

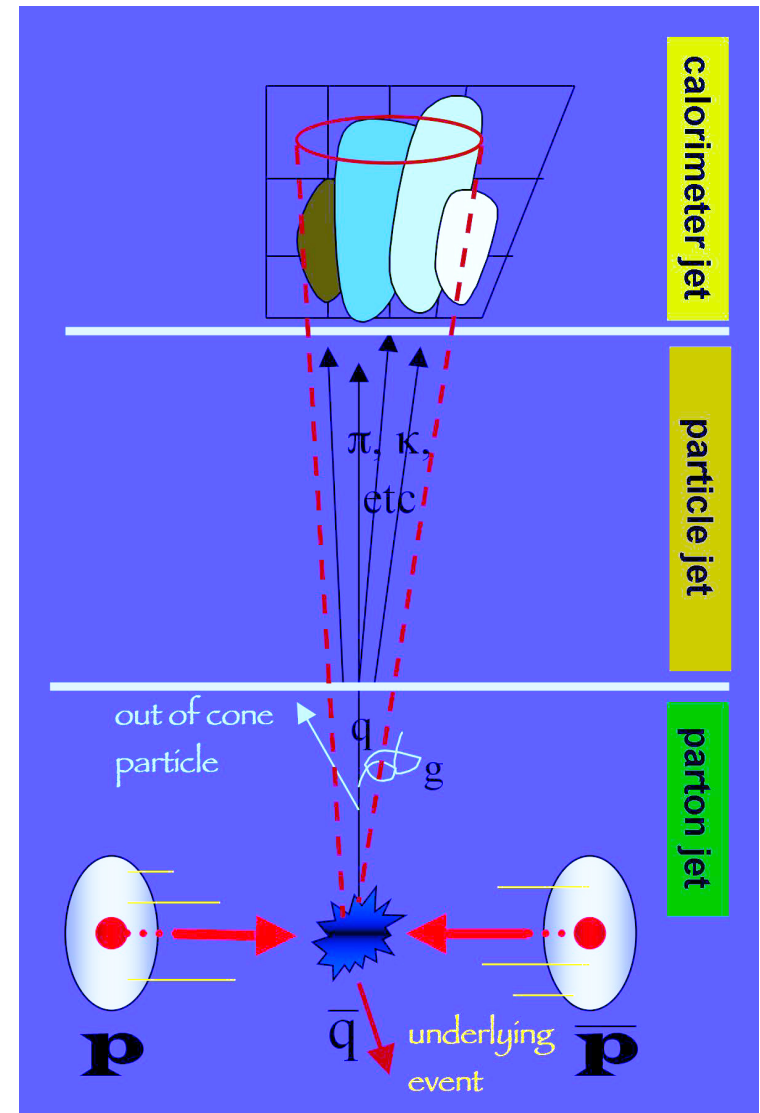
Hard Scattering in Hadron-Hadron Collisions: Physics and Anatomy

Section 4: Production & Identification of Jets

- 1. Definitions of Basic Physics Processes**
- 2. Anatomy of a Jet**
- 3. Jet-Finding Algorithms**
- 4. Resolutions and Efficiencies**
- 5. Heavy Quark Tagging**
- 6. Example: Quark Substructure**

Definitional Issues

- **Confinement in QCD ensures that high P_T quarks & gluons undergo**
 - **Fragmentation -- ie, dissociation into a “jet” of coloured partons**
 - **Hadronization -- ie, the partons form colourless, observable hadrons**
- **Study of jets motivated by**
 - **Understanding QCD**
 - **Studying of heavy quarks**
 - > b/c quarks that fragment & hadronize before decay
 - > Top quarks that decay before fragmentation/hadronization
 - **Searching for new interactions that couple to quarks/gluons**
 - **Background source to e, μ, γ & τ**

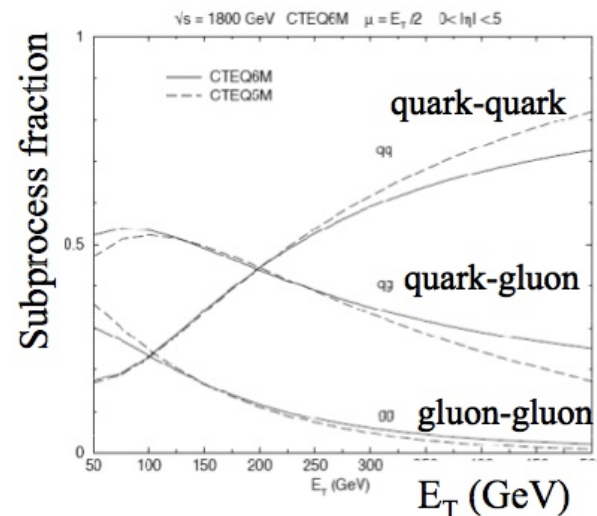
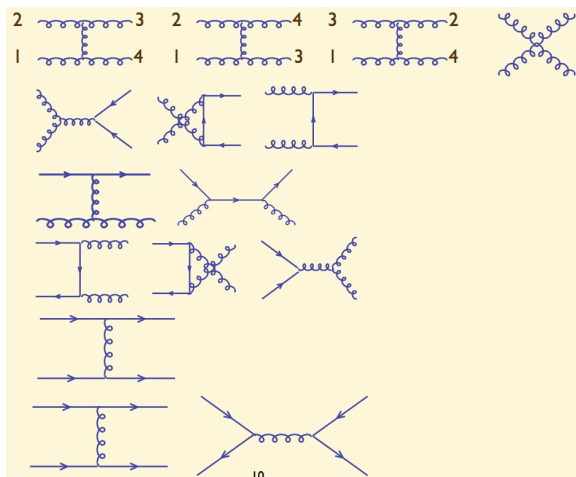


Fundamentals of Jet Physics

- Basic production mechanism in pQCD starts with

$$\sigma = \sum_{\substack{\text{partons } i \\ \text{colour } j}} C_{ij} \int_0^1 d\tau \int_{\tau}^1 \frac{dx_1}{\tau} [f_1(x_1) f_2(\tau/x_1)] \hat{\sigma}(\tau s)$$

- Leading-order (LO) diagrams already complex



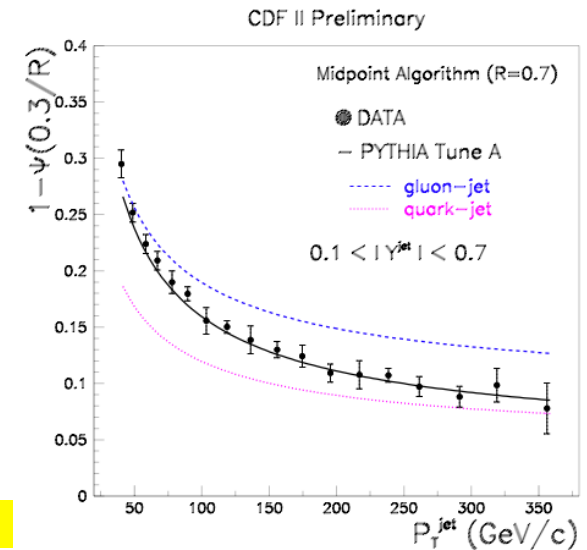
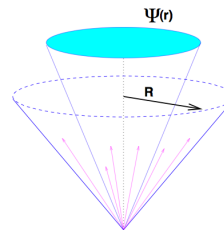
Process	$\bar{\Sigma} \mathcal{M} ^2 / g^4$	Numerical value for 90°
$qq' \rightarrow qq'$	$\frac{4}{9} \frac{s^2 + u^2}{t^2}$	2.22
$q\bar{q}' \rightarrow q\bar{q}'$	$\frac{4}{9} \frac{s^2 + u^2}{t^2}$	2.22
$qq \rightarrow qq$	$\frac{4}{9} \left(\frac{s^2 + u^2}{t^2} + \frac{s^2 + t^2}{u^2} \right) - \frac{8}{27} \frac{s^2}{ut}$	3.26
$q\bar{q} \rightarrow q'\bar{q}'$	$\frac{4}{9} \frac{t^2 + u^2}{s^2}$	2.22
$q\bar{q} \rightarrow q\bar{q}$	$\frac{4}{9} \left(\frac{s^2 + u^2}{t^2} + \frac{t^2 + u^2}{s^2} \right) - \frac{8}{27} \frac{u^2}{st}$	2.59
$q\bar{q} \rightarrow gg$	$\frac{32}{27} \frac{t^2 + u^2}{tu} - \frac{8}{3} \frac{t^2 + u^2}{s^2}$	1.04
$gg \rightarrow q\bar{q}$	$\frac{1}{6} \frac{t^2 + u^2}{tu} - \frac{3}{8} \frac{t^2 + u^2}{s^2}$	0.15
$gq \rightarrow gq$	$\frac{4}{9} \frac{s^2 + u^2}{su} + \frac{u^2 + s^2}{t^2}$	6.11
$gg \rightarrow gg$	$\frac{9}{2} \left(3 - \frac{tu}{s^2} - \frac{su}{t^2} - \frac{st}{u^2} \right)$	30.4

What Have We Learned?

- **Definition of jets critical**
 - Much evolution in algorithms
 - Driven in large measure by theoretical considerations
- **Calibration of jets requires data-driven techniques**
 - Developed several techniques to calibrate *in situ*
 - Still “work in progress”
- **Approach to jet-finding and calibration driven by physics**
 - Best example is comparison between
 - > QCD tests
 - > Reconstruction of heavy objects (top and Higgs)

- **Need data to understand jets as backgrounds**
 - Examples include
 - > Lepton ID
 - > MET measurement
 - > Heavy quark tagging
 - Use to “calibrate” MC/simulation
- **Bottom line: SM Picture of QCD works well**

$$\Psi(r) \equiv \frac{1}{N_{jet}} \sum_{jets} \frac{P_T(0,r)}{P_T(0,R)}$$

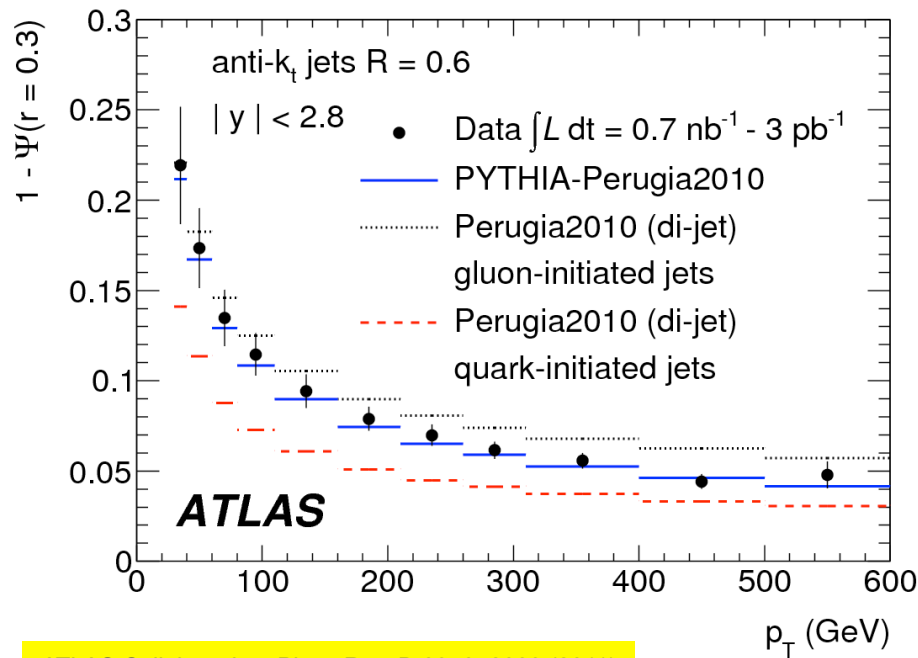


D. Acosta et al. (CDF), Phys. Rev. D 71, 112002 (2005)

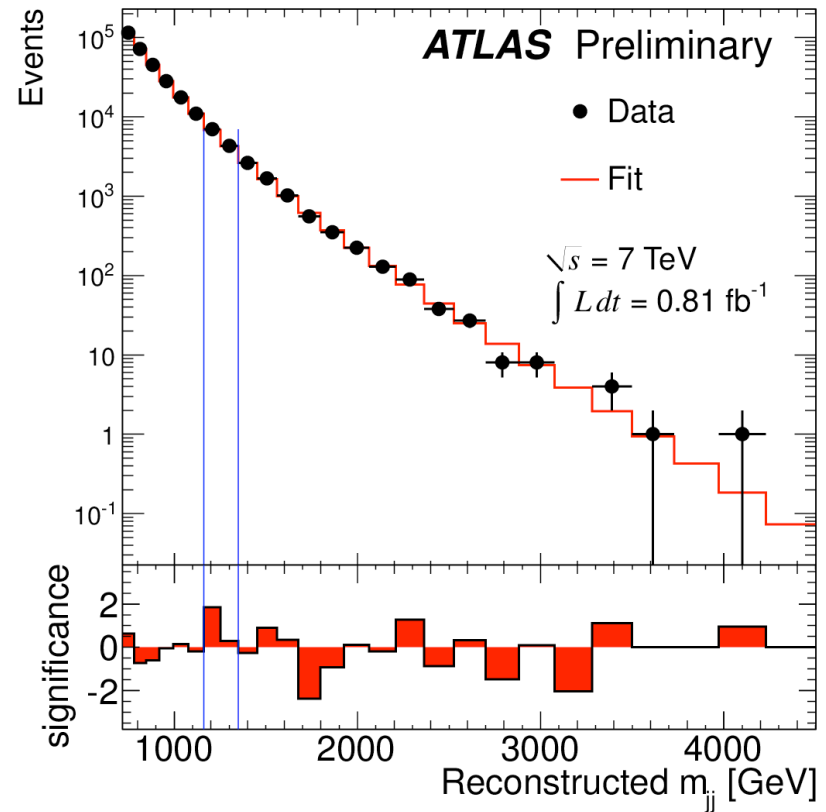
LHC Lessons

■ LHC studies have reproduced many of these effects

- **However, much higher jet momenta**
 - > Jets with $p_T > 2$ TeV
- **Focus has been on searches**
 - > Looking for resonances in dijet mass
 - > Sensitive to excited quarks

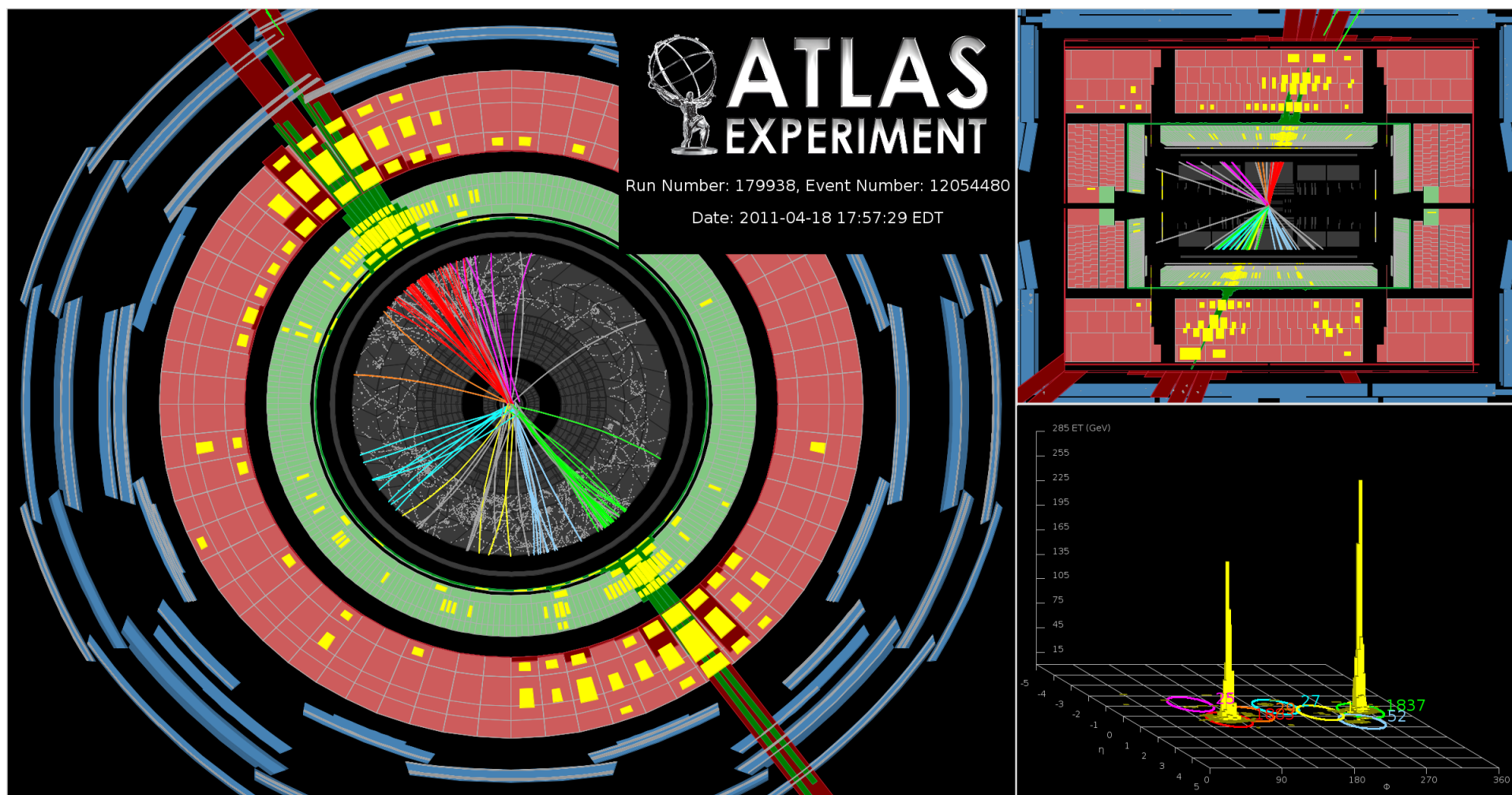


ATLAS Collaboration, Phys. Rev. D 83, 052003 (2011)



ATLAS-CONF-2011-096, 18 July 2011

Highest Dijet Mass Event



Jet Anatomy

■ A jet arises from 2 different physical phenomena

- Happen at different energy scales
 - > Fragmentation of initial parton
 - QCD radiation of a coloured object
 - Creates a “cluster” of coloured partons
 - In principle, not independent of rest of event
 - Energy scale $\gg 1$ GeV
 - > Hadronization of “cluster”
 - Formation of colourless objects -- mesons & baryons
 - Responsible for the real observables
 - Energy scale ~ 1 GeV

■ Have to worry about

- What defines a jet (algorithm)?
- What its properties are (recombination scheme)?

■ First, tackle easiest part: What is a jet's observable properties?

- Assume you have a collection of final state mass-less “particles” detected in calorimeter towers i

$$\vec{p}_J \equiv (E^J, p_x^J, p_y^J, p_z^J) \equiv \sum_i (E^i, p_x^i, p_y^i, p_z^i)$$

$$p_T^J \equiv \sqrt{(p_x^J)^2 + (p_y^J)^2}$$

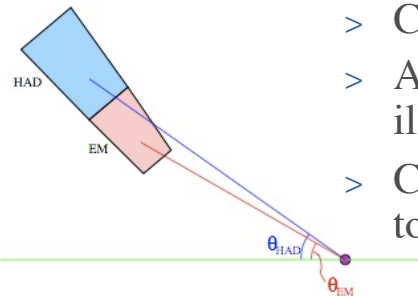
$$M^J \equiv \sqrt{(E^J)^2 - (p^J)^2}$$

$$y^J \equiv \frac{1}{2} \ln \frac{E^J + p_z^J}{E^J - p_z^J}$$

$$\varphi^J \equiv \tan^{-1} \frac{p_y^J}{p_x^J}$$

– Advantages:

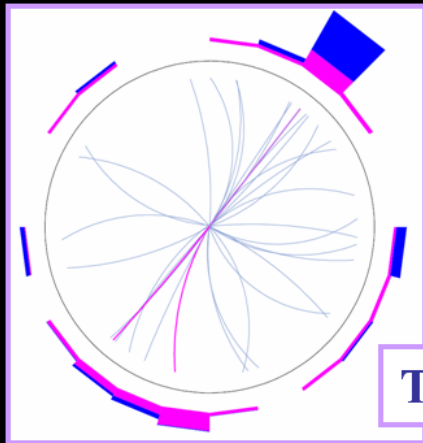
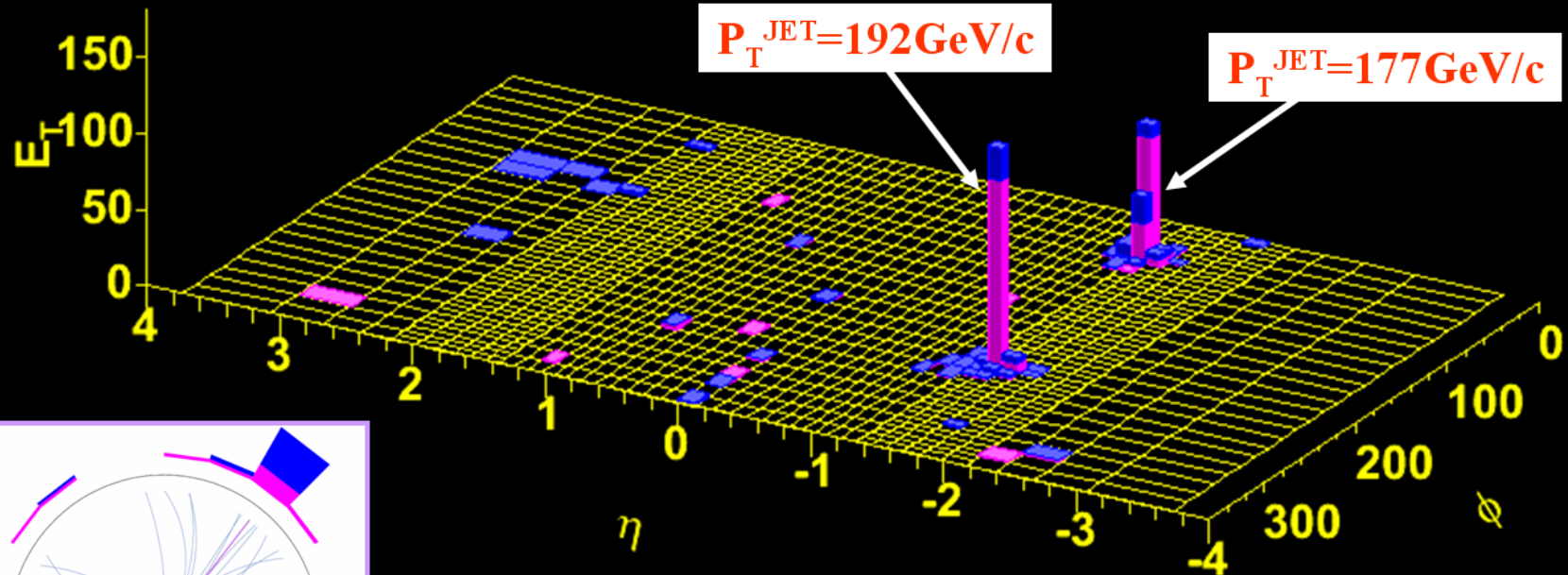
- > Clear Lorentz behaviour
- > Avoids use of E_T which has ill-defined definition
- > Can generalize to “cells”, towers, charged particles, etc.



A Real Jet Event

Tower $E_T > 0.5$ GeV

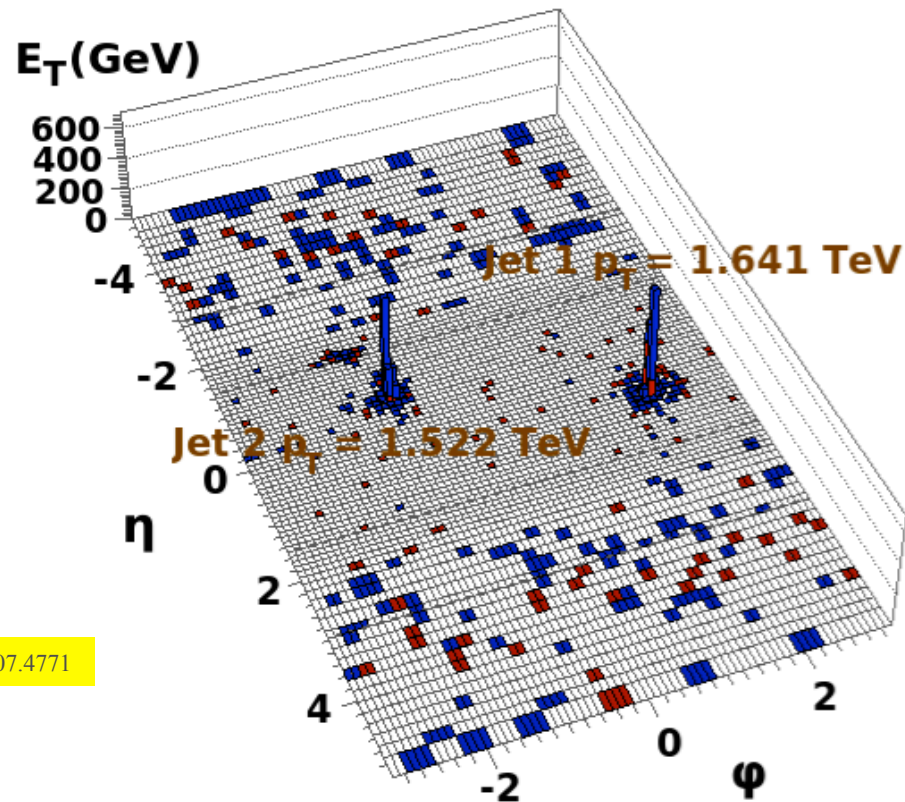
K_T D=0.7: Raw P_T^{JET}



Track $p_T > 0.5$ GeV/c

CDF RUN II
Run 163064
Event 6753986

An LHC CMS Jet Event



Highest dijet
Mass event
Observed by
CMS at 7 TeV

CMS Collaboration, hep-ex/1107.4771

Parton Shower Evolution

- Start with a parton (q/g) with virtuality μ^2

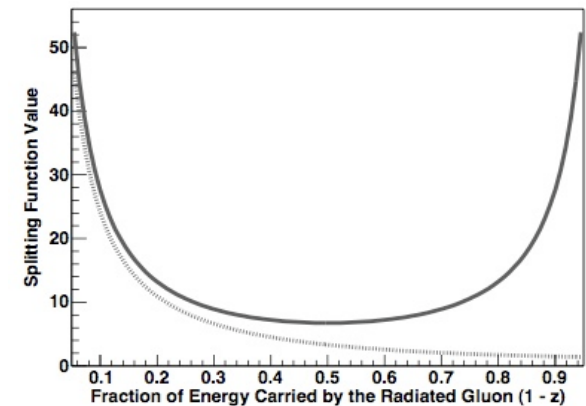
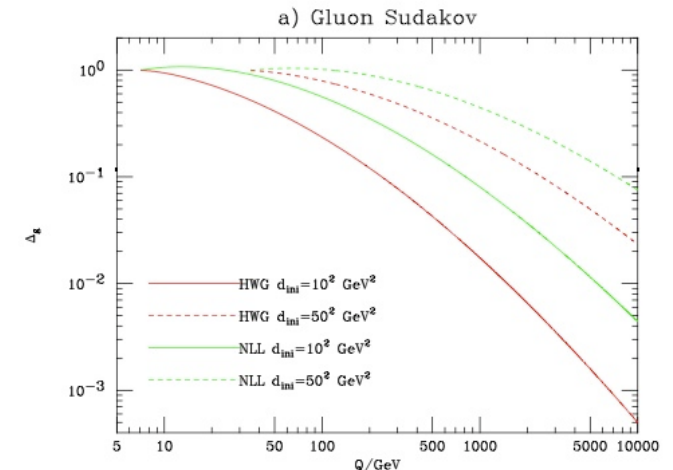
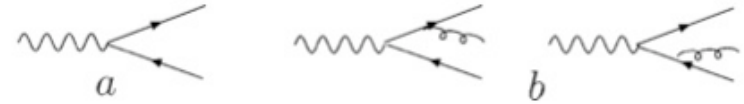
- Probability of emission with daughter carrying z fraction of parent momentum

$$d^2 P_a(z, \mu^2) = \frac{d\mu^2}{\mu^2} \frac{\alpha_s}{2\pi} P_{a \rightarrow bc}(z) dz$$

- Order these using Sudakov factor, relating $\mu^2 \sim Q^2$

$$P_a^{no}(Q_{\max}^2, Q^2) = \exp\left(-\int_{Q^2}^{Q_{\max}^2} dQ' \int_{z_{\min}}^{z_{\max}} dz' P_a(z', Q'^2)\right)$$

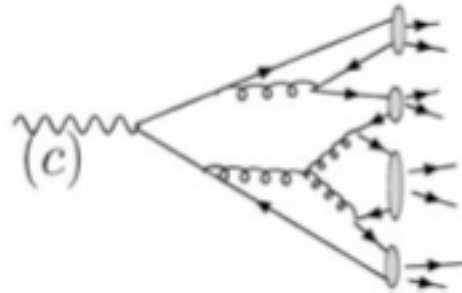
- Deal with infrared & collinear divergences
 - > Define minimum $\mu - \mu_0$
- Ensure colour coherence of multiple emissions
 - > Typically do this by angular ordering, selective vetoing, etc.
 - > Must be respected when hadronization is performed



Hadronization of Showers

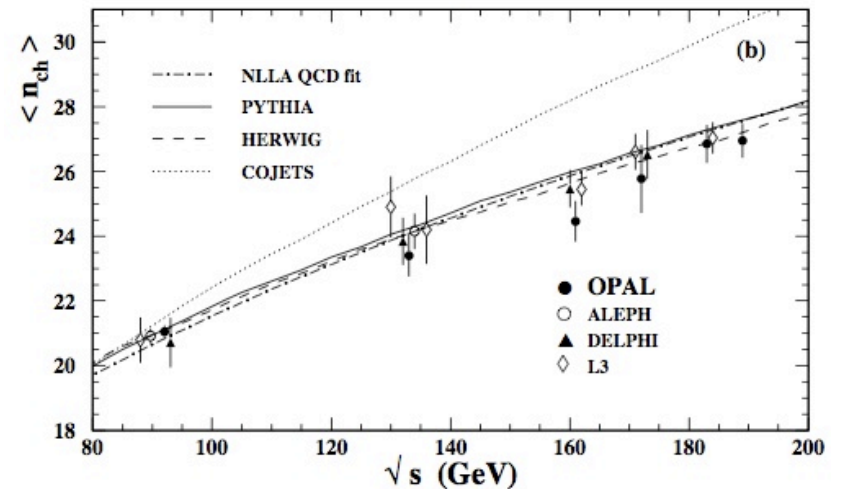
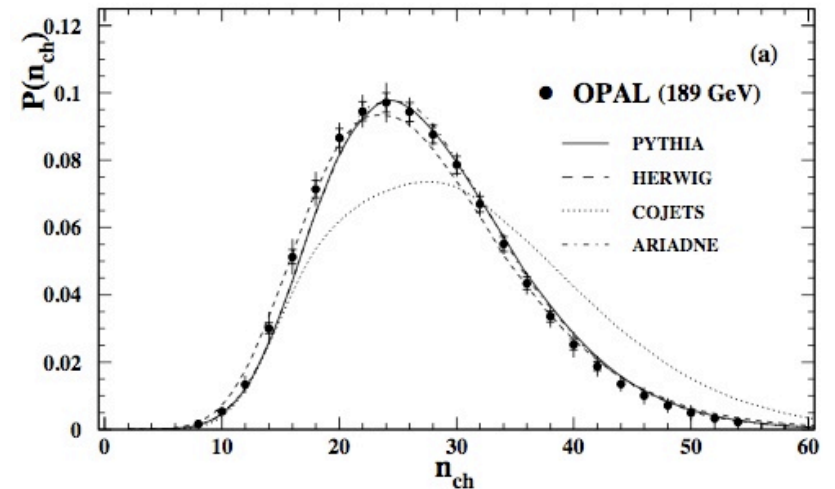
Hadronization is then performed

- Invoke “parton-hadron duality”
- Several models
 - > String fragmentation (eg., PYTHIA)
 - > Cluster fragmentation (eg. HERWIG)



- Have various parameters that need to be tuned to data
 - > Best constraints from LEP
 - Tevatron results confirm these, but don't really add much power
 - Challenging to measure without significant systematics
 - > Remains a source of systematic uncertainty

OPAL, Eur. Phys. J C16, 185 (2000)



Jet Algorithms

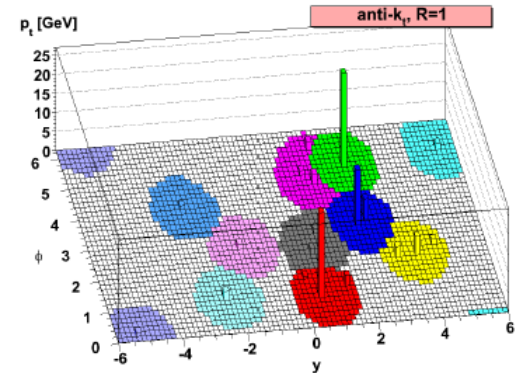
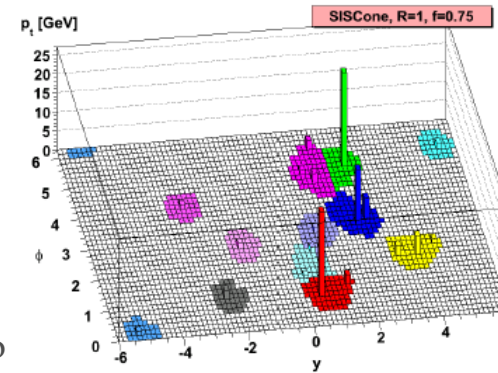
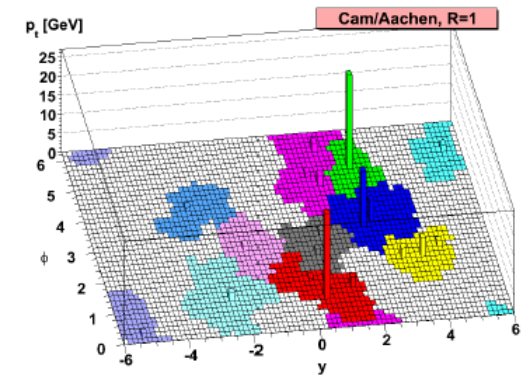
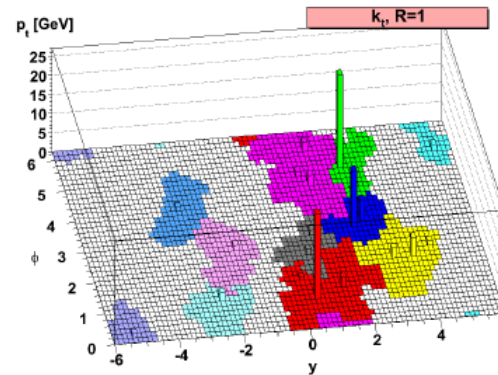
- **Jet clustering algorithms have been focus of much effort**
 - **Goals of any algorithm can be divided into**
 - > **Theoretically motivated:**
 - **Fully specified**
 - **Detector independent**
 - **Theoretically well-behaved**
 - **Order independent**
 - > **Experimentally motivated:**
 - **Fully specified**
 - **Detector independent**
 - **Optimal resolution and efficiency**
 - **Ease of calibration**
 - **Computationally efficient**
- **Various efforts to develop consistent frameworks**
 - **Snowmass Accord (1990)**
 - **Les Houches Accord (1999)**
- **Raz Alon (see talk below) has done a nice job of summarizing current Jet Algorithm codes**
 - **Key observations:**
 - > **In principle, prefer some algorithms over others**
 - **Seedless cone-based algorithms**
 - **K_T algorithms**
 - > **Computational efficiency is a concern in some cases**
 - **But largely an issue of optimization**
 - > **Selection of “best” algorithm requires evaluation of ultimate systematic uncertainties**
 - **Need data, as certain choices will depend on performance of calorimeter**
 - **Example is noise and pileup**
 - **Good news is that we are not limited by lack of ideas**

R. Alon, <http://indico.cern.ch/conferenceDisplay.py?confId=52628>

Clustering Effects

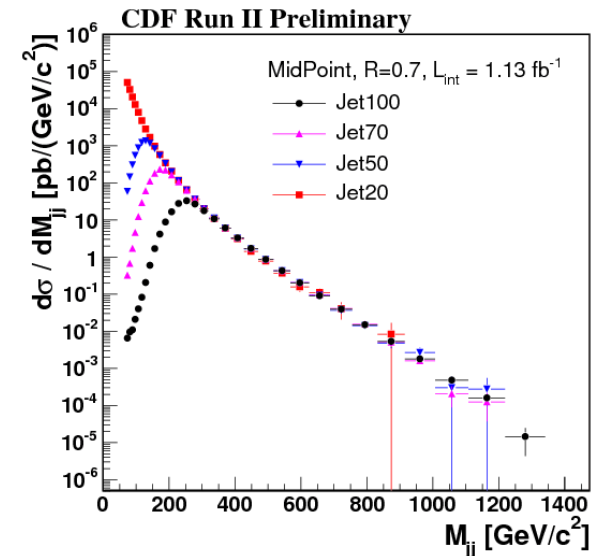
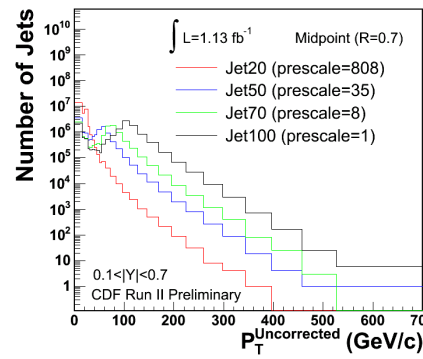
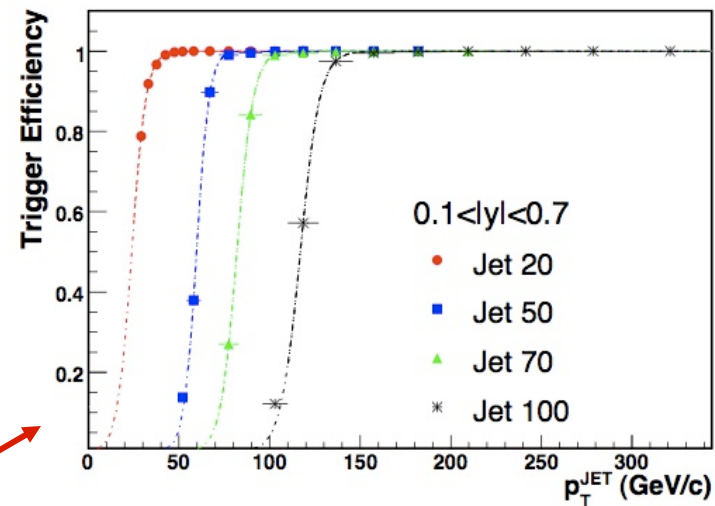
■ Illustrate by one example (from ATLAS studies)

- Compare results of several different algorithms
 - > K_T with $R=1$
 - > Angular-ordering (Cam/Aachen)
 - > SIScone
 - > Anti- K_T
- Things to be concerned about
 - > Cluster sizes determined by data will present challenges to calibrate
 - > Cluster merging/splitting will continue to be a challenge
 - > Optimization of resolution/systematic uncertainties will require effort
- Things not to worry about
 - > Angular resolution (though need to check for any biases)!



Jet Finding Efficiencies

- **Efficiency of finding jets limited primarily by two effects:**
 - Detector energy response & resolution
 - Physical size of jets
 - > For cone algorithms, these two compete with each other
- **Further complicated by the fact that jets are produced with sharply falling spectrum**
 - Means that efficiencies become an issue already at the trigger level
 - Manage these at Tevatron with variety of triggers
 - > Prescale lower-energy jet triggers
 - > Lower energy jets used primarily for
 - Background studies
 - Calibration



Jet Energy Resolutions

- MC + simulation give estimates of energy resolution

- Resolution is determined primarily by convolution of
 - > Intrinsic calorimeter response
 - > Jet fragmentation & hadronization effects
 - > Jet algorithm + pileup +
- In reality, need to measure the resolution in data

- Four *in situ* measurements of resolution developed at Tevatron

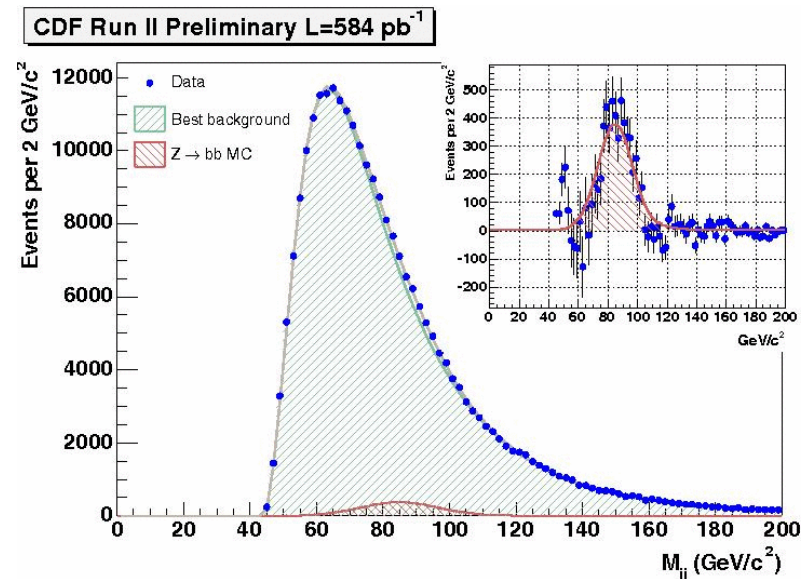
- γ +jet balancing
- W to qq in top quark decays
- Dijet balancing (more of a constraint than anything else)
- Z to bb decays
 - > Require two jets, each with secondary vertex b-tag
 - Possible due to L2 vertex trigger

- Taking the FWHM $\sim 25 \text{ GeV}/c^2$, obtain

$$\sigma_Z \sim 12\% M_Z$$

$$\Rightarrow \frac{\sigma}{P_T^J} \sim 17\%$$

- Or about 50% more than intrinsic energy resolution of calorimeter

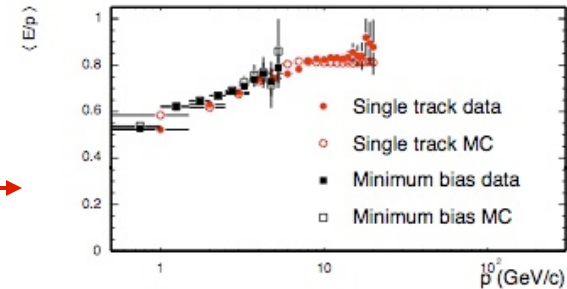


Jet Energy Calibration

■ To calibrate jet energy scale:

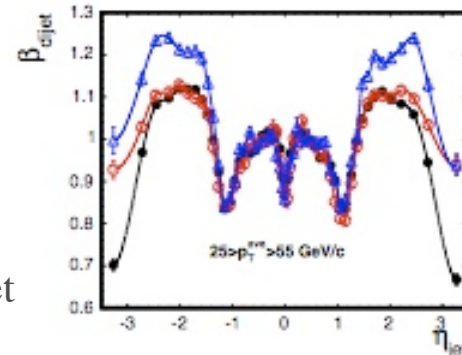
– 1. Determine intrinsic response to particles

- > Combination of in situ measurements & test beam data



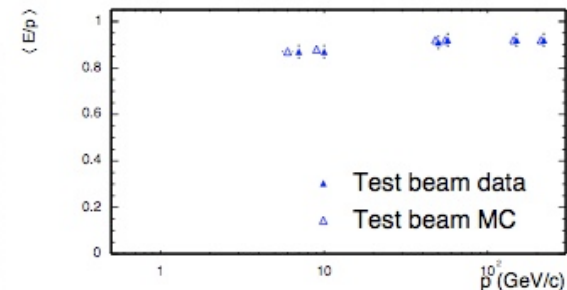
– 2. Dijet balancing to get uniform η response

- > Primarily dijet data
- > “Tune” MC and simulation



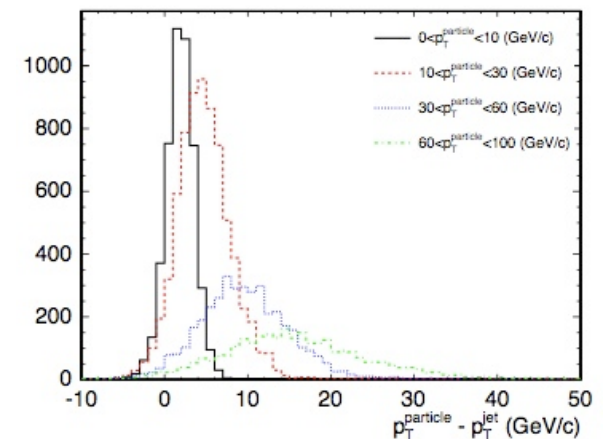
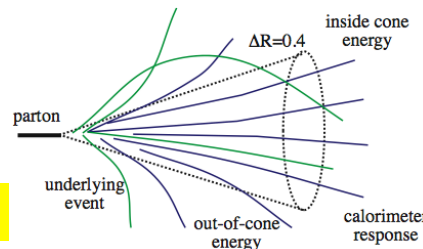
– 3. Determine absolute response to “particle jet”

- > Define particle jet as all real particles in cone of jet
- > Account for calorimeter nonlinearity, cracks, etc.



– 4. Take into account “out-of-cone” effects, multiple interactions

- > Use combination of MC and data

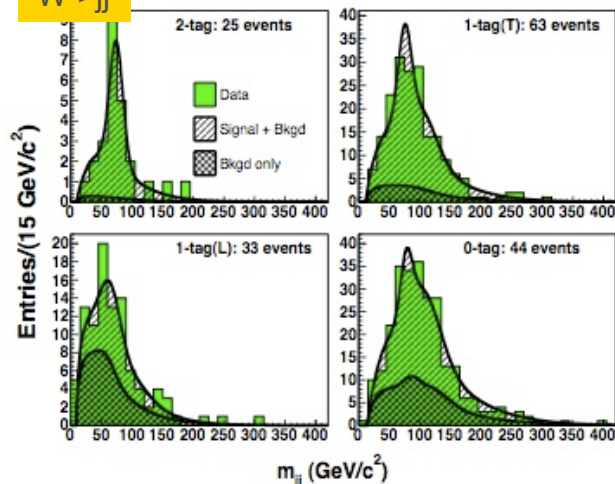


Final Steps in Energy Calibration

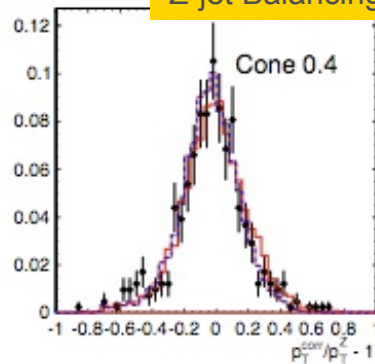
■ Cross check using, for example,

- Z+jet & γ +jet balancing
- Dijet balancing
- W \rightarrow jj in ttbar events

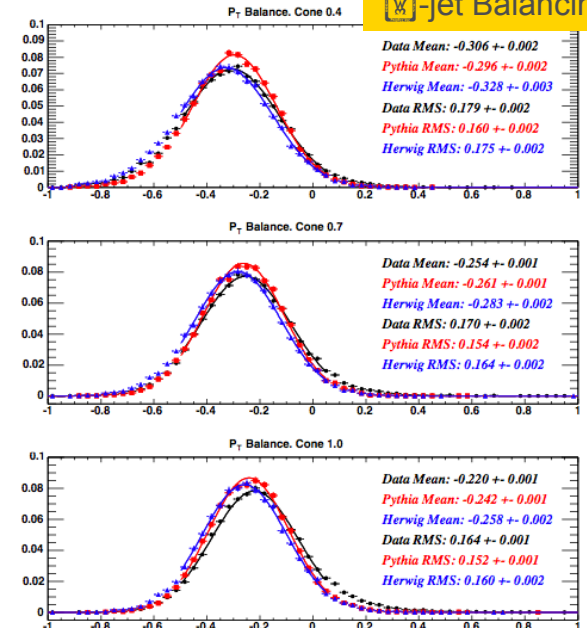
W \rightarrow jj



Z-jet Balancing

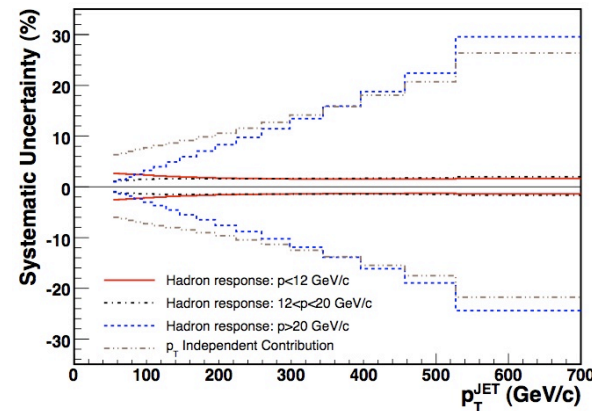


γ -jet Balancing



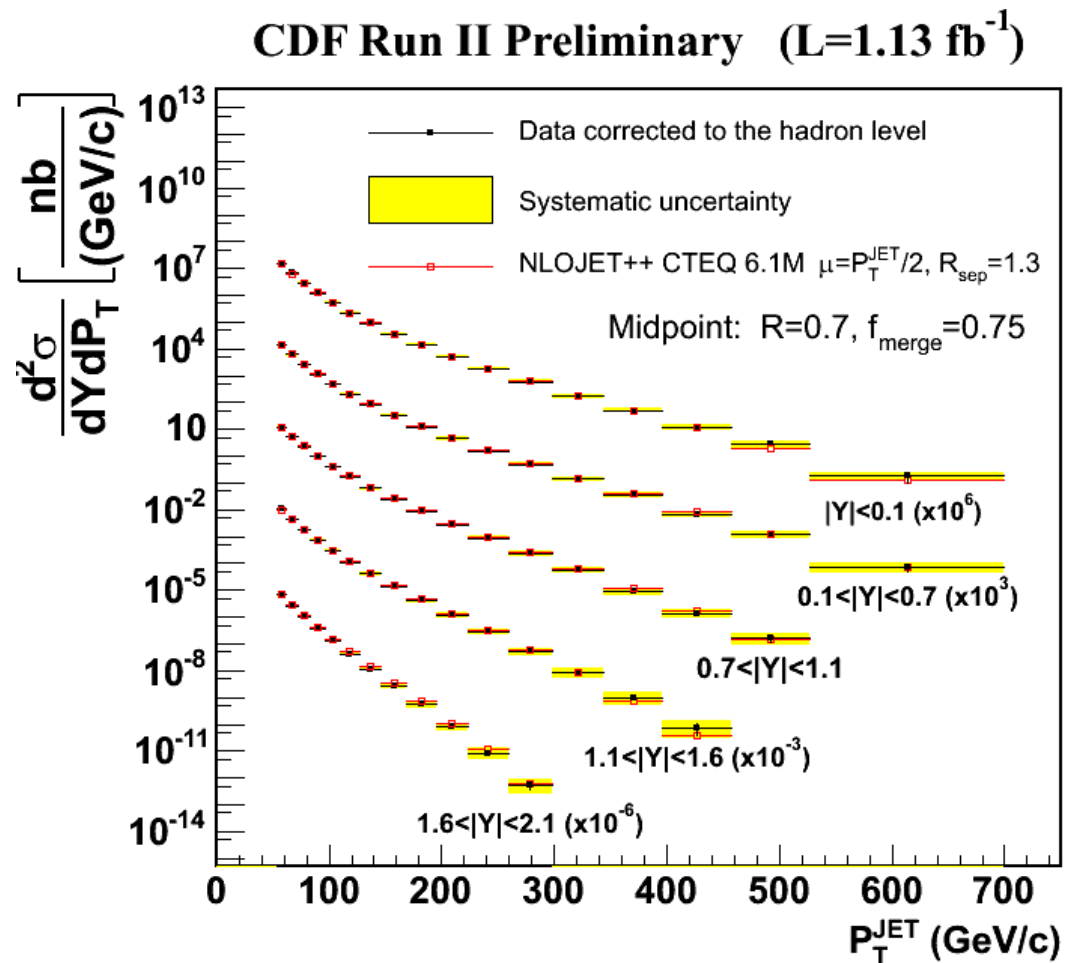
■ Estimate systematic uncertainties

- Estimate each source independently
- Struggle with the fact that we cannot measure high P_T jet response



Production Cross Sections

- CDF analysis of $\sim 1.13 \text{ fb}^{-1}$ of jet data
 - Used mid-point algorithm with $R=0.7$, $f_{\text{merge}}=0.75$
 - Data is scaled in plot to avoid overlapping
- Provide a strong test of QCD
 - Theoretically “clean” to model
 - Compare with NLO calculations
 - > Fill in details!
 - Generally a trend of small excess of events at higher P_T
 - Not statistically significant given systematic uncertainties



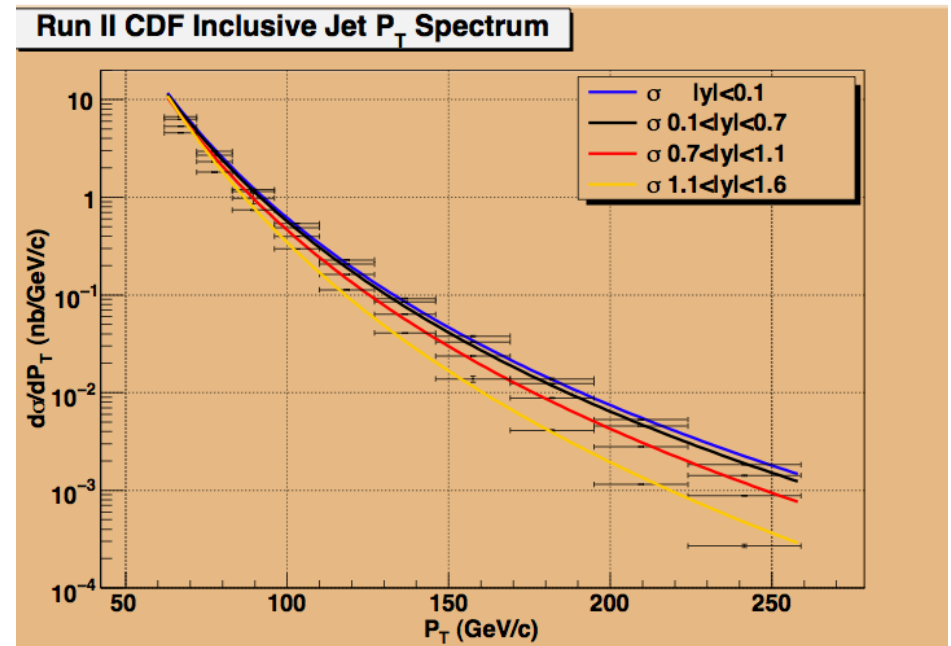
Total Jet Production Rates

■ An “Exercise to Reader” – what is total cross section?

- To answer this question
 - > Fit the spectrum in each y bin to power law using ROOT
 - > Use fit to extrapolate over various P_T ranges
 - Was lazy, only did the first four bins
 - > Generally, differential cross section falls with $(P_T)^{-6}$
 - And gets a little steeper as P_T increases
 - Means that higher P_T jets tend to be more central

■ Note large cross section at low P_T

- This is the source of backgrounds to other objects
- Also note that these are quite uncertain given the extrapolation!
 - > Eg., just changing range of fit
 - $\Delta\sigma(P_T > 10) \sim 30\%$



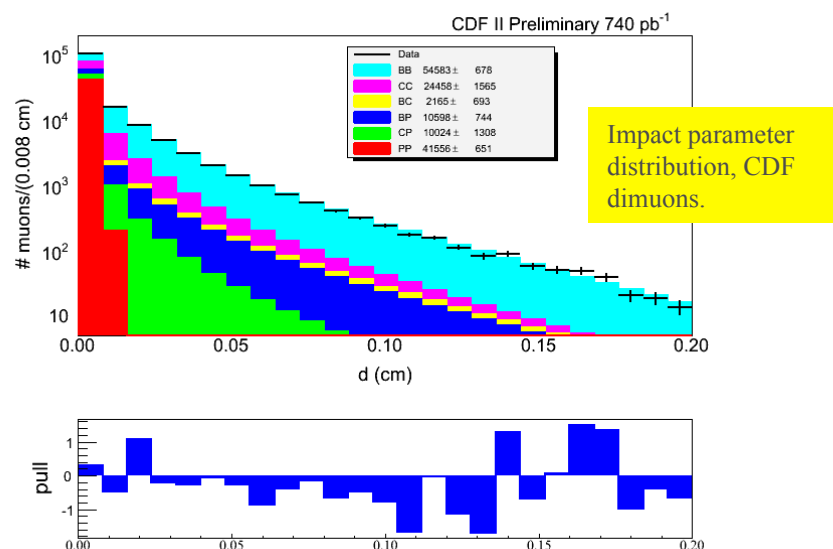
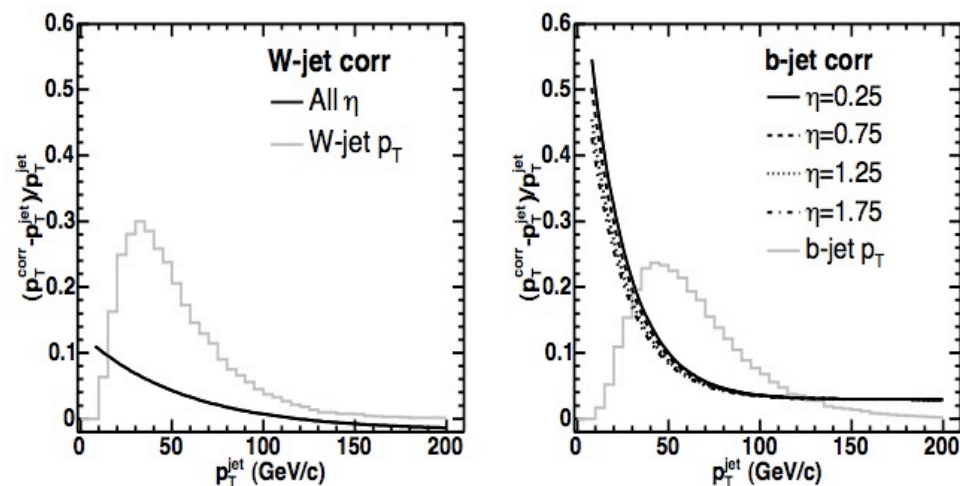
Cross Section (in nb)

	PT > 62 GeV	PT > 30 GeV	PT > 10 GeV
$ y < 0.1$	122	5,600	1,800,000
$0.1 < y < 0.7$	111	5,600	2,000,000
$0.7 < y < 1.1$	96	6,100	3,000,000
$1.1 < y < 1.6$	93	8,900	8,900,000
	422	26,200	15,700,000

Note: Another ~5-10% in rapidity interval $1.6 < |y| < 2.1$

Heavy Quark Jets

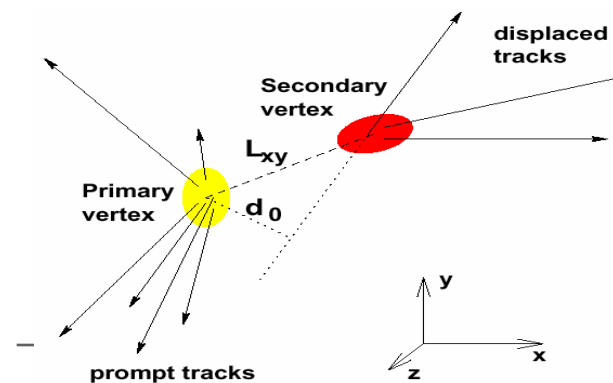
- **Heavy quarks (b/c) also manifest themselves as jets**
 - **Different fragmentation process**
 - **Different hadronization**
 - > Result in kinematics that differ from light quark & gluon jets
 - **“rich” in ν 's and charged leptons**
 - > Used for identification
 - > But also affect efficiency and energy resolution
 - **Relatively long lifetimes allow for tagging using secondary vertices**
 - > Become “standard” technique
- **Bottom quarks have been particularly important**
 - **Essential for top quark studies**
 - **Result in unique capabilities at hadron colliders**
 - > Good example is B_s studies



Heavy Flavour Tagging

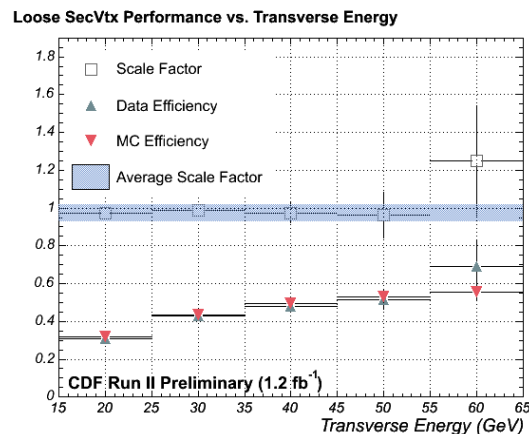
- Heavy flavour tagging has been essential tool at Tevatron
 - Top quark search
 - Search for Higgs
 - Studies of bottom/charm production
- Two methods developed
 - Semileptonic tagging
 - > 20% of b' 's decay inclusively to μ or e
 - Another 20% have leptons from charm decay
 - > Challenge is purity of tagging scheme
 - CDF couldn't get fake rates below about 3-4%
 - Secondary vertex tagging most powerful

- Basic strategy is to use well-measured tracks
 - Select those with large impact parameter
 - > Typically reconstruct average primary beam position in (x,y)
 - Require 2+ tracks with impact parameter $> 2s$ and high quality
 - > Attempt to create a secondary vertex
 - > If successful, see if secondary vertex is sufficiently far from primary
 - Tag when secondary vtx found
 - Also “fake tag” when tag found, but in wrong direction

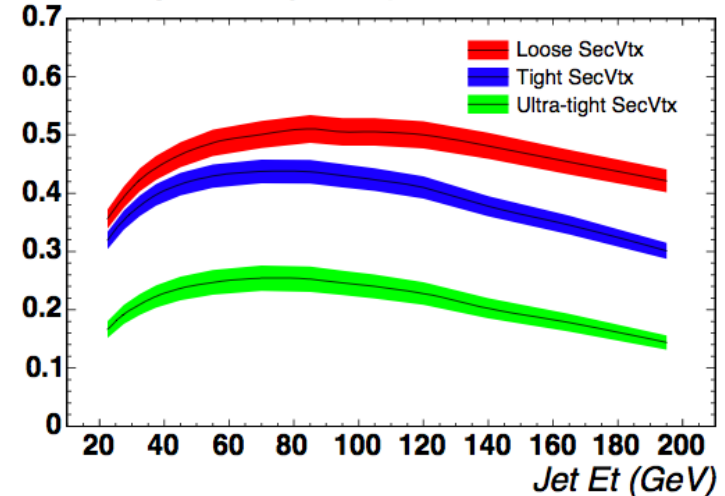


Tagging Efficiencies

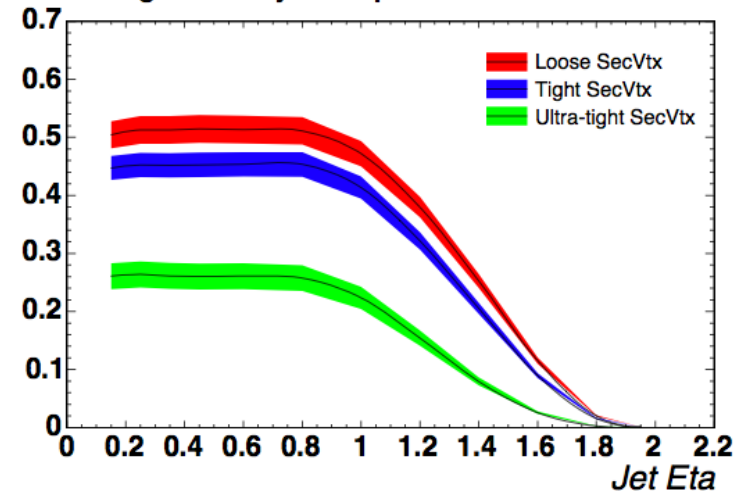
- **Tagging efficiency difficult to model via simulation**
 - Requires excellent knowledge of tracking resolution & efficiency
 - **Strategy:**
 - > Measure efficiency and “mistag” rates in data
 - Inclusive electrons and muons
 - Estimate b quark fraction
 - Tag fully reconstructed Bs
 - > Compare with simulation & compute a scale factor
 - $SF = \epsilon_{\text{data}} / \epsilon_{\text{MC}} \sim 0.95 \pm 0.05$ for “tight” SECVTX



SecVtx Tag Efficiency for Top b-Jets



SecVtx Tag Efficiency for Top b-Jets

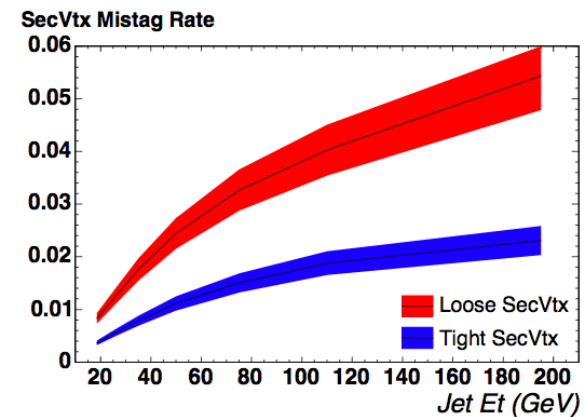
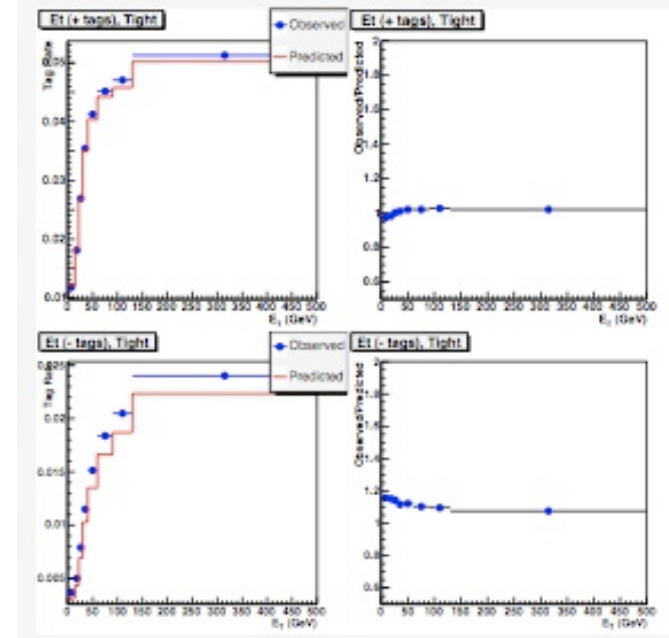


Tagging Fake Rates

- **B tagging fake rates measured from data**
 - Take samples of dijet data, and then create a “fake matrix”
 - > Function of 6 variables
 - > Measure both +ve and -ve tag rates for “taggable jets”
 - Use -ve tag rates as mistag rate
 - > Apply mistag rate to the jets in data sample before tagging

Variable	Bin Edges
Jet E_T (GeV)	0, 15, 22, 30, 40, 60, 90, 130, 1000
Num Tracks/jet	0, 1, 2, 3, 4, 5, 6, 7, 8, 10, 13, 100
$ \eta_{\text{jet}} $	0.0, 0.4, 0.8, 1.1, 2.4
$n_{\text{PrimaryVertex}}$	1 – 6
$\sum E_t$ (GeV)	0, 80, 140, 220, 1000
z_{prim} (cm)	-25, -10, 10, 25

Table 1: Variables and binning used in the mistag matrices.



Example: Quark Substructure

■ Search for quark substructure a long-standing tradition at high energies

- Eichten, Lane & Peskin
 - > PRL 50, 811 (1983)
- Introduced “contact term” Λ_C
- CDF obliged in 1996
 - > $\Lambda_C \sim 1.6$ TeV

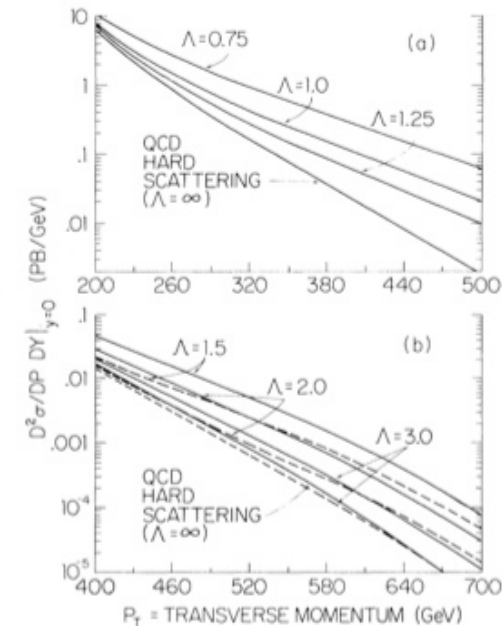
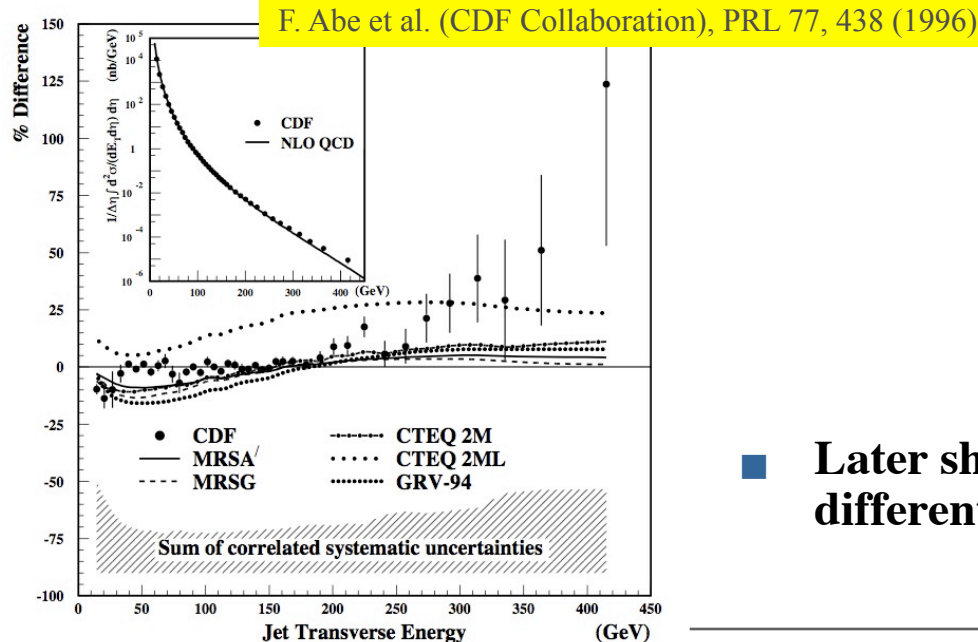


FIG. 3. The jet production cross section (in picobarns/gigaelectronvolt) at rapidity $y = 0$ vs transverse momentum at $\sqrt{s} = 2$ TeV in (a) $\bar{p}p$ collisions and (b) pp collisions for various Λ (in teraelectronvolts). The solid and dashed lines in (b) refer, respectively, to the plus and minus signs in Eq. (5). As a result of a cancellation near $y = 0$, the interference is negligible in (a).

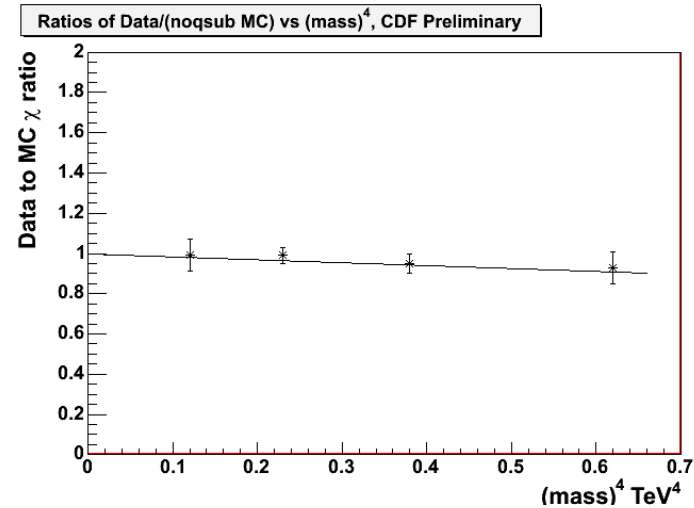
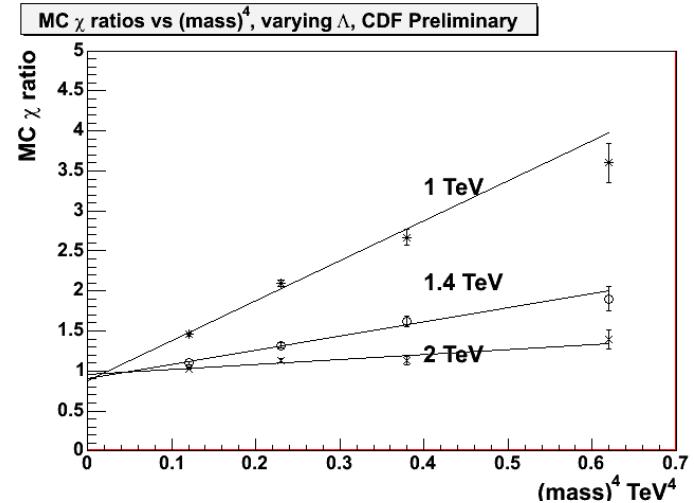
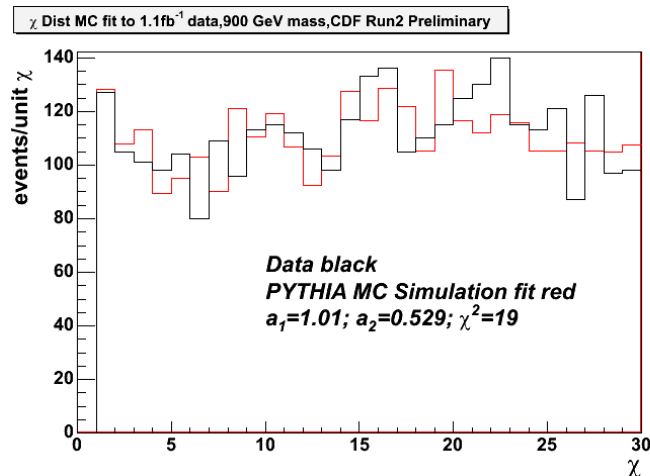
■ Later shown to be described by different PDF behaviour at large x

More Sensitive Study

- Employ angular distribution in dijet scattering

$$\chi \equiv \exp|\eta_1 - \eta_2|$$

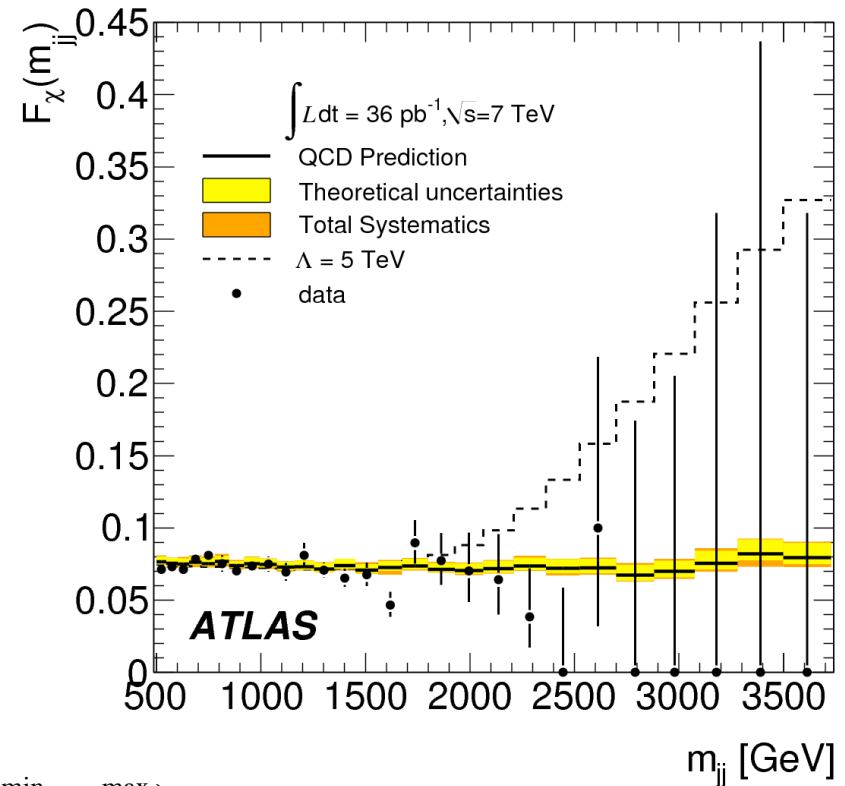
- Look at this as a function of dijet invariant mass
 - > 100 GeV mass bins
- More sensitive to Λ_C
 - > Less sensitive to PDFs
 - > $\Lambda_C > 2.4$ TeV at 95% CL



And Even More Sensitive!

- **ATLAS has further improved sensitivity**
 - Look at fraction of centrally produced jets relative to larger angular range
 - See behaviour as dijet mass increases
 - Expect QCD background to have flat ratio
- **More sensitive to Λ_C**
 - $\Lambda_C > 9.5$ TeV at 95% CL

ATLAS Collaboration, New Jounr Phys. 13,053004 (2011)



$$F_{\chi}([m_{jj}^{\max} + m_{jj}^{\max}] / 2) \equiv \frac{N_{\text{events}}(|y^*| < 0.6, m_{jj}^{\min}, m_{jj}^{\max})}{N_{\text{events}}(|y^*| < 1.7, m_{jj}^{\min}, m_{jj}^{\max})}$$

$$y^* \equiv \frac{1}{2}(y_1 - y_2)$$