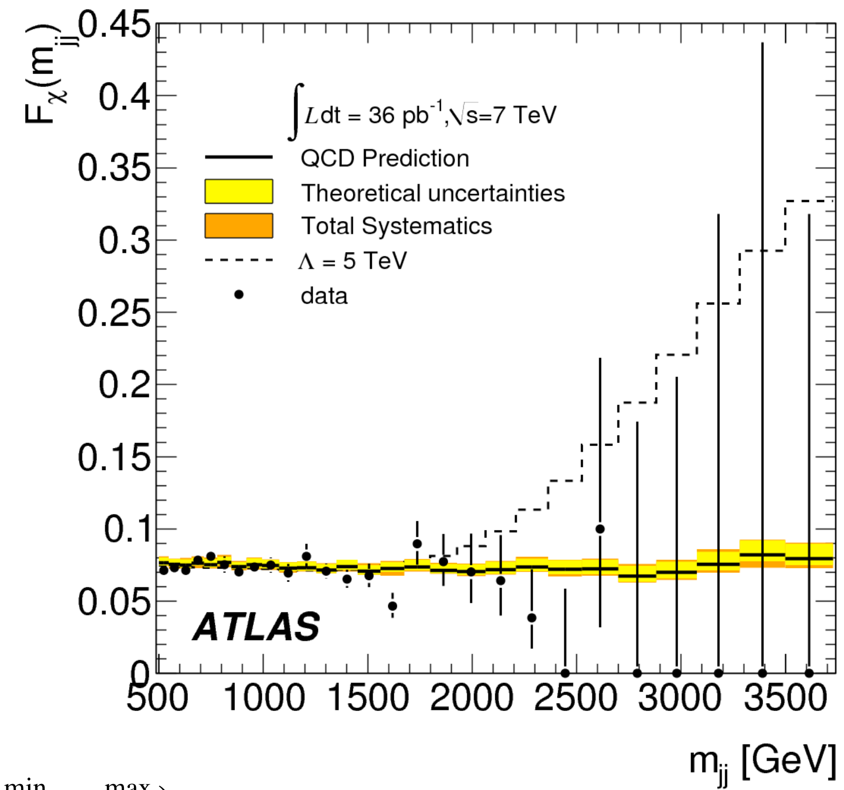




CDF Public Note 9609, November 2008

# And Even More Sensitive!

- **ATLAS has further improved sensitivity**
  - **Look at fraction of centrally produced jets relative to larger angular range**
  - **See behaviour as dijet mass increases**
  - **Expect QCD background to have flat ratio**
- **More sensitive to $\Lambda_C$**
  - **$\Lambda_C > 9.5$ TeV at 95% CL**

ATLAS Collaboration, New Jounr Phys. 13,053004 (2011)




$$F_\chi([m_{jj}^{max} + m_{jj}^{max}]/2) \equiv \frac{N_{events}(|y^*| < 0.6, m_{jj}^{min}, m_{jj}^{max})}{N_{events}(|y^*| < 1.7, m_{jj}^{min}, m_{jj}^{max})}$$

$$y^* \equiv \frac{1}{2}(y_1 - y_2)$$